

The QCD Equation of State with Thermal Properties of phi-mesons

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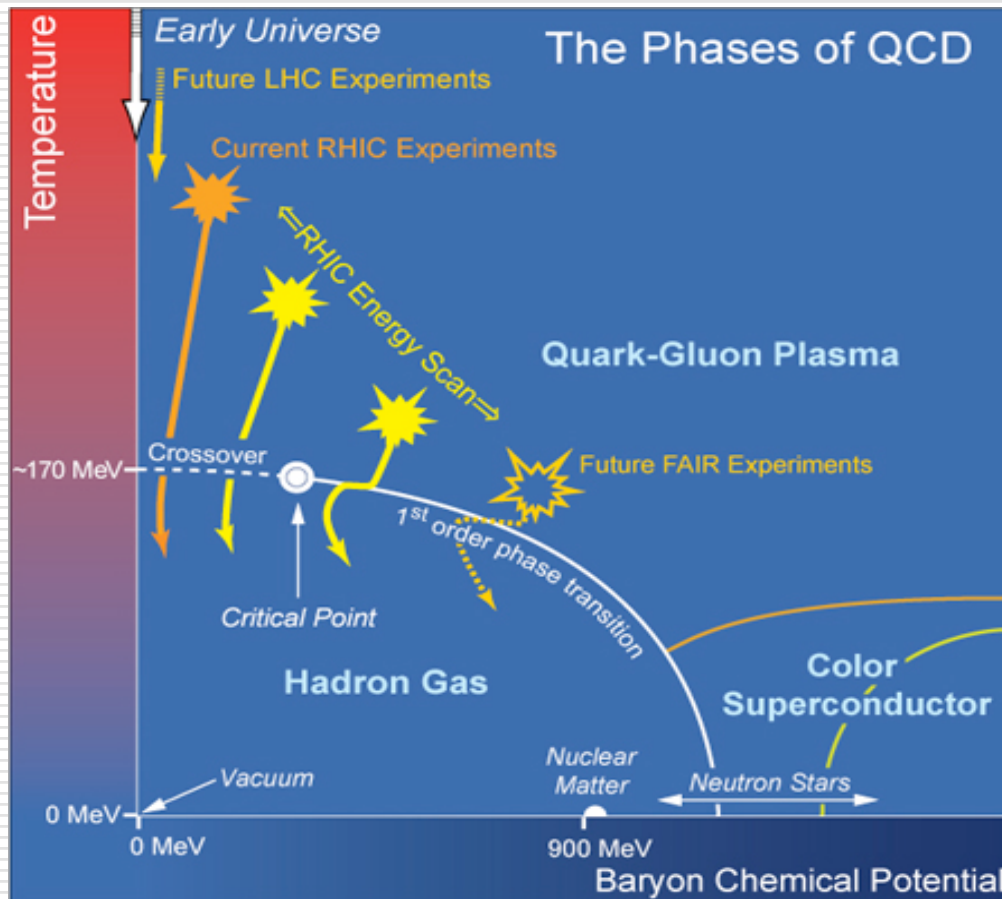
Excited QCD 2009, 8-14 February 2009, Zakopane, Poland

Outline

- ✓ Introduction
- ✓ Results:
 - ✓ Van Hove Signal and order of transition with Φ
 - ✓ Lattice QCD vs data using the Φ -meson
 - ✓ Approach to T_c
- ✓ Summary and outlook

Introduction

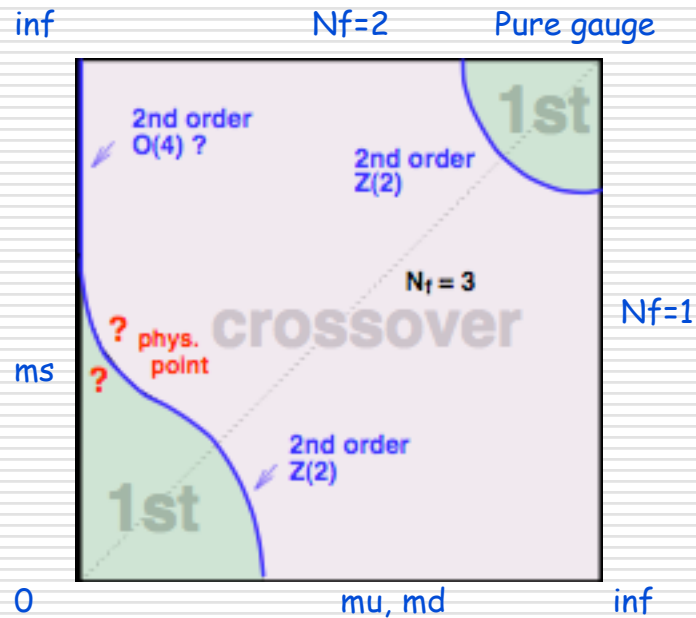
The QCD Phase Diagram



Existence of critical point not completely clear from $\mu_B > 0$ lattice calculations, Allton et al, hep-lat/0501030

The order of the transition changes with μ_B , therefore with the \sqrt{s} of the collision

Order of transition from lattice QCD at $\mu_B=0$

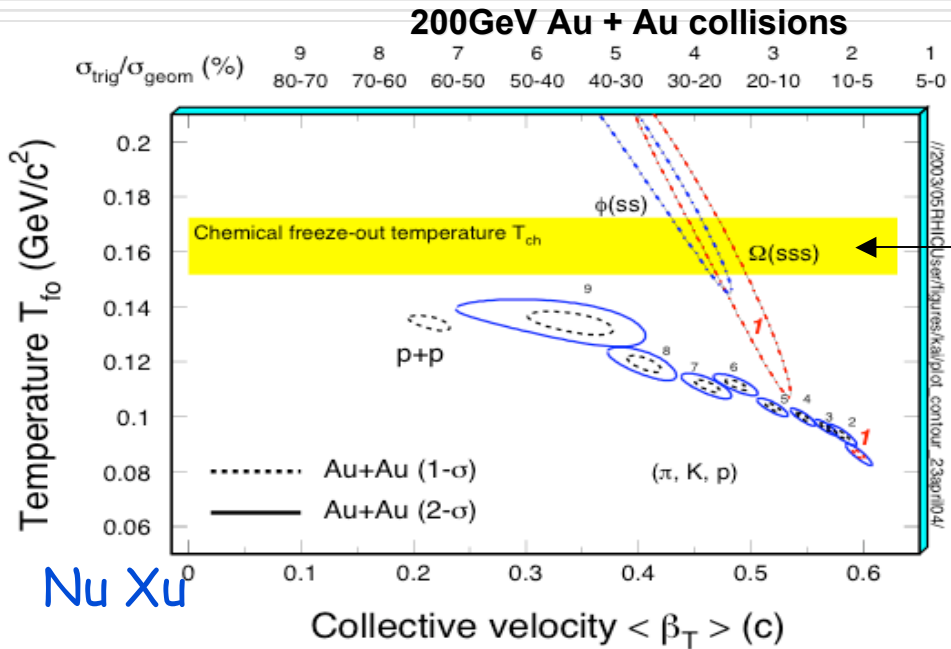


A Peikert, PhD Thesis, Univ. Bielefeld

Order of transition from lattice QCD at $\mu_B=0$ varies with nr of flavours and masses of quarks

--> Address the order of transition using data vs theory

Why the phi meson ?



Small hadronic rescattering cross-section, $\sigma(\phi N) = 10$ mb.

ϕ (and Ω) decouple early: they show larger T_{fo} and smaller $\langle \beta_T \rangle$ than other hadrons (see 1 for central collisions on the left fig.)

ϕ T_{th} (thermal fr. out) $\sim T_{ch}$ (chem. fr. out) within errors

life time = 45 fm/c \rightarrow decay products ($\phi \rightarrow K+K$ -channel) are not influenced much by rescattering

$\rightarrow \phi$ (as well as Ω) are interesting probes for studying EoS and nature of phase transition

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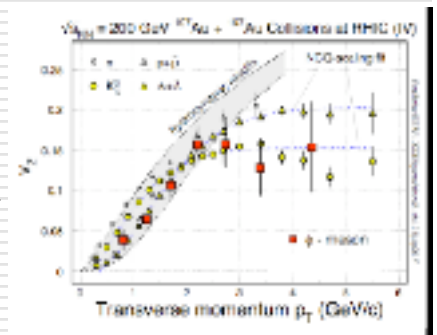
STAR: NPA715, 458c(03); PRL 92, 112301(04); 92, 182301(04).

Observables for phi:

T_{th} (thermal freeze out) : reflect information below or near T_c

Yields, T_{eff} , flow eg v_T (radial flow), v_2 (elliptic flow) etc: can reflect information below and above T_c (e.g. suggested by n_q scaling of v_2)

\rightarrow Here we attempt to use phi data and confront them to lattice and simple thermodynamic expectations like the Van Hove signature



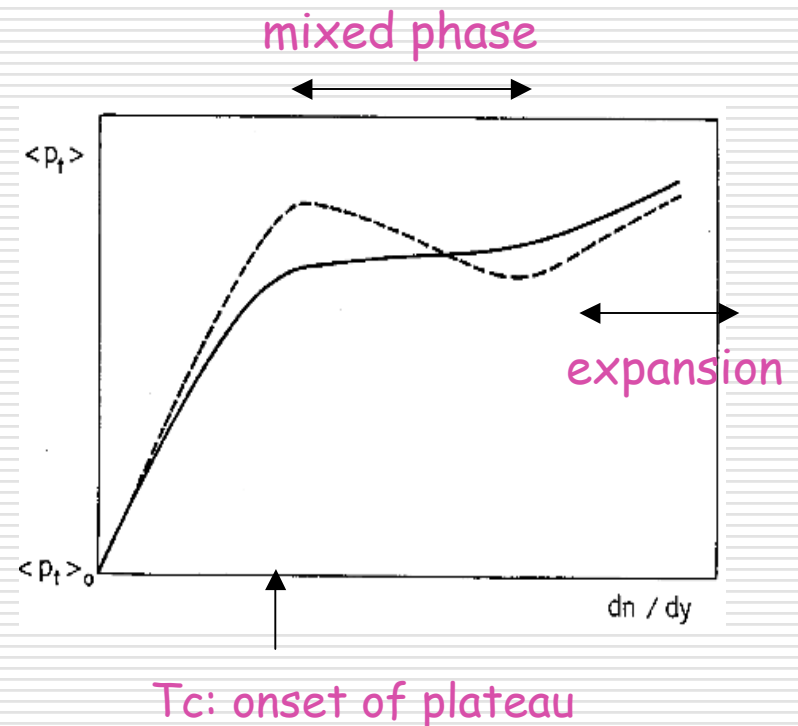
The Van Hove Signal of 1st order Phase Transition

L. Van Hove:

Phys. Lett. B 118, 138 (1982)

The flattening of the $\langle p_T \rangle$ with increasing multiplicity in proton-anti-proton collisions, at mid-rapidity, may serve as a probe for the equation of state of hot hadronic matter.

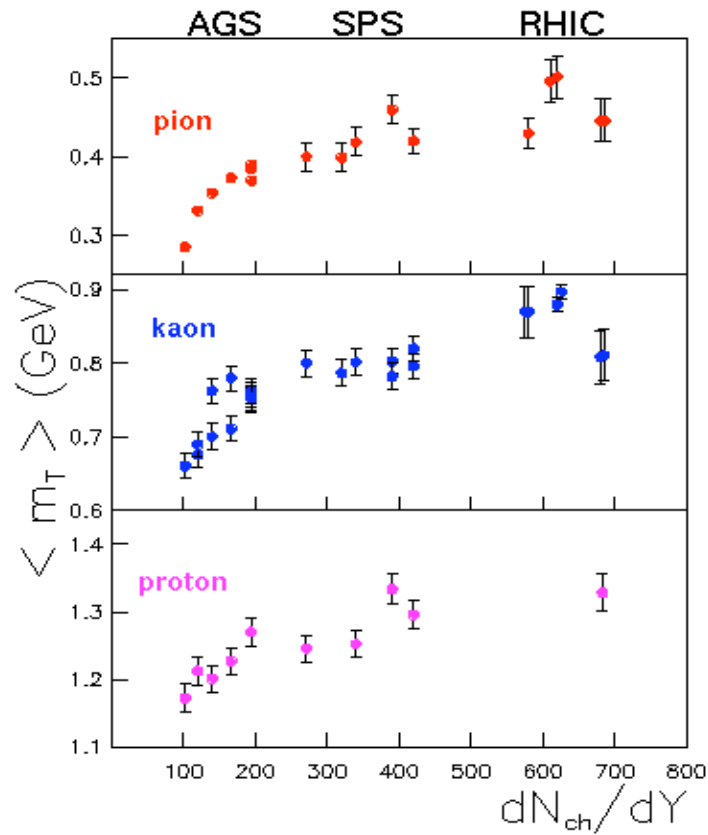
Multiplicity reflects the entropy and $\langle p_T \rangle$ the effect of temperature and transverse expansion.



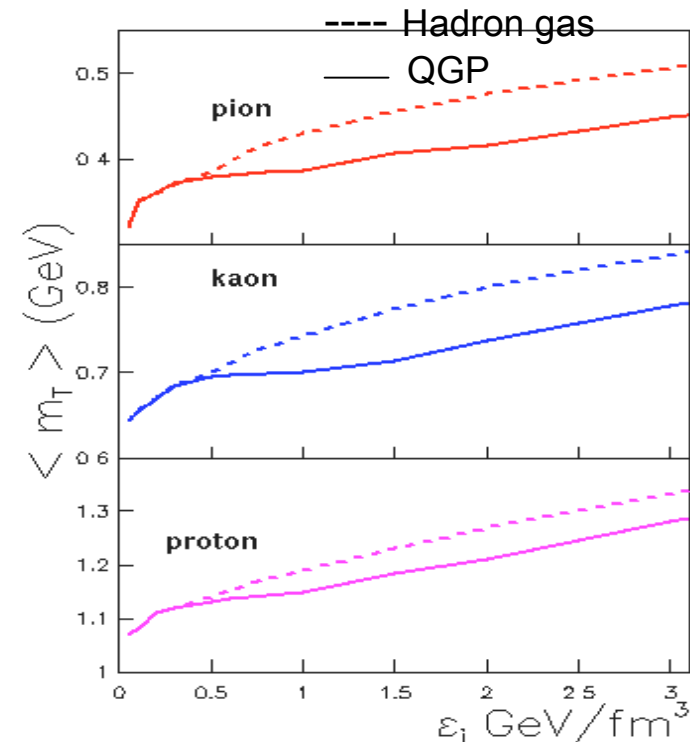
The observation of an increase, a central plateau, followed by a second increase in $\langle p_T \rangle$ as a function of dN/dy is suggested as signature of latent heat (1st order phase transition).

Van Hove Signal in π, k, p

Mohanty, Alam, Sarkar, Nayak, Nandi
 PRC 68 (2003) 021901 (R)



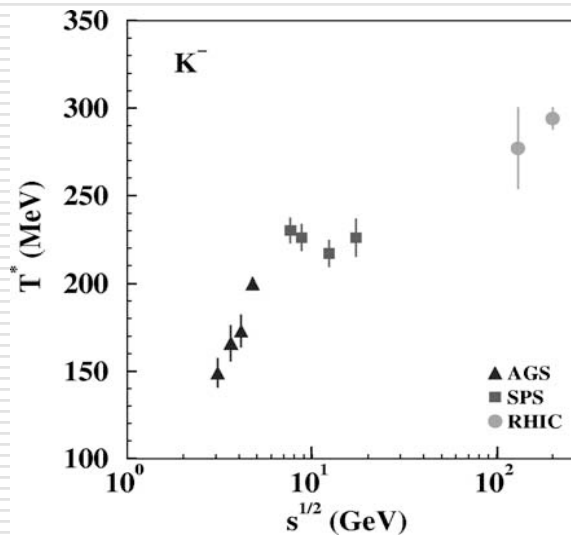
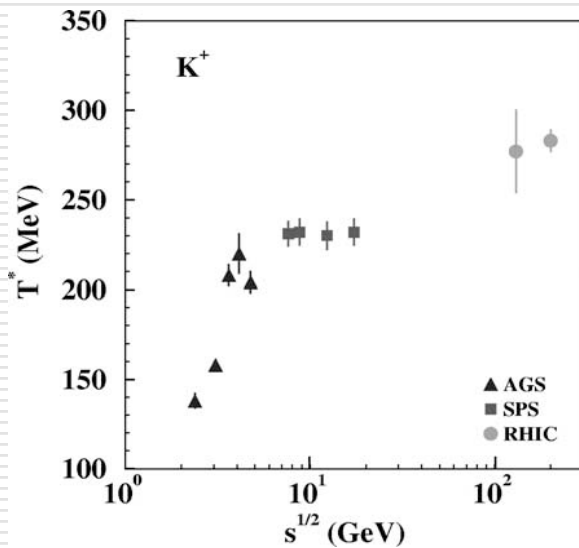
Hydrodynamic calculation with
 (3+1) dimensional expansion



- Rapid rise in $\langle m_T \rangle$ going from AGS to SPS energies and indication of flattening towards RHIC.

Hydrodynamic calculation with QGP reproduces the general trend of data
 Observation of a central plateau indicates 1st order phase transition

Collision energy dependence of inv. slope of K[±]



M. Gazdzicki et al

PLB 567 (2003) 175

T^* = inv. slope = T + flow
instead of $\langle pt \rangle$

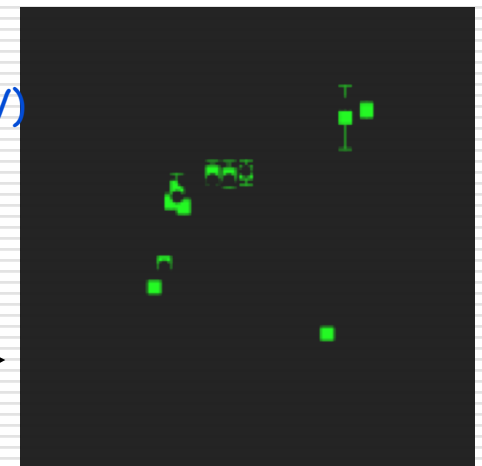
\sqrt{s} instead of
 dN/dy

Indication of phase co-existence (1st order phase transition) in Kaon production in A+A collisions.

Hydrodynamic calculation with 1st order phase transition can reproduce these data.

Y. Hama et al, Acta Phys Pol B, 35, (2004).

T^*
(MeV)



\sqrt{s} (GeV)

Results I

Van Hove Signal and order of transition with Φ

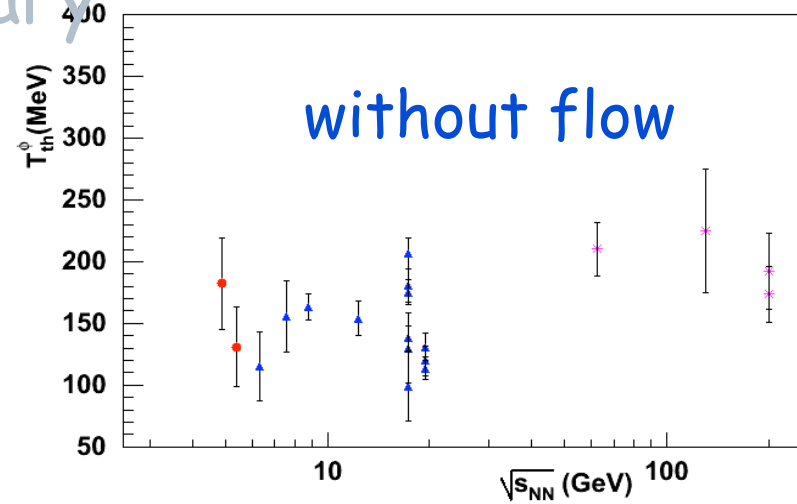
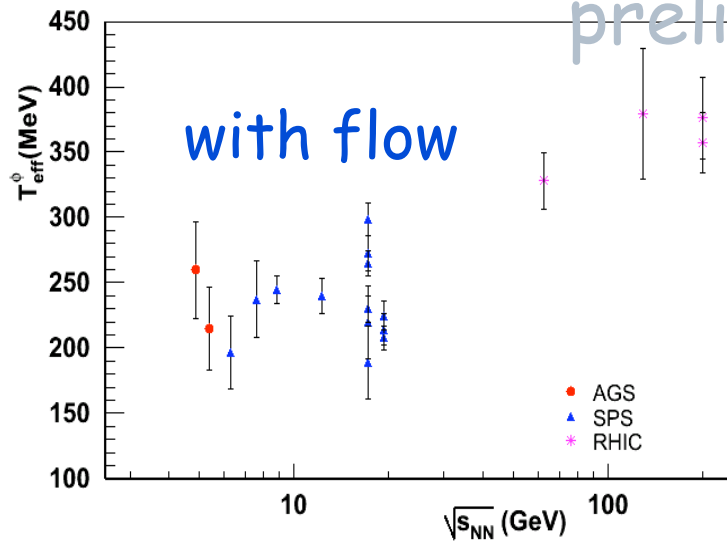
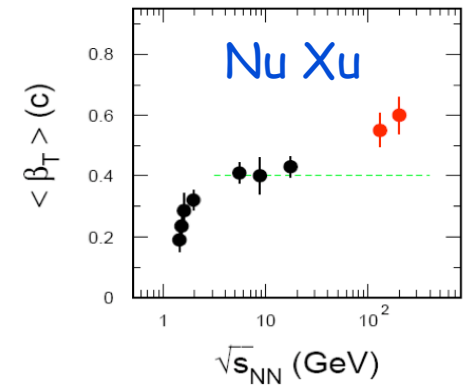
Van Hove Signal in phi-meson : \sqrt{s} dependence

T_{eff} : inverse slope

T_{th} ; T (thermal freeze out)

$$T_{eff} = T_{th} + 0.5 m_{\phi} \langle \beta_T \rangle^2,$$

$\langle \beta_T \rangle$: average coll. radial flow velocity



T_{eff} of phi shows a plateau and increases at $\sqrt{s}=62$ GeV --> indicating a possible 1st order transition

T (th. fr. out) of phi at RHIC exhibits high values, while an enhancement remains after flow subtraction.

T_{th} is approx. near T (chem. fr. out).

Van Hove Signal in ϕ -meson: ϵ (Bjorken) dependence

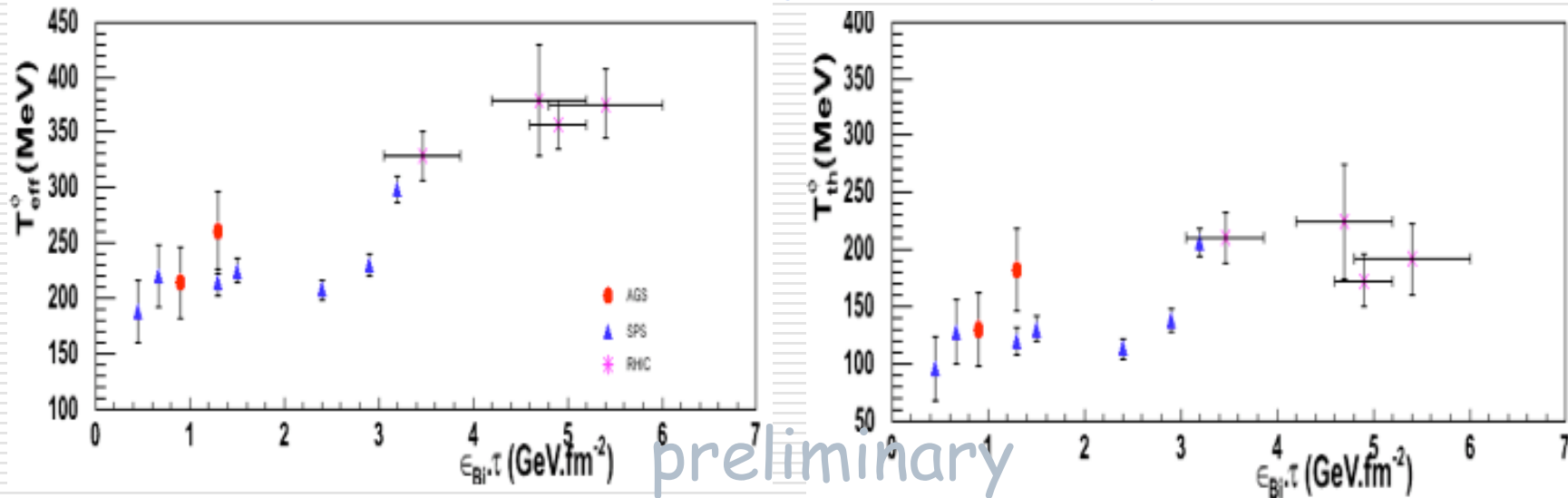
T_{eff} : inverse slope

T_{th} ; T(thermal freeze out)

$e(Bj) = (dE_t/dy) / (\pi R^2 \tau)$

$T_{eff} = T_{th} + 0.5 m_\pi \langle \beta_T \rangle^2$, $\langle \beta_T \rangle$: average coll. radial flow velocity

1 fm/c formation time



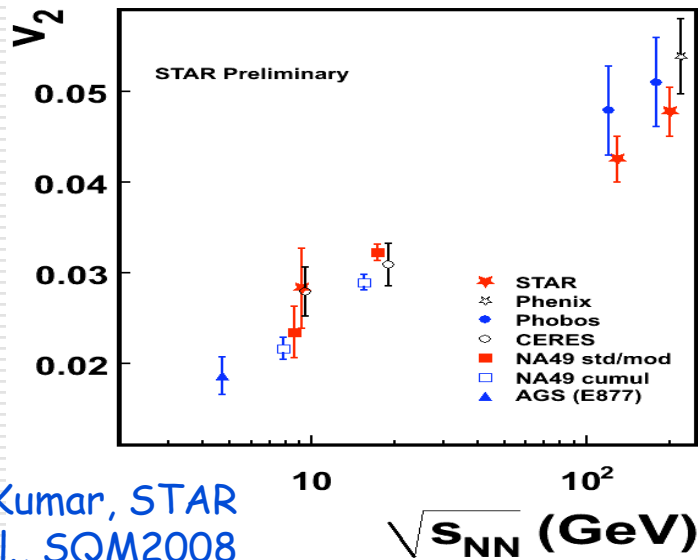
T_{eff} of ϕ shows a plateau below 3 GeV/fm^3 , and subsequent increase above $\sim 3.2-3.8 GeV/fm^3$.

More ϕ data are needed below 1 GeV/fm^3 to assess the increase and onset of plateau \rightarrow or one can look also other particles than ϕ for this.

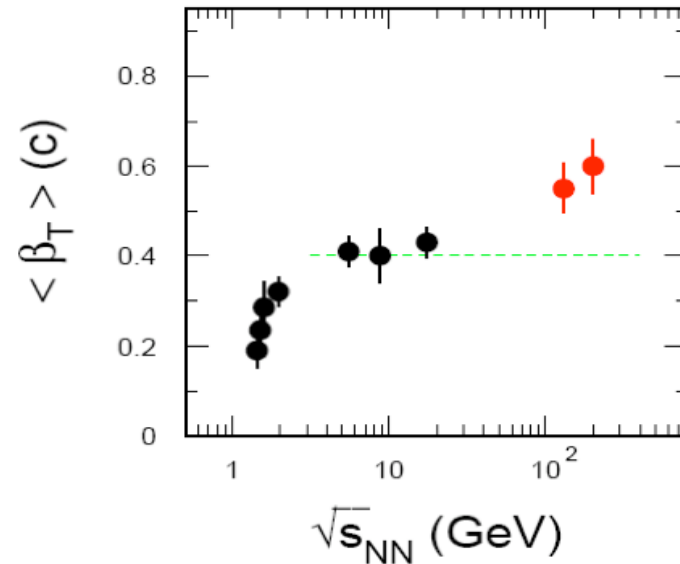
$T(th)$ shows a plateau, while a small increase is still observed above 3 GeV/fm^3 after radial flow is subtracted.

Note that taking the measured $\beta_T(\phi)$ will enhance the $T(th)$ ϕ points at RHIC (work in progress).

Van Hove signal in flow



L. Kumar, STAR coll., SQM2008



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- The central plateau is seen also in parameters characterizing collective expansion namely in β_T as a function of \sqrt{s} .
- The v_2 as a function of \sqrt{s} may indicate a plateau structure too. Need more data and plot versus $e(B_j)$ will be useful.
- Those parameters may reflect the initial pressure, the latter expected to remain constant during mixed phase in 1st order transition.

[For v_T hadronic rescattering phase may be important: Y. Hama et al, Acta Phys Pol B, 35, (2004).]

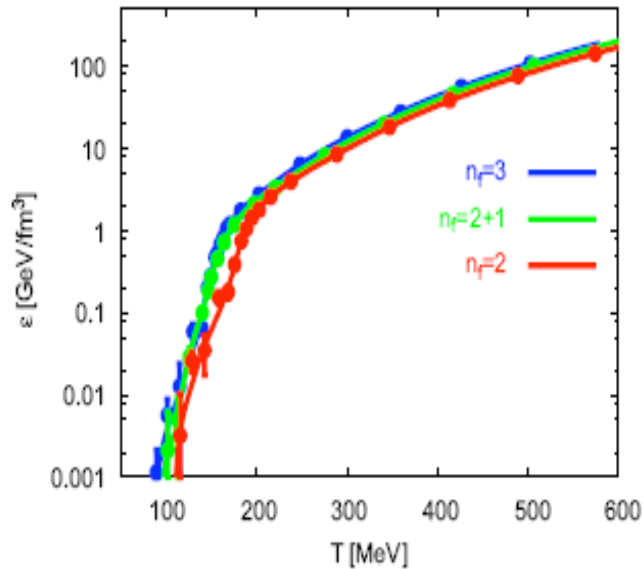
Results II

Lattice QCD vs data using the Φ -meson

Lattice QCD & EoS: F-meson

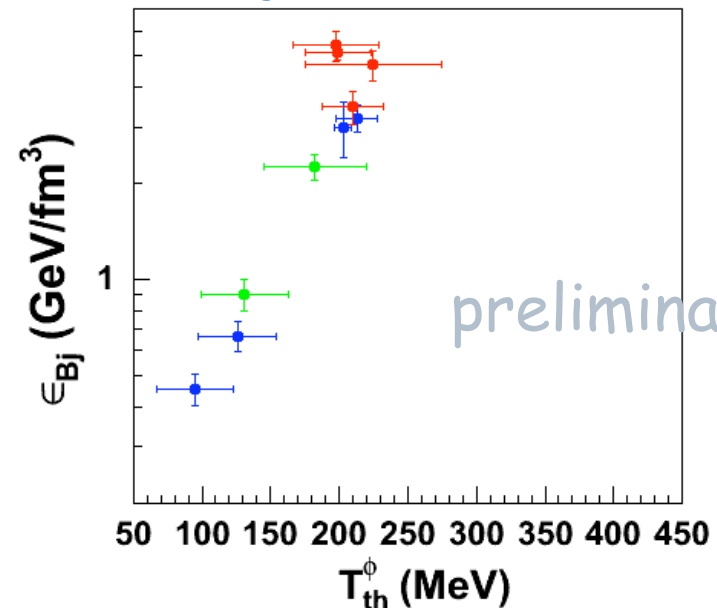
The energy density as a function of the T

D.E. Miller / Physics Reports 443 (2007) 55–96



T is about 192 MeV and density about 1 GeV/fm³ where the discontinuity in slope occurs. (cross over at $\mu_B=0$)

Assuming $t = 1 \text{ fm}/c$



Energy density combines the effects of the initial conditions which are controlled by the impact parameter the nuclei used, and the collision energy.

- $T_{th}(\phi)$ is hadronic thermal freeze out T, probably near chemical freeze out T, and below or at T_c .

- $e(B_j)$ and T_{th} in data (right) are 'measured' at different times

---> Note that the curve is artificially approaching T_c .

(Because the high e points correspond in reality to T_{init} of QGP above T_c)

--> To overcome this we use as next a 'measure' of T_{init} or Pressure($init$) from data

Lattice QCD & EoS: F-meson

The energy density as a function of the T_{eff}

D.E. Miller / Physics Reports 443 (2007) 55–96

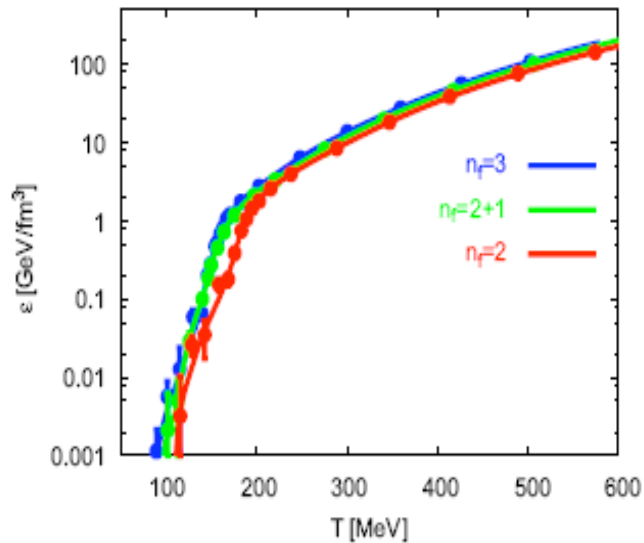
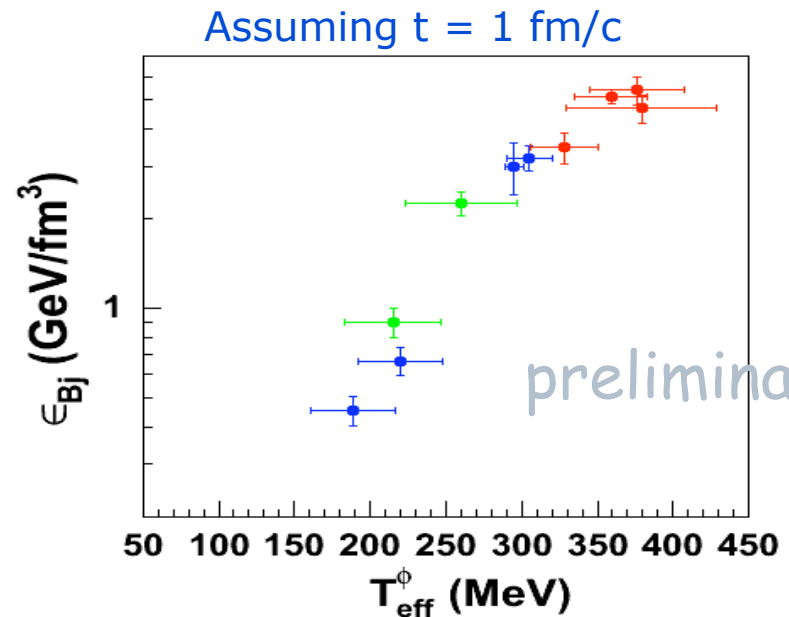


Fig. 8. The energy density for the three values of $n_f = 2, 2 + 1, 3$ with massive dynamical quarks.



Since $T(\text{init})$ QGP above T_c , is not measurable with ϕ , we use T_{eff} in order to reflect the behaviour of the initial state variables which can be above T_c , and in particular of the pressure (flow) term in T_{eff} .

T_{eff} is not a temperature but reflects the initial energy ($T+\text{flow}$).

(In the right plot, e and T_{eff} from data are not measured at the same time. In lattice they are.)

--> Some similarity is exhibited in lattice vs data.

--> More data needed (especially in low T_{eff}).

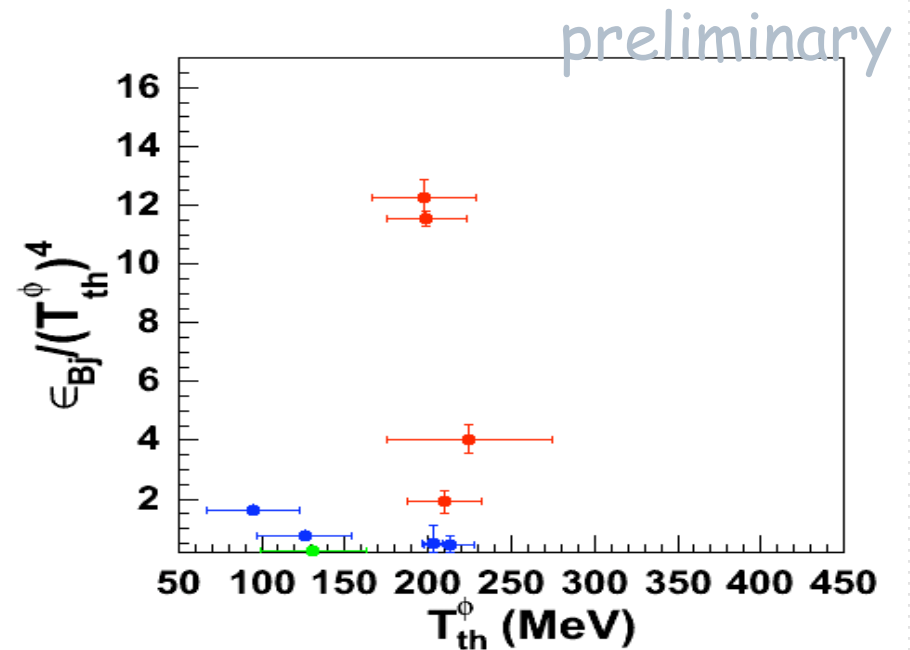
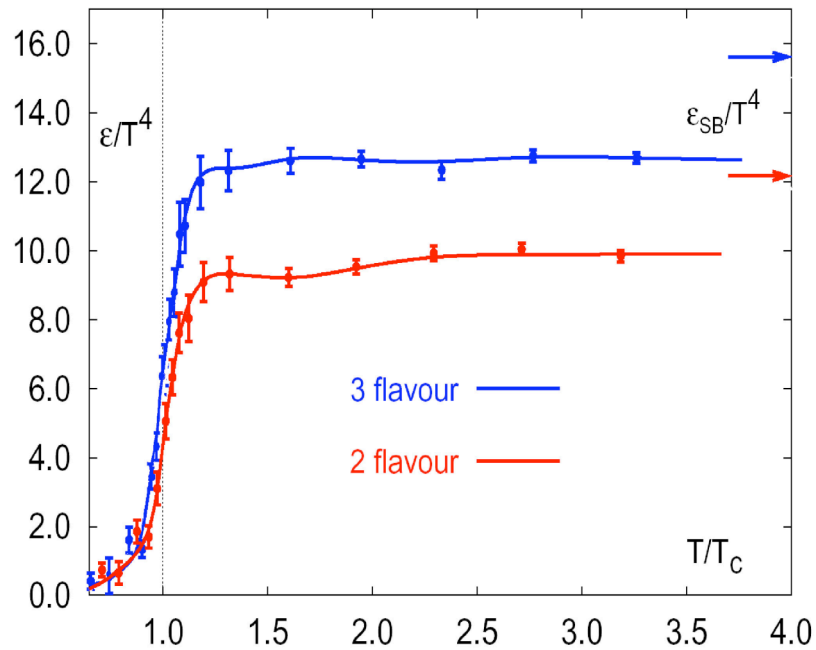
→ Use hydro for this study

→ study also for $e(B_j)$ versus v_2 and for $e(B_j)$ vs β_{T_2}

(work in progress)

Lattice QCD & EoS: phi-meson

The energy density/ T^4 as a function of the T



formation time parametrized as a function of E
y axis arbitrarily normalized

T_{th} is below or at T_c , even if $e(B_j)$ is above T_c , causing an artificial increase approaching T_c .

To access the $T > T_c$ part, one can try to use T_{eff} , reflecting qualitatively properties of the $T > T_c$ range.

Results III

Lattice QCD vs data:

Approach to T_c using $T(\text{chem. freeze out})$ of all hadrons at $\mu_B=0$

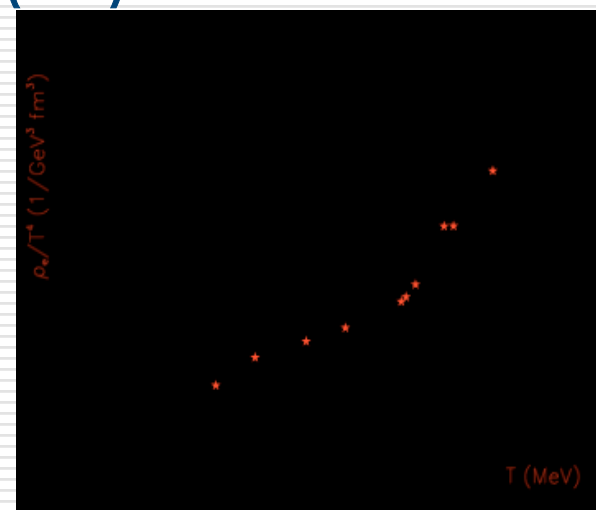
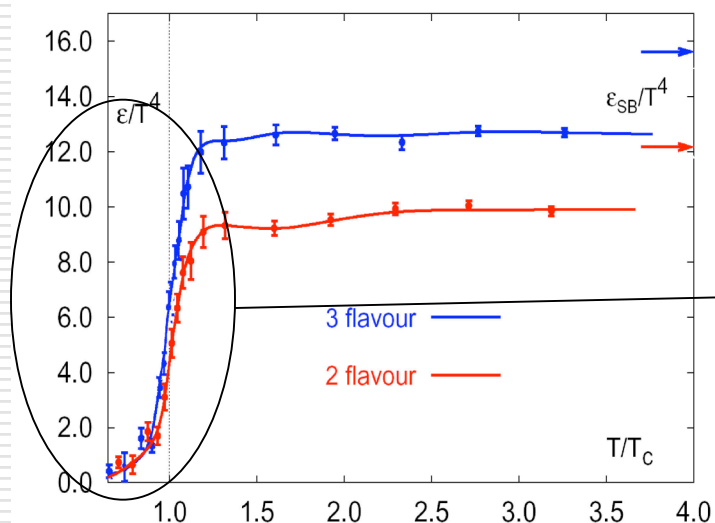
Lattice QCD & EoS:

The energy density/ T^4 as a function of the T at $\mu_B=0$,
at the chemical freeze out

We now look at the e/T^4 vs T dependence, using e, T at the chemical freeze out extracted with a thermal model from particle ratios, and extrapolated to $\mu_B=0$

S. Kabana, P. Minkowski, New J Phys 3:4, 2001, hep-ph/0010247.

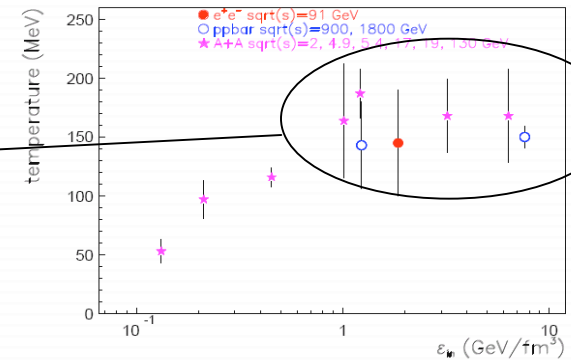
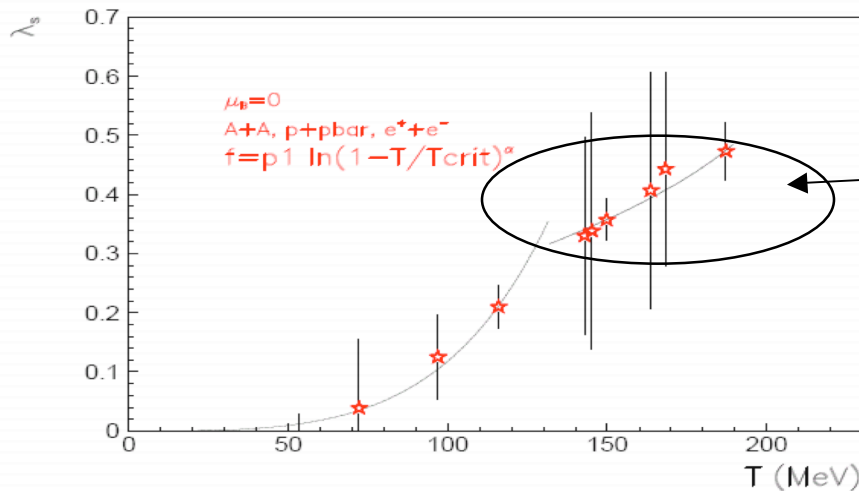
S. Kabana, hep-ph/0104001, Eur. Phys. J. C 21 (2001) 545



The e, T here are measured at same time, namely at chem. freeze out,
and $\mu_B=0$

The increase of e/T^4 , approaching T_c is observed

An attempt to extract a critical exponent



**S. Kabana, hep-ph/0104001,
 Eur. Phys. J. C 21 (2001) 545**

First study of the approach of λ_s to the Temperature, both extracted from data using a thermal model, for the colliding systems which are near T_c (SPS and RHIC) is fitted by a function

$$F = c |\ln(1 - T/T_c)|^a$$

We find $T_c = 218 \pm 70 \text{ MeV}$ and $a = 0.54 \pm 0.47$

Despite the large errors one can observe that the T_c value is within the error in agreement with previous T_c estimates

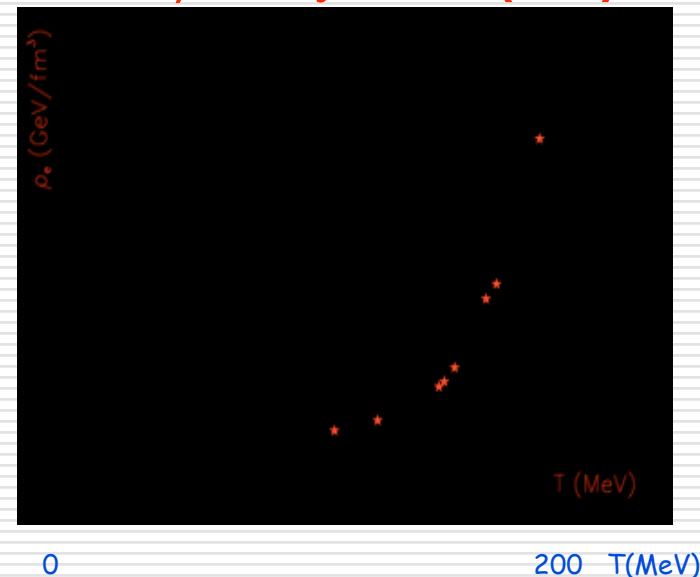
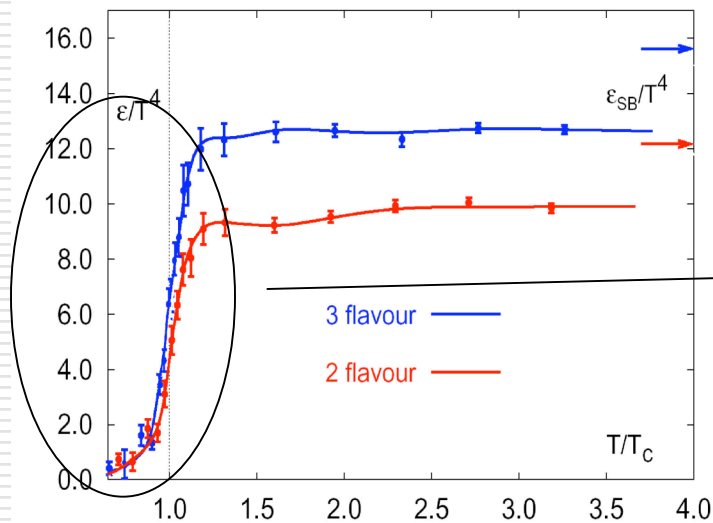
(Errors can be improved with new data)

Lattice QCD & EoS:

The energy density as a function of the T at $\mu_B=0$ and at chemical freeze out

e vs T dependence, using e , T at the chemical freeze out, and at $\mu_B=0$ (as a normalization to same μ_B for all collisions, and to compare to $\mu_B=0$ lattice results)

S. Kabana, hep-ph/0104001, Eur. Phys. J. C 21 (2001) 545



The e , T here are measured at same time, namely at chem. freeze out, and $\mu_B=0$

The increase of e approaching T_c is observed, and can be studied in terms of a critical exponent. Same can be done for ρ_s , ρ_h .

Summary and outlook

- ❖ The ϕ shows the Van Hove signature in the inverse slope ($= T + \text{flow}$) versus $e(B_j)$ or \sqrt{s} dependence (like π, K, p too).
- ❖ Hydrodynamic calculations for π, K, p assuming 1st order transition reproduce the plateau and the latter enhancement of the slopes.
- ❖ Flow (and therefore T_{eff} too) could reflect initial conditions above and below T_c .
- ❖ β_T shows a signature of a 1st order transition a la van Hove, in terms of a equivalent of pressure P : $\beta_T = f(P)$ vs $e(B_j)$.

In particular the 'pressure'-equivalent remains \sim constant, while $e(B_j)$ increase from 1 to 3 GeV/fm^3 .

Lattice calculations suggest that the order of the transition changes with μ_B and therefore along the data points shown versus \sqrt{s} or e .

- ❖ The above observations suggest a mixed phase up to $e(B_j) = 3.2 - 3.8 \text{ GeV}/\text{fm}^3$, where the plateau seems to end.
- ❖ This does not exclude the change of the order at higher \sqrt{s} (RHIC), as predicted by lattice towards $\mu_B = 0$

Summary and outlook

- ✓ The thermal properties of the phi-meson have been used to compare data to lattice calculations and to study the nature of the QCD phase transition.
- ✓ The $e(B_j)$ vs $T_{\text{eff}}(\text{phi}) = T_{\text{th}} + \text{flow}$: both variables may relate to initial observables below and above T_c --> Comparison to lattice QCD (preliminary results).
- ✓ Study **approach to T_c from below** and near T_c , and at $\mu_B=0$: e , T measured at same time namely at chemical freeze out from hadron ratios. **Extract critical exponents.**

Outlook :

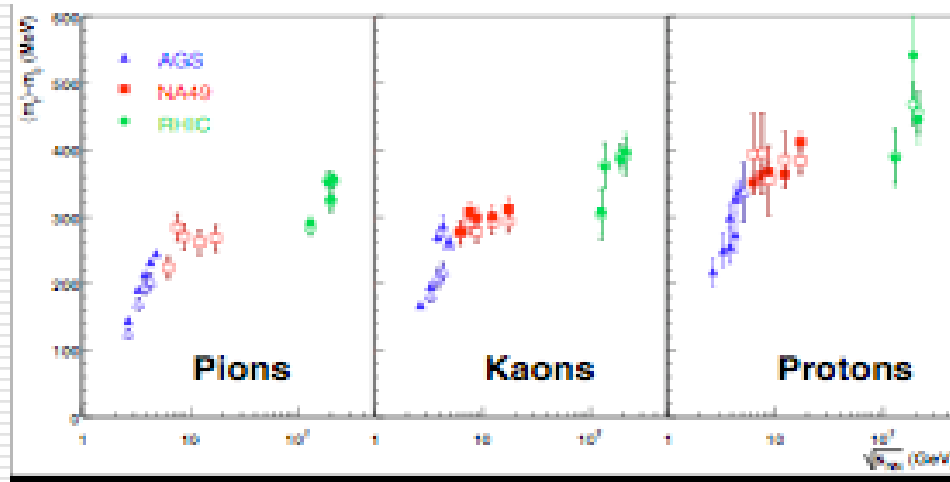
- ✓ Work in progress with hydrodynamic calculations for the phi.
- ✓ Study approach to T_c (for $e=1-3 \text{ GeV}/\text{fm}^3$, $e > 3 \text{ GeV}/\text{fm}^3$)
- ✓ Use $\beta_T(\text{phi})$ for $T(\text{th})$ phi estimate
- ✓ Use more phi data, especially at low \sqrt{s} if available.
- ✓ Use v_2 , v_T for qualitative lattice comparisons.

Thank you very much !

BACKUPS

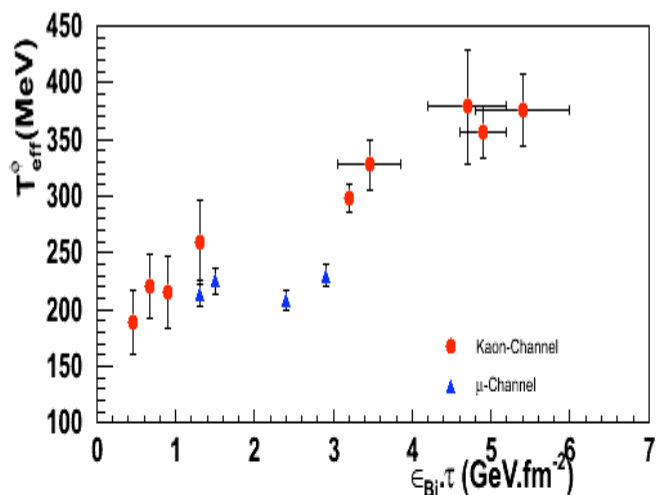
Van Hove signal in pi,K,p NA49 plot

Central AuAu or PbPb collisions



Ch. Blume, J Phys G31 (2005) 57.

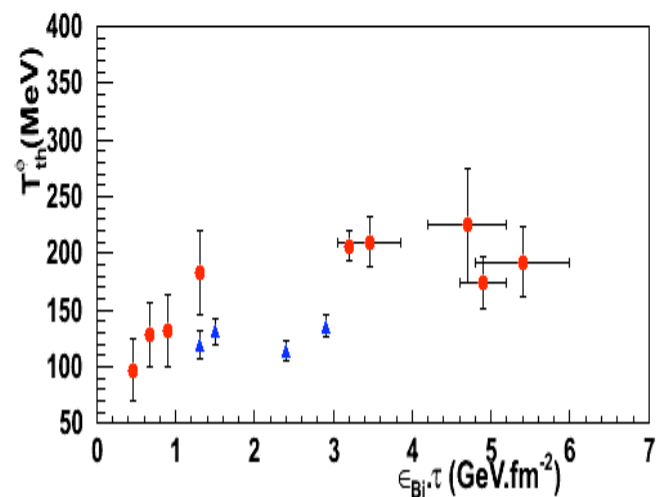
Van Hove Signal & Φ -puzzle



Φ is reconstructed from its hadronic and leptonic decay channels.

The puzzle:

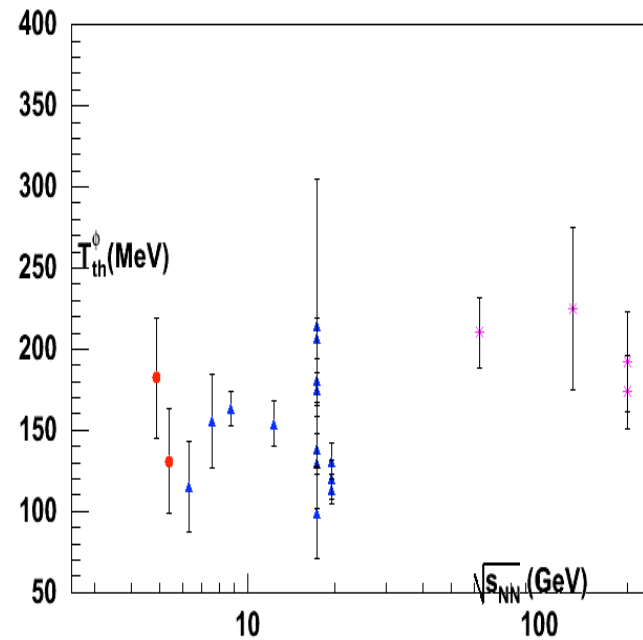
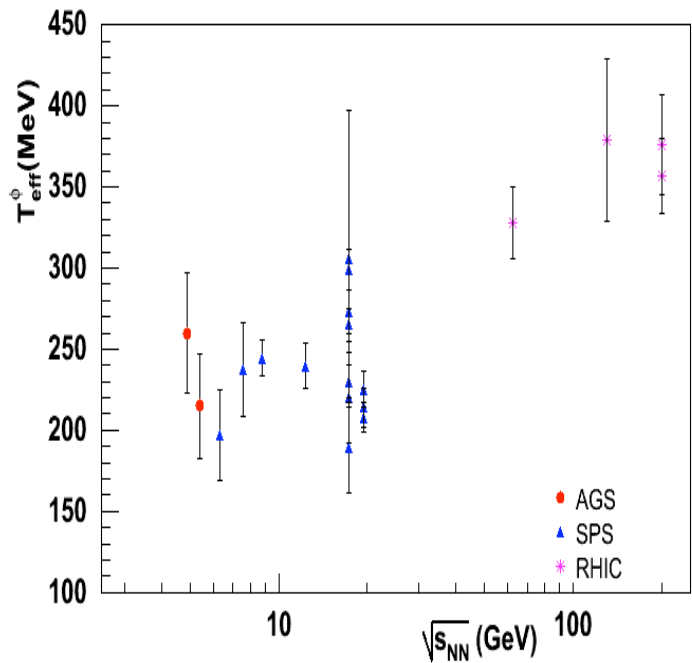
- ❖ the yield is lower for hadronic channel compared to the leptonic channel of reconstruction
- ❖ the inverse slope parameter T_{eff} is higher for hadronic channel



Does Φ -puzzle make the central plateau?
However this plateau is seen in π , K , p too

Note: Φ Thermal freeze out temperature near $\sim T_c$

Backups



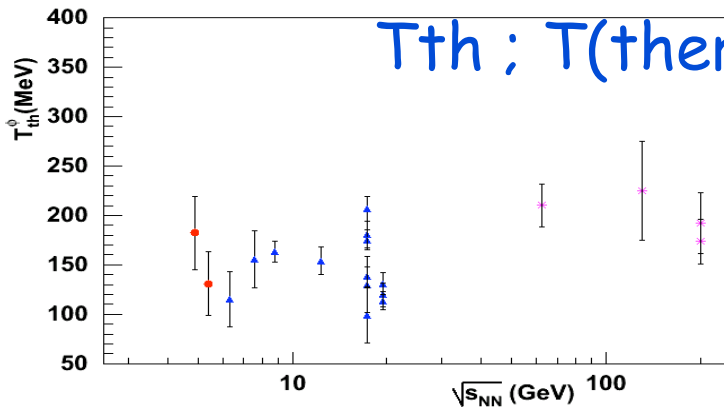
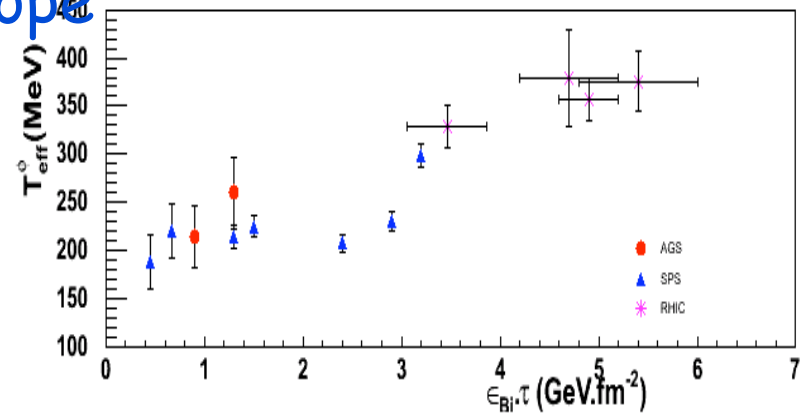
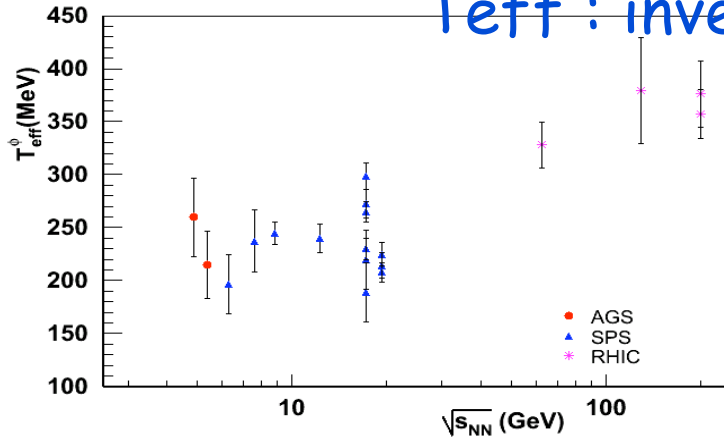
Inclusion of CERES leptonic channel data which has error 92 MeV

There plots are with new NA49 and CERES data at SPS.

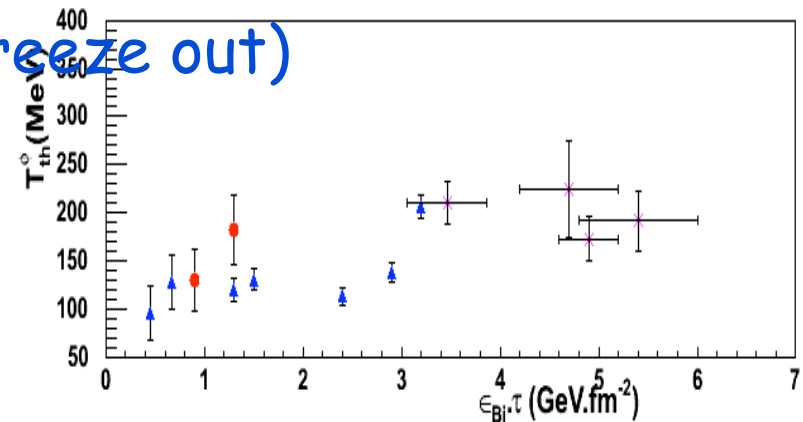
Van Hove Signal in phi: $e(Bj)$ and \sqrt{s}

T_{eff} : inverse slope

$$T_{eff} = T_{th} + 0.5 m_{\phi} \langle \beta_T \rangle^2$$



T_{th} ; T (thermal freeze out)

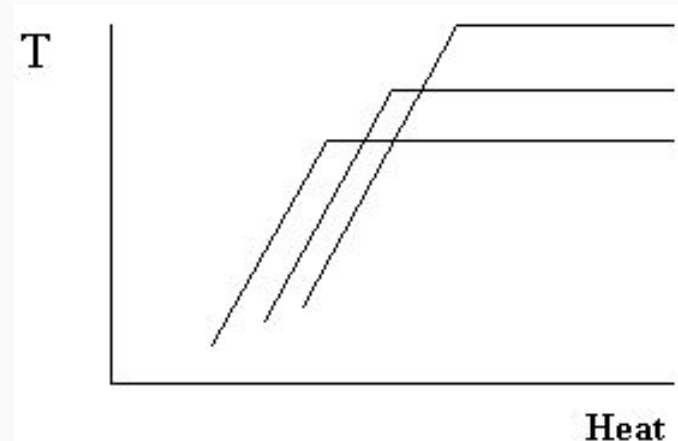
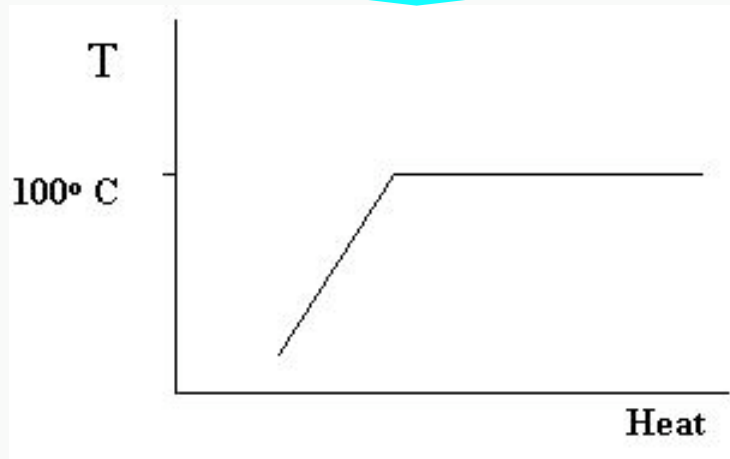


$e(Bj)$ dependence improves from \sqrt{s} in taking also into account effects like the different nuclei and impact parameters

Gap of \sqrt{s} between SPS and RHIC is less pronounced on $e(Bj)$ scale

Gedanken experiment to identify the water steam phase transition:

We heat a box with water more and more and measure its temperature T . We can only measure the T of the water (Had. Gas) and not of the steam (QGP). We plot T versus heat. T will rise until we heat enough to reach $T=100^\circ\text{C}$. From then on, it will remain the same, namely $T_{\text{lim}} \sim 100^\circ\text{C}$. Each time steam is present, we have to wait until it is again water, to measure its T . (E.g. R.Hagedorn (1965), H. Stocker et al (1981) etc.)



Now we repeat the experiment adding each time salt to the water. The T versus heat curve will not be as before, and we can not find the $T_{\text{lim}} = 100^\circ\text{C}$.

S. Kabana, J. Phys. G27 (2001) 497

The baryochemical potential is like salt for hadronic systems.

Therefore, in order to measure a unique curve of T at freeze-out as a function of ϵ (init) in hadronic particle systems, one has to use the same conditions, with the same μ_B , the simplest one being $\mu_B=0$.