

Baby-steps beyond rainbow-ladder

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Excited QCD09 – Zakopane, 9/02/09



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Christian S. Fischer, Dominik Nickel, RW, Eur.Phys.J.C,2008 [arXiv:0807.3486]

Christian S. Fischer, RW, Phys.Rev.D78:074006,2008 [arXiv:0808.3372].



Desire:

- Poincaré covariant description of mesons
- formulated in the continuum
- description in terms of fundamental quantities of QCD

Natural framework:

- Bethe-Salpeter equations
- Schwinger-Dyson equations

Studied in detail for many years

- Rarely extended beyond simplest truncations
- Ad-hoc 'improvements' used.
- Severe approximations (e.g. M-N)



Rainbow-Ladder:

- successful description of light-mesons subject to an apposite phenomenological ansatz for the interaction.
e.g. Maris-Tandy model.

[P. Maris, P. C. Tandy, PRC **60** (1999) 055214]

Lacking in many regards:

- no unquenching effects – pion cloud
- no η/η' splitting – $U_A(1)$ anomaly
- admits $(\bar{3})_c$ coloured diquark bound-states – but useful in studying Baryons

[A. Bender, C. Roberts, L.v Smekal, PRLB **380** (1996) 7]



Moreover:

- Describes only pure $q\bar{q}$ -states:
 - No flavour mixing
 - No decay channels
 - No exotics

Rainbow-Ladder provides ONLY

- $\gamma_\mu \otimes \gamma_\mu$ couplings
 - Simplicity of interaction means higher spin states of mesons are poorly represented.
 - No variety in attraction/repulsion.



Goal:

- Consistent Green's function approach
 - Ghost/Gluon solutions of DSE
 - Quark-Gluon vertex beyond γ^μ
 - BSE kernel satisfying axWTI

“Break the ladder:”

- Unquenching effects
- Leading Yang-Mills corrections

- 1 Introduction
 - Bethe-Salpeter equations
 - Schwinger-Dyson equations
 - Rainbow-Ladder
- 2 Quark-gluon vertex
 - Basic structure
- 3 Beyond rainbow-ladder: Unquenching effects
 - Modelling the pion-cloud
- 4 Beyond rainbow-ladder: Yang-Mills sector
 - Gluonic Corrections
- 5 Outlook/Conclusions

Outline

- 1 **Introduction**
 - Bethe-Salpeter equations
 - Schwinger-Dyson equations
 - Rainbow-Ladder
- 2 **Quark-gluon vertex**
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- 3 **Beyond rainbow-ladder: Unquenching effects**
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- 5 **Outlook/Conclusions**

Bethe-Salpeter equations

Bound states:

- poles in $n \geq 3$ -point colour singlet Green's functions

$$\Gamma_H(p, P) = \frac{r_H \Gamma_h(p, P)}{P^2 + m_H^2} + \text{regular terms}$$

$\Gamma_h(p, P)$ solves homogeneous Bethe-Salpeter Equation:



Required inputs

- Quark propagator
- Gluon propagator
- Quark-Gluon vertex
- Scattering kernel K

Schwinger-Dyson equations

Basic objects are the propagators of the theory.

Quark

$$\langle \bar{\psi}^a \psi^b \rangle \equiv S_F^{ab}(p) = \delta^{ab} \frac{i\not{p} + M(p^2)}{p^2 + M^2(p^2)} Z_f(p^2)$$

Gluon[†]

$$\langle A_\mu^a A_\nu^b \rangle \equiv D_{\mu\nu}(p) = \delta^{ab} \left(\delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) \frac{Z(p^2)}{p^2}$$

Ghost

$$\langle \bar{c}^a c^b \rangle \equiv D_G^{ab}(p) = -\delta^{ab} \frac{G(p^2)}{p^2}$$

Each satisfy a SDE in terms of higher-Green's fns.

([†]in Landau gauge)

Schwinger-Dyson equations

Basic objects are the propagators of the theory.

Quark

$$\begin{array}{c} -1 \\ \text{---} \bullet \text{---} \end{array} = \begin{array}{c} -1 \\ \text{---} \end{array} + \begin{array}{c} -1 \\ \text{---} \bullet \text{---} \circ \text{---} \bullet \text{---} \end{array}$$

The diagram shows the Schwinger-Dyson equation for a quark propagator. On the left is a quark propagator with a self-energy insertion (a grey dot on a solid line). This is equal to the sum of a bare quark propagator (a solid line) and a quark propagator with a ghost loop (a solid line with a loop of ghost lines, represented by a circle with a white vertex and two grey vertices).

Gluon (truncated)

$$\begin{array}{c} -1 \\ \text{---} \bullet \text{---} \end{array} = \begin{array}{c} -1 \\ \text{---} \end{array} - \frac{1}{2} \begin{array}{c} -1 \\ \text{---} \bullet \text{---} \circ \text{---} \bullet \text{---} \end{array} + \begin{array}{c} -1 \\ \text{---} \bullet \text{---} \circ \text{---} \bullet \text{---} \end{array}$$

The diagram shows the Schwinger-Dyson equation for a truncated gluon propagator. On the left is a truncated gluon propagator with a self-energy insertion (a grey dot on a wavy line). This is equal to the sum of a bare truncated gluon propagator (a wavy line), minus half of a truncated gluon propagator with a ghost loop (a wavy line with a loop of ghost lines, represented by a circle with a white vertex and two grey vertices), plus a truncated gluon propagator with a ghost loop (a wavy line with a loop of ghost lines, represented by a circle with a white vertex and two grey vertices).

Ghost

$$\begin{array}{c} -1 \\ \text{---} \bullet \text{---} \end{array} = \begin{array}{c} -1 \\ \text{---} \end{array} + \begin{array}{c} -1 \\ \text{---} \bullet \text{---} \circ \text{---} \bullet \text{---} \end{array}$$

The diagram shows the Schwinger-Dyson equation for a ghost propagator. On the left is a ghost propagator with a self-energy insertion (a grey dot on a dashed line). This is equal to the sum of a bare ghost propagator (a dashed line) and a ghost propagator with a gluon loop (a dashed line with a loop of gluon lines, represented by a circle with a white vertex and two grey vertices).

Each satisfy a SDE in terms of higher-Green's fns.

(†in Landau gauge)

Bethe-Salpeter equations

Rainbow-Ladder truncation

Symmetries help constrain system

Axial-vector WTI

$$P_\mu \Gamma_{5\mu}^a(k; P) = S^{-1}(k_+) \frac{1}{2} \lambda_f^a i \gamma_5 + \frac{1}{2} \lambda_f^a i \gamma_5 S^{-1}(k_-) - M_\zeta i \Gamma_5^a(k; P) - i \Gamma_5^a(k; P) M_\zeta .$$

- Symmetry preserving truncation in DSE and BSE
→ preserve Goldstone character of the pion

BSE

$$\Gamma_{tu}^{(\mu)}(p; P) = \lambda(P^2) \int \frac{d^4 k}{(2\pi)^4} K_{tu;rs}(p, k; P) \left[S(k_+) \Gamma^{(\mu)}(k; P) S(k_-) \right]_{sr}$$

Solve by introducing an eigenvalue $\lambda(P^2)$

Bethe-Salpeter equations

Rainbow-Ladder truncation

Replace quark-gluon vertex by tree-level form.

Quark-gluon vertex

$$\Gamma_{\nu}^{\text{qg}}(p_1, p_2, p_3) = \gamma_{\nu} Z_2 / \tilde{Z}_3 \Gamma^{\text{YM}}(p_3^2)$$

BSE Kernel constructed by considering AVWTI:

Quark scattering kernel

$$K_{tu;sr}(q, p; P) = \frac{g^2 Z(k^2) \Gamma^{\text{YM}}(k^2) Z_{1F}}{k^2} \left(\delta_{\mu\nu} - \frac{k_{\mu} k_{\nu}}{k^2} \right) \left[\frac{\lambda^a}{2} \gamma_{\mu} \right]_{ts} \left[\frac{\lambda^a}{2} \gamma_{\nu} \right]_{ru}$$

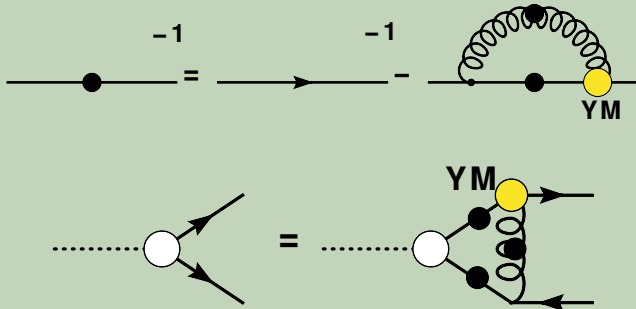
BSE obtained from quark SDE by the substitution:

$$\gamma^{\mu} S(k) \gamma^{\nu} \longrightarrow \gamma^{\mu} S(k_-) \Gamma_M^{(\rho)}(k; P) S(k_+) \gamma^{\nu}$$

Bethe-Salpeter equations

Rainbow-Ladder truncation

Pictorially



- Satisfies AV-WTI
- Reproduces:
 - masses of light pseudoscalar, vectors
 - leptonic decay constants
 - electromagnetic form factors, pion charge radius.

Beyond Rainbow-Ladder



Moving beyond Rainbow-Ladder:

- akin to looking for a pot of gold at the end of a *rainbow*

Technically very challenging:

- Coupled integral equations
- Must preserve symmetries
- Computationally involved:
 - Calculate input Green's functions
 - Solve normalisation conditions.
- Want:
 - Unquenching (quark-loops)
 - Yang-Mills corrections

Beyond Rainbow-Ladder



Be more **humble** and ask for some pi at the end of our rainbow.

- Arises from unquenching (pion-cloud)
- Hadronic contribution (decay widths?)
- additional tensor structure
→ beyond the rainbow

Challenging, but tractable within further simplifying approximations.

Beyond Rainbow-Ladder



Also mandatory to think about additional contributions from Yang-Mills sector.

- Use of *ab initio* quantities to determine:
 - gluon propagator
 - quark-gluon interaction

model determined dynamically

- non vector-vector couplings

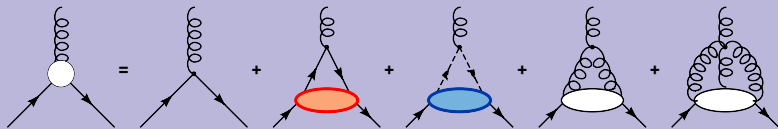
**richer pattern of chiral
symmetry breaking exhibited by
meson masses**

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Quark-Gluon vertex

Quark-Gluon vertex



- **Quark diagram** Hadronic contributions

[C. Fischer, D. Nickel, J. Wambach, PRD **76** (2007) 094009]

[C. Fischer, D. Nickel, RW, arXiv:0807.3486, [hep-ph]]

[C. Fischer, RW, arXiv:0808.3372, [hep-ph]]

- **Ghost diagram** Infrared leading

[R. Alkofer, C. Fischer, F. Llanes-Estrada, MPL A **23**, 1105 (2008)]

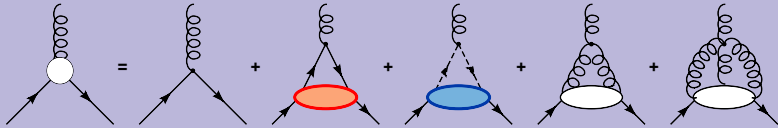
[R. Alkofer, C. Fischer, F. Llanes-Estrada, K. Schwenzer, Annals of Physics, arXiv:0804.3042 [hep-ph]]

For all scales vanishing symmetrically, exhibit power law solutions

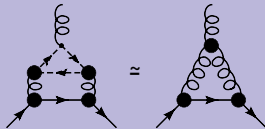
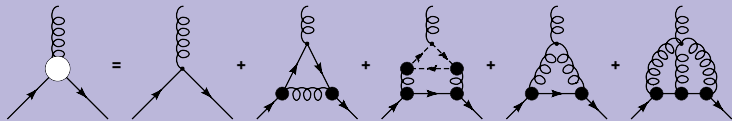
$$\Gamma^{n,m,l} \sim \left(p^2 / \Lambda_{\text{QCD}}^2 \right)^{(n-m)\kappa - l/2}$$

Quark-Gluon vertex

Quark-Gluon vertex

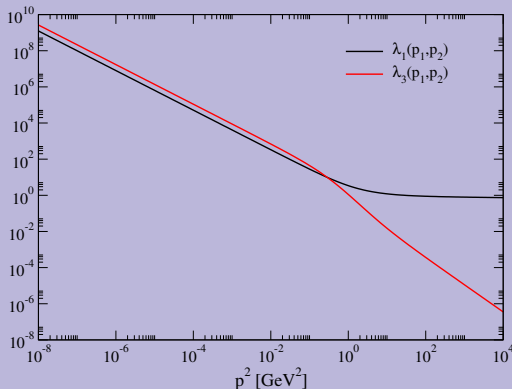


Quark-Gluon vertex – leading skeleton expansion



Quark-Gluon vertex: Numerical solutions

Vector and scalar dressing functions



[R. Alkofer, C. Fischer, F. Llanes-Estrada, K. Schwenzer, Annals of Physics, arXiv:0804.3042 [hep-ph]]

$$\Gamma^\mu(p_1, p_2) = \sum_{k=1}^{12} \lambda_k(p_1, p_2) L^\mu(p_1, p_2)$$

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Hadronic Unquenching Effects (Pion Cloud)

First steps

- 1 Subsume all Yang-Mills corrections into a vertex dressing
 - form inspired by quark-gluon vertex calculations
 - scales left free (but constrained) for parameter fitting
- 2 Gluon propagator obtained from SDE solutions
- 3 Focus on and quantify hadronic effects.

For investigatory purposes we simplify the truncation further:

- Gives idea of necessary difficulty
- Allows techniques to be refined:
 - Solving quark in the complex plane (Euclidean space)
 - Normalisation condition for non-trivial Kernel.

Hadronic Quenching Effects (Pion Cloud)

Modelling the YM part

Yang-Mills part of quark-gluon interaction Γ_μ :

- $Z(k^2)\Gamma_{\text{YM}} \sim \alpha(k^2)$: for large momenta
- $\Gamma_{\text{YM}} \sim (k^2)^{-1/2-\kappa}$: IR soft-singularity in gluon momentum.

[R. Alkofer, C. Fischer, F. Llanes-Estrada, K. Schwenzer, Annals of Physics, arXiv:0804.3042 [hep-ph]]

- For consistency with axWTI, use $\Gamma_\mu \sim \Gamma_{\text{YM}}\gamma_\mu$

Soft-Divergence

$$\Gamma_{\text{YM}}(k^2) = \left(\frac{k^2}{k^2 + d_2} \right)^{-1/2-\kappa} \times \left(\frac{d_1}{d_2 + k^2} + \frac{k^2 d_3}{d_2^2 + (k^2 - d_2)^2} + \frac{k^2}{\Lambda_{\text{QCD}}^2 + k^2} \right) \times \left[\frac{4\pi}{\beta_0 \alpha_\mu} \left(\frac{1}{\log\left(\frac{k^2}{\Lambda_{\text{QCD}}^2}\right)} - \frac{\Lambda_{\text{QCD}}^2}{k^2 - \Lambda_{\text{QCD}}^2} \right) \right]^{-2\delta}$$

Hadronic Quenching Effects (Pion Cloud)

Modelling the YM part

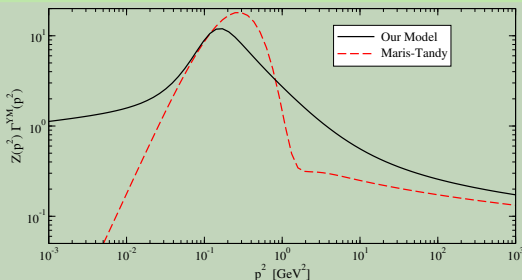
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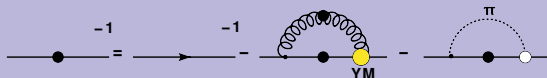
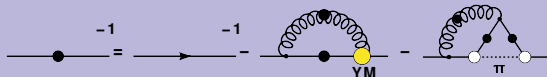
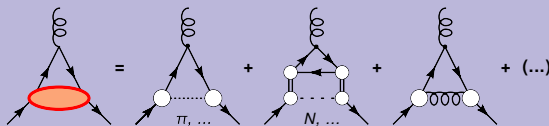
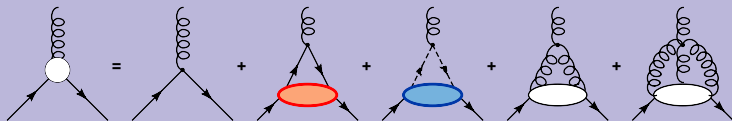
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Soft-Divergence



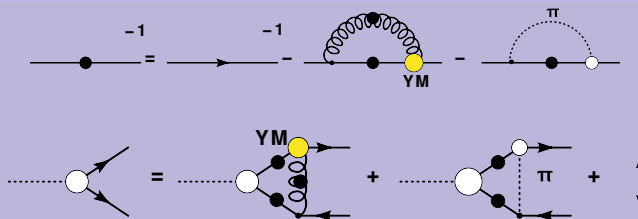
Quark-Gluon vertex - Hadronic unquenching

Truncation



Pion-cloud effects in light mesons

Coupled SDE/BSE system



- AxWTI satisfied in χ -limit:

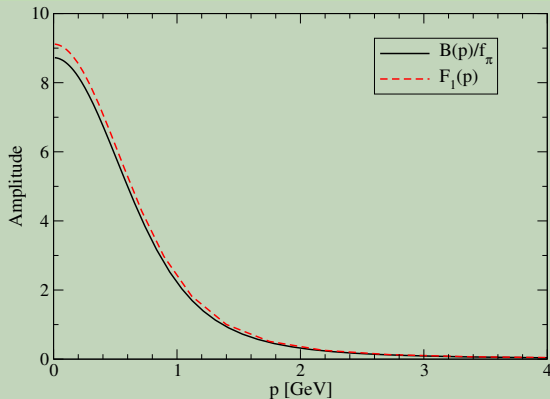
$$P_\mu \Gamma_{5\mu}^a(k; P) = \mathbf{S}^{-1}(k_+) \frac{1}{2} \lambda_f^a i \gamma_5 + \frac{1}{2} \lambda_f^a i \gamma_5 \mathbf{S}^{-1}(k_-) - M_\zeta i \Gamma_5^a(k; P) - i \Gamma_5^a(k; P) M_\zeta .$$

- Generalised GMOR relation well-satisfied:

$$f_\pi m_\pi^2 = r_\pi (m_u(\mu^2) + m_d(\mu^2)) ,$$

Pion-cloud effects in light mesons

Simple off-shell prescription



$$\Gamma_\pi^j(p; P) = \tau^j \gamma_5 \frac{B_\chi(p^2)}{f_\pi}$$

Pion-cloud effects in light mesons

Normalisation

$$\delta^{ij} = 2 \frac{\partial}{\partial P^2} \text{tr} \int \frac{d^4 k}{(2\pi)^4} \left[3 \left(\bar{\Gamma}_\pi^i(k, -Q) S(k + P/2) \Gamma_\pi^j(k, Q) S(k - P/2) \right) + \int \frac{d^4 q}{(2\pi)^4} [\bar{\chi}_\pi^i]_{sr}(q, -Q) K_{tu;rs}^{\text{pion}}(q, k; P) [\chi_\pi^j]_{ut}(k, Q) \right],$$

Canonical condition:

Demand residue of bound-state in inhomogeneous Bethe-Salpeter equation is equal to unity.

Pion-cloud effects in light mesons

Normalisation

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$$\delta^{ij} = \frac{\partial}{\partial P^2} \left[\text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \right]$$

Results

Spectrum of light mesons

	Maris-Tandy		Our Model		Experiment
	w/o pi	incl. pi	w/o pi	incl. pi	
M_π	140	138 [†]	125	138 [†]	138
f_π	104	93 [†] (90)	102	93 [†] (90)	93
M_σ	746	598	638	485	400 – 1200
M_ρ	821	720	795	703	776
f_ρ	160	167 (167)	159	162 (165)	156
M_{a_1}	979	913	941	873	1230
M_{b_1}	820	750	879	806	1230

Results

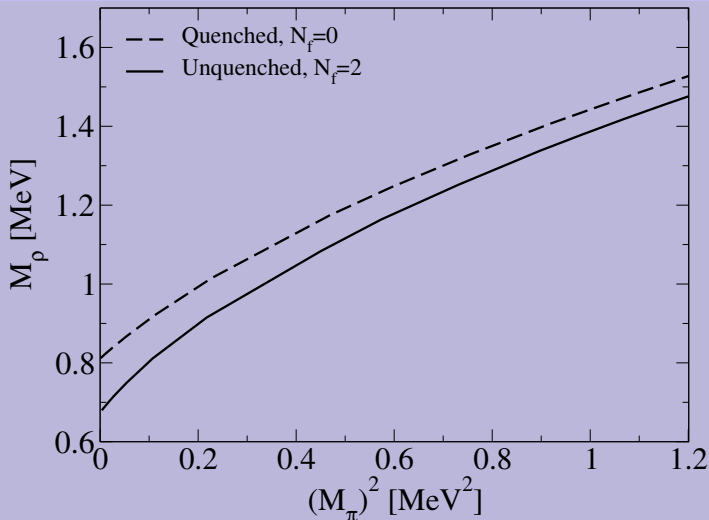
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M_{a_1}	979	913	941	873	1230
M_{b_1}	820	750	879	806	1230
M_η			493	497	548
$M_{\eta'}$			949	963	948

Yang-Mills part of vertex too simple.

Vector vs. Pseudoscalar mass

Mass plots

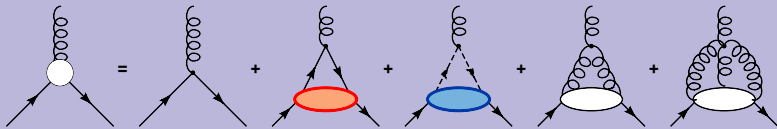


Outline

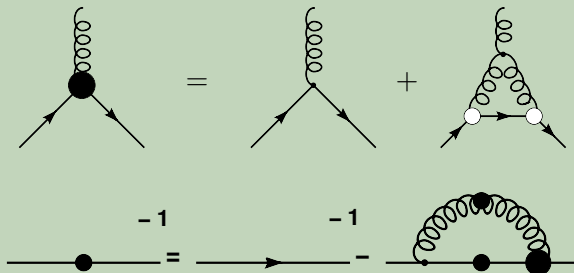
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Gluonic corrections

Quark-Gluon vertex



Consider the approximated system:



Gluonic corrections

Challenges

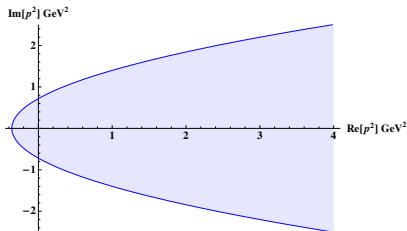
Solving vertex SDE for real Euclidean Momenta

- Nowadays ROUTINE.

all basic QCD vertices tackled to date within some approximation

Bound-states in Euclidean space $\rightarrow P^2 = -M^2$.

- Need both quark propagator and Quark-Gluon vertex for \mathbb{C} -momenta
- Only a (very) technical difficulty – surmounted.



Gluonic corrections

Bethe-Salpeter equation

$$\text{---} \bullet \text{---}^{-1} = \text{---} \text{---}^{-1} + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

Cutting propagators/substitution rules give BSE Kernel:

$$\text{---} \circ \text{---} = \text{---} \circ \text{---} + \text{---} \circ \text{---} + \text{---} \circ \text{---} + \text{---} \circ \text{---}$$

- We *really* calculate these two-loop diagrams
- No Munczek-Nemirovsky
- Axial-Vector WTI preserving truncation.

Gluonic corrections

Simple exploratory model

- Replace three-gluon vertex with tree-level form

$$\Gamma_{\mu\nu\rho}^{(0)abc}(p, k, q) = g f^{abc} \left(\delta_{\mu\nu} (p - q)_\rho + \delta_{\nu\rho} (q - k)_\mu + \delta_{\rho\mu} (k - p)_\nu \right)$$

- Replace internal quark-gluon vertices with γ_μ .
- Replace gluon with some integrated strength

$$Z(p^2) = \frac{g^2}{4\pi} \frac{\pi D}{\omega^2} p^4 \exp\left(-p^2/\omega^2\right)$$

NOT representative of what we expect in nature:

- think of effective gluon interaction compensating for lack of internally dressed vertices.
- no ultraviolet support - renormalisation trivial.

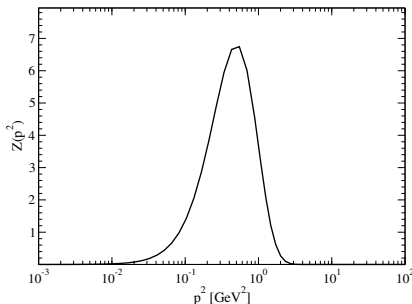
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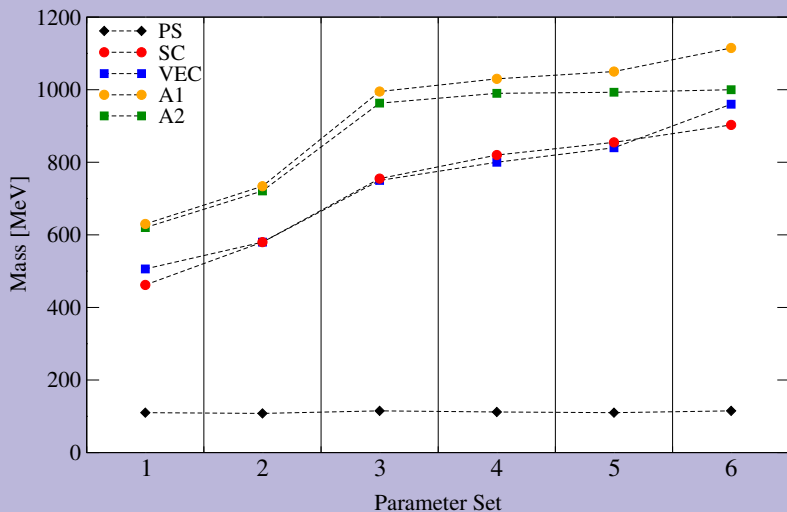
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Gluonic corrections

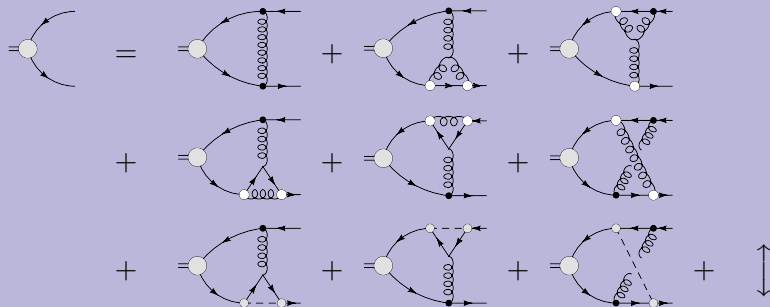
Simple exploratory model – Results



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Outlook



- Unquenching: Improve pion off-shell prescription
- Inputs: Employ solutions from SDE solutions
- Results: Meson spectrum, EM form factors

Seven down - two to go. But we must still solve ...

Outlook

Normalisation

$$\delta^{ij} = \frac{\partial}{\partial P^2} \left[\begin{array}{c} \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \\ + \text{Diagram 4} + \text{Diagram 5} + \text{Diagram 6} \\ + \text{Diagram 7} + \text{Diagram 8} + \text{Diagram 9} \\ + \text{Diagram 10} \end{array} \right]$$

The diagrams represent various Feynman diagrams for the self-energy of a fermion loop. They include:

- Diagram 1: A simple fermion loop with two external lines.
- Diagram 2: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the scalar line.
- Diagram 3: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line.
- Diagram 4: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 5: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 6: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 7: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 8: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 9: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.
- Diagram 10: A fermion loop with a scalar loop (dashed line) and a fermion loop (solid line) attached to the fermion line, with a different internal structure.

Need in order to determine leptonic decay constants

Conclusions

Summary

Quark-Gluon vertex **critical** object. Contains

- Hadronic unquenching effects
- Yang-Mills corrections

Demonstrated that effects from the pion-cloud:

- are generally attractive
- generate effects of right size
- can be successfully modelled in a simple model

Presented progress on state-of-the-art calculations:

- incorporation of leading corrections to vertex
- full two-loop calculations - no kinematic restrictions.