

The Phase Structure of Dense QCD from Chiral Models

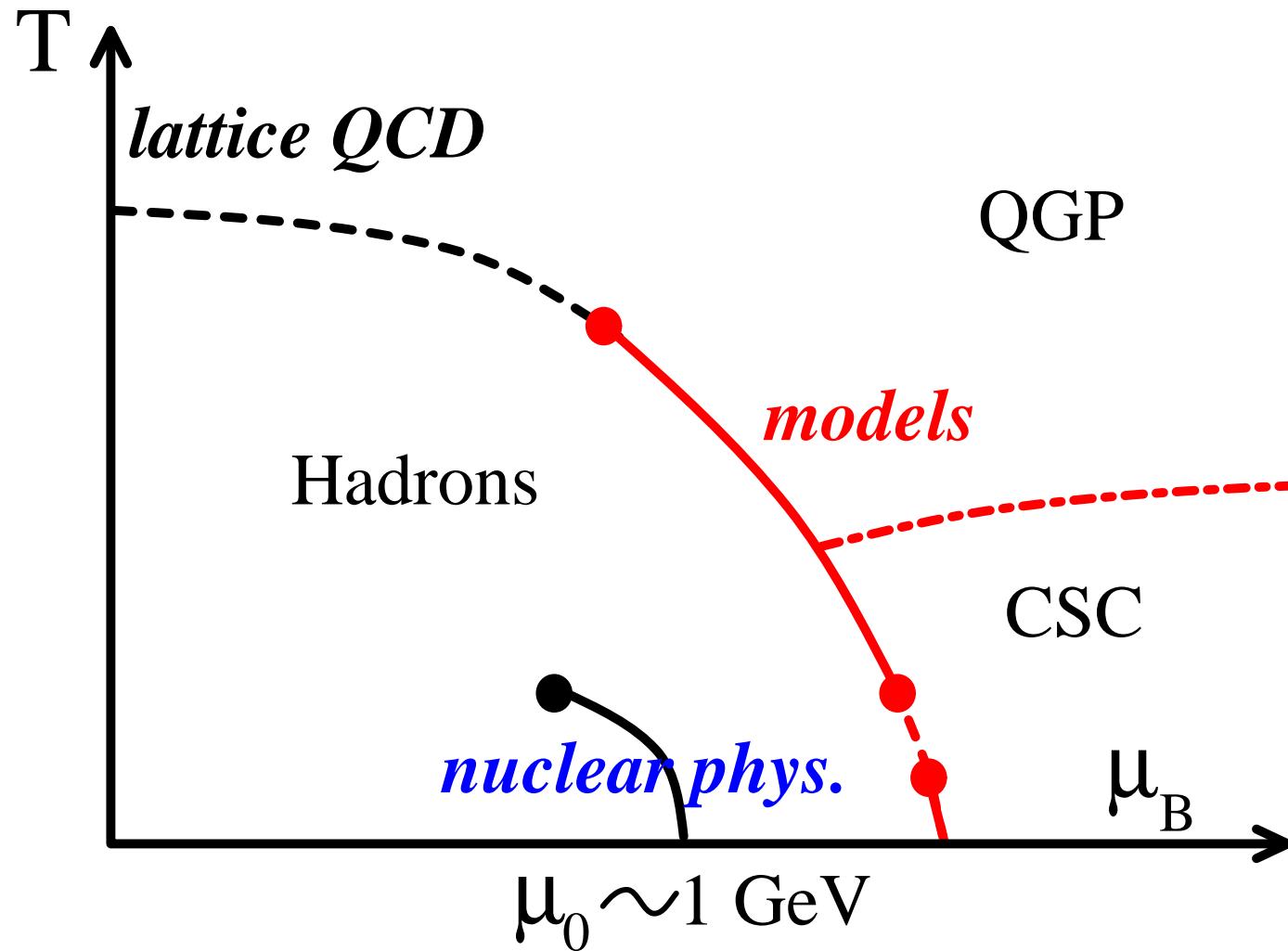
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references

- L. McLerran, K. Redlich and CS, Nucl. Phys. A **824**, 86 (2009);
- M. Harada, CS and S. Takemoto, Phys. Rev. D **81**, 016009 (2010);
- A. Andronic *et al.*, arXiv:0911.4806 [hep-ph].

Phases of QCD at finite temperature and density

- what we know and what we don't know



order of phase transition? chiral and deconfinement? critical point(s)?

- **strong interaction: spontaneous χ SB and color confinement**

- 2 order parameters: $\underbrace{\text{chiral condensate } \langle \bar{q}q \rangle}_{\text{quark dynamics}}$ & $\underbrace{\text{Polyakov loop } \langle \Phi \rangle}_{\text{gluon dynamics}}$

- * left-right mixing: $\langle \bar{q}q \rangle = \langle \bar{q}_L q_R + \bar{q}_R q_L \rangle$
- * free energy of a static quark: $\langle \Phi \rangle = e^{-F_q/T}$

- NJL model with Polyakov loops (PNJL model) [Fukushima (03), Ratti et al. (05)]

- * quark (NJL) and gluon (potential $\mathcal{U}(\Phi)$) minimally coupled via covariant derivative: $\mathcal{L} = \mathcal{L}_{\text{NJL}}(\psi, \Phi[A_0]) + \mathcal{U}(\Phi)$

$$\mathcal{L}_{\text{kin}} = \bar{\psi} i \not{D} \psi, \quad D^\mu = \partial^\mu - i A^\mu, \quad A^\mu = \delta_{\mu 0} A^0$$

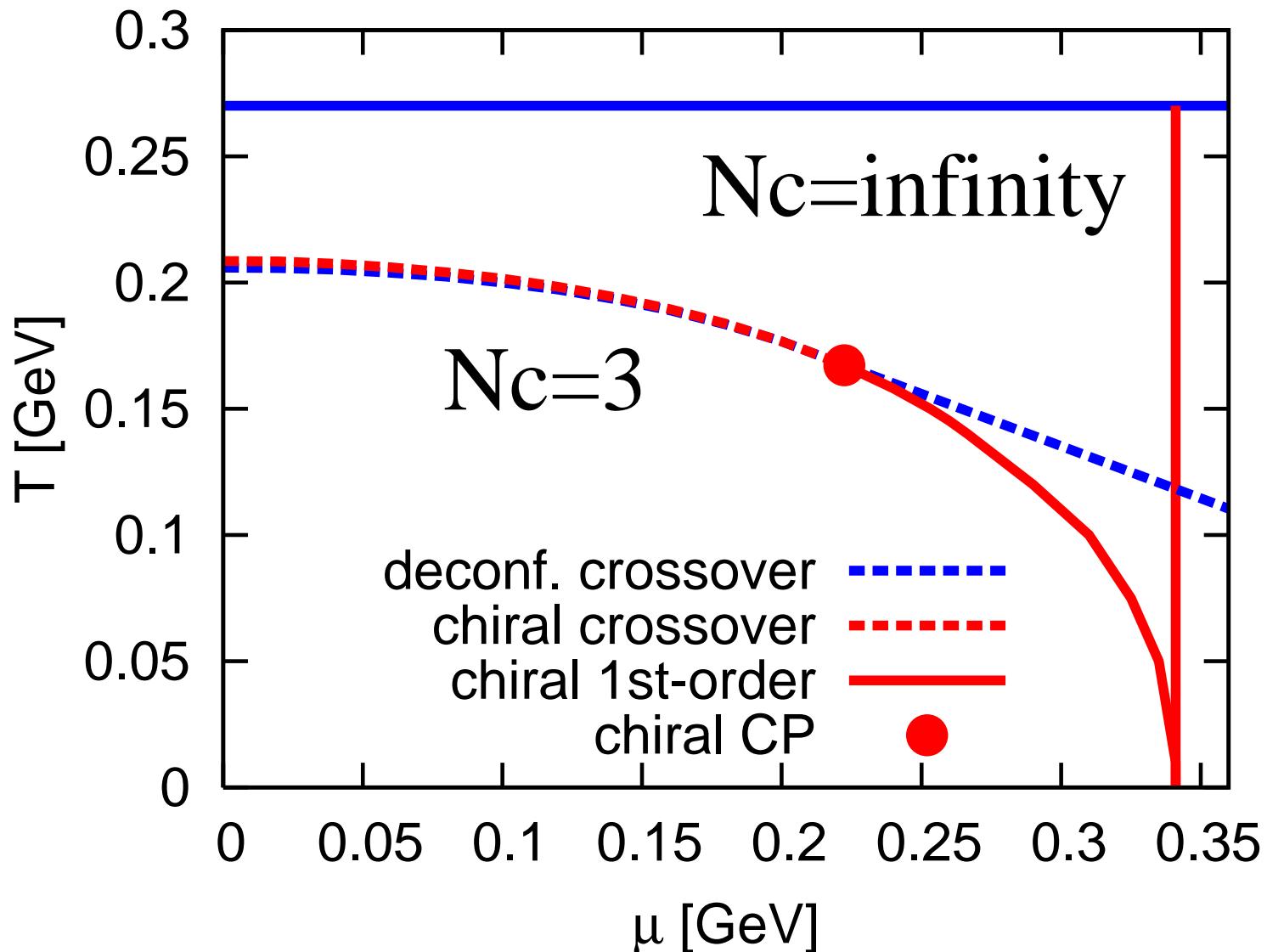
- * model mimics confinement at low temperature $\langle \Phi \rangle \sim 0$

$$\Omega \sim \langle \Phi \rangle \cdot [1\text{-quark states} + 2\text{-quark states}] + 3\text{-quark}$$

\Rightarrow 1- and 2-quark states are thermodynamically irrelevant.

- chiral & deconf. trs for $N_c = 3$ & ∞ from a PNJL

[McLerran-Redlich-CS (08)]



- “transition” from hadronic to quarkyonic phase for $N_c = 3$?

quarkyonic phase $\stackrel{??}{\sim}$ chirally restored & confined phase

- large N_c limit: clear distinction of 2 different confined-phases
baryon number density $\langle N_B \rangle = 0$ (mesonic) and $\neq 0$ (quarkyonic)
- finite N_c : no clear definition
but it may separate meson dominant from baryon dominant region
 \Rightarrow quarkyonic “transition” as meson-baryon “transition”

[A. Andronic *et al.* (09)]

- enhancement of baryon number susceptibility in chiral models

baryons near chiral symmetry restoration?

- standard LSM (naive): $D\chi$ SB generates masses $m_N \xrightarrow{\sigma \rightarrow 0} 0$
- parity doublet model (mirror): $D\chi$ SB generates mass difference
 $m_{N_+} \xrightarrow{\sigma \rightarrow 0} m_{N_-} = m_0 \neq 0$ [Detar-Kunihiro (89)]

anomaly matching? w/o Lorentz invariance

- no WZW term, massless excitations...

Dense baryonic matter in chiral models

- **nuclear matter: known properties**

- binding energy: $E/A(\rho_0) - m_N = -16 \text{ MeV}$
- saturation density: $\rho_0 = 0.16 \text{ fm}^{-3}$
- incompressibility: $K = 9\rho_0^2 \partial^2(E/A)/\partial\rho^2|_{\rho=\rho_0} = 200\text{-}400 \text{ MeV}$

- **standard LSM vs. parity doublet model**

- LSM: no stable ground state corr. to nuclear matter saturation
 - [Kerman-Miller (74)]
 - no direct interaction among $N_+ - N_- - \pi$ at tree
 - cf. vacuum decay width $\Gamma(N_-(1535) \rightarrow N_+\pi)^{(\text{exp})} = 70 \text{ MeV}$
 - PDM: possible to describe saturation [Zschiesche-Tolos-Schaffner-Bielich-Pisarski (07)]

- cold nuclear matter in **SU(2) parity model** [Zschiesche et al. (07)]

- 2 nucleon fields

$$\begin{aligned}\psi_{1L} &: (1/2, 0) & \psi_{1R} &: (0, 1/2) \\ \psi_{2L} &: (0, 1/2) & \psi_{2R} &: (1/2, 0)\end{aligned}$$

- Lagrangian

$$\begin{aligned}\mathcal{L} = & \bar{\psi}_1 i\cancel{D}\psi_1 + \bar{\psi}_2 i\cancel{D}\psi_2 + m_0 (\bar{\psi}_2 \gamma_5 \psi_1 - \bar{\psi}_1 \gamma_5 \psi_2) \\ & + a \bar{\psi}_1 (\sigma + i\gamma_5 \vec{\tau} \cdot \vec{\pi}) \psi_1 + b \bar{\psi}_2 (\sigma - i\gamma_5 \vec{\tau} \cdot \vec{\pi}) \psi_2 \\ & - g_\omega \bar{\psi}_1 \psi \psi_1 - g_\omega \bar{\psi}_2 \psi \psi_2 + \mathcal{L}_M,\end{aligned}$$

$$\begin{aligned}\mathcal{L}_M = & \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} \partial_\mu \pi \partial^\mu \pi \\ & - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu + (g_4)^4 (\omega_\mu \omega^\mu)^2 \\ & + \frac{1}{2} \bar{\mu}^2 (\sigma^2 + \vec{\pi}^2) - \frac{\lambda}{4} (\sigma^2 + \vec{\pi}^2)^2 + \epsilon \sigma\end{aligned}$$

- masses: $m_\pm = \frac{1}{2} \left[\sqrt{(a+b)^2 \sigma^2 + 4m_0^2} \mp (a-b)\sigma \right]$

- parameters ($g_4 = 0$): $m_0 = 790 \text{ MeV}$, $m_\sigma = 371 \text{ MeV}$, $g_\omega = 6.79$

- **thermodynamics of SU(2) parity model** [CS (2010)]

- mean-field approx.: $\langle \sigma \rangle$ and $\langle \omega \rangle$ from $\frac{\partial \Omega}{\partial \sigma} = \frac{\partial \Omega}{\partial \omega} = 0$
- baryon and meson number densities:

$$\rho_B(T, \mu_B) = \sum_{i=\pm} d_i \int \frac{d^3 p}{(2\pi)^3} f(T, \mu; m_i),$$

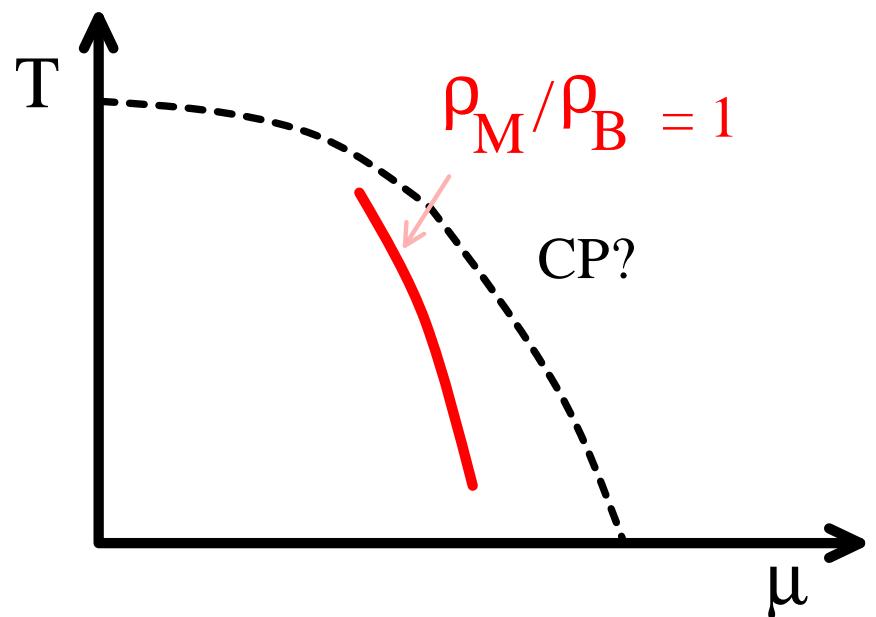
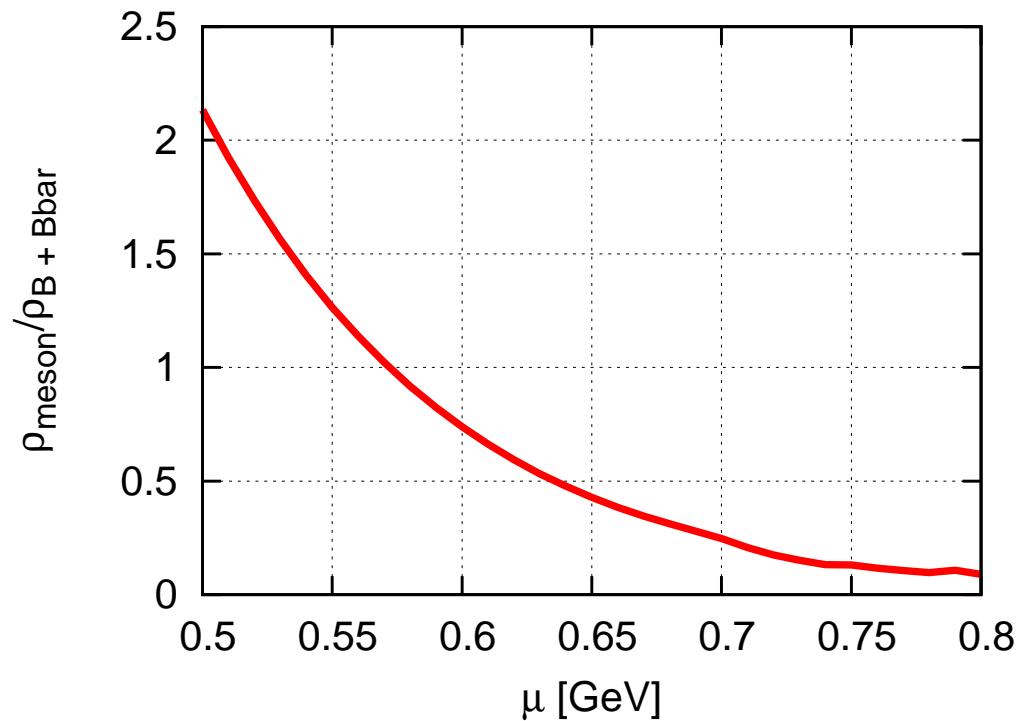
$$\rho_{\bar{B}}(T, \mu_B) = \sum_{i=\pm} d_i \int \frac{d^3 p}{(2\pi)^3} \bar{f}(T, \mu; m_i),$$

$$\rho_M(T) = \sum_{i=\sigma, \pi, \omega} d_i \int \frac{d^3 p}{(2\pi)^3} b(T; m_i)$$

- in-medium meson masses

$$m_\sigma^2 = \frac{\partial^2 \Omega}{\partial \sigma^2}, \quad m_\pi^2 = \frac{\partial^2 \Omega}{\partial \pi^2}$$

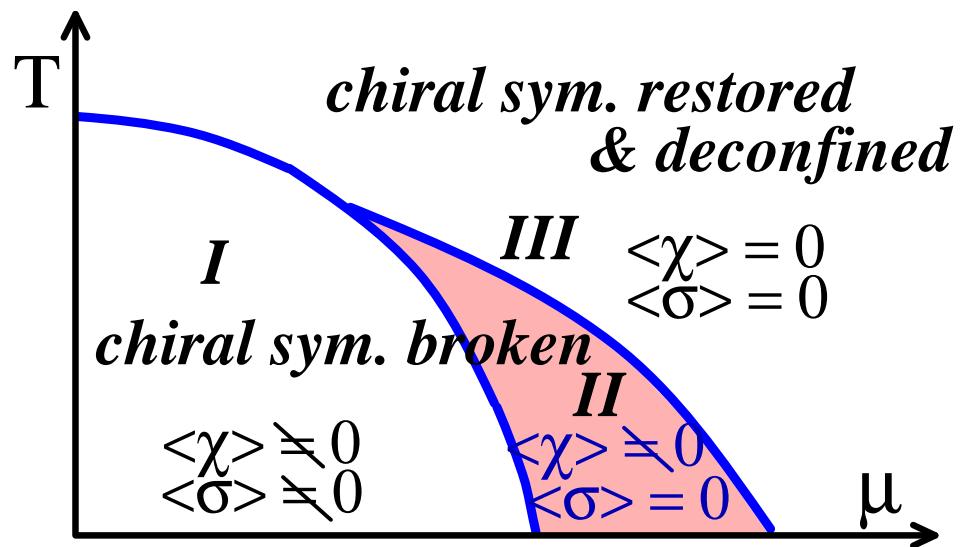
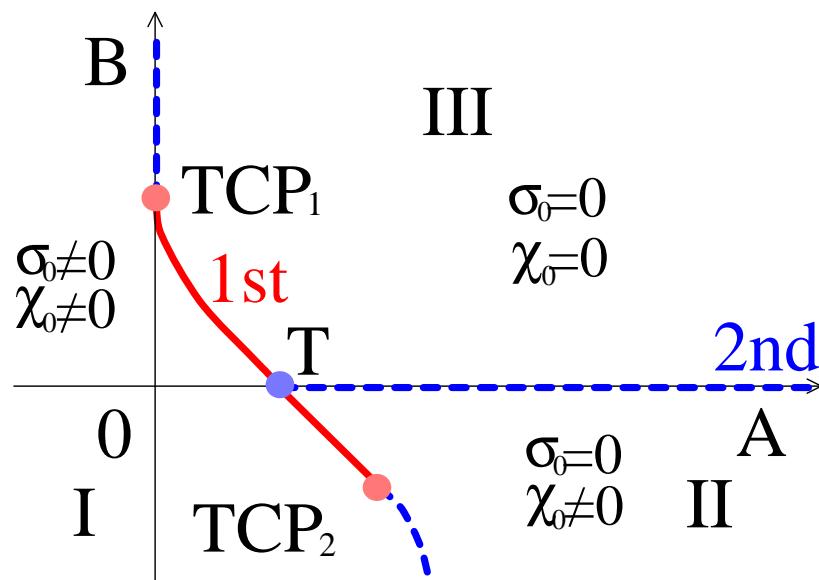
– EoS at finite temperature ($T = 100$ MeV)



- * chiral crossover $\mu_\chi \sim 0.7$ GeV
- * $\rho_{\text{meson}}/\rho_{\text{baryon}} \sim 1$ or $s_{\text{meson}}/s \sim 1$ at $\mu \sim 0.55$ GeV
- * from meson dominant to baryon dominant

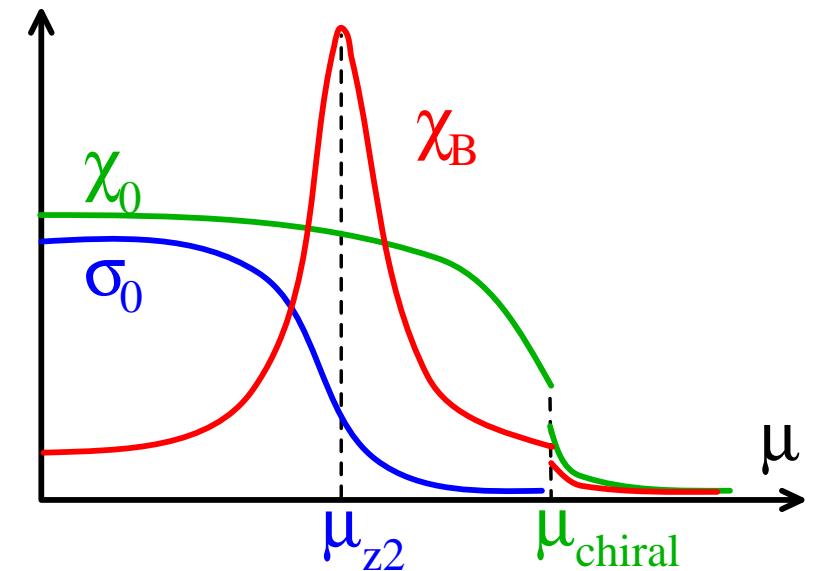
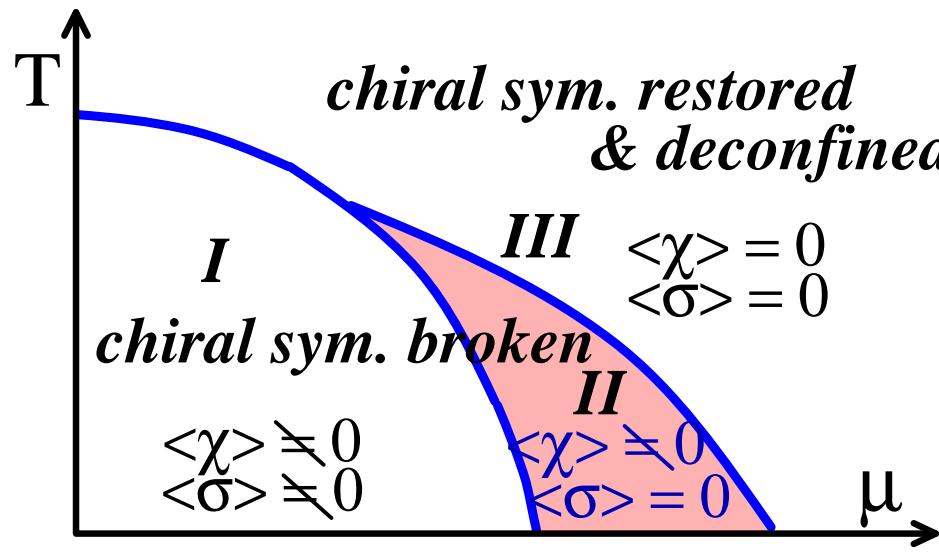
- 2 phases with broken symmetry: distinguished by n_B

[Harada-CS-Takemoto (09)]



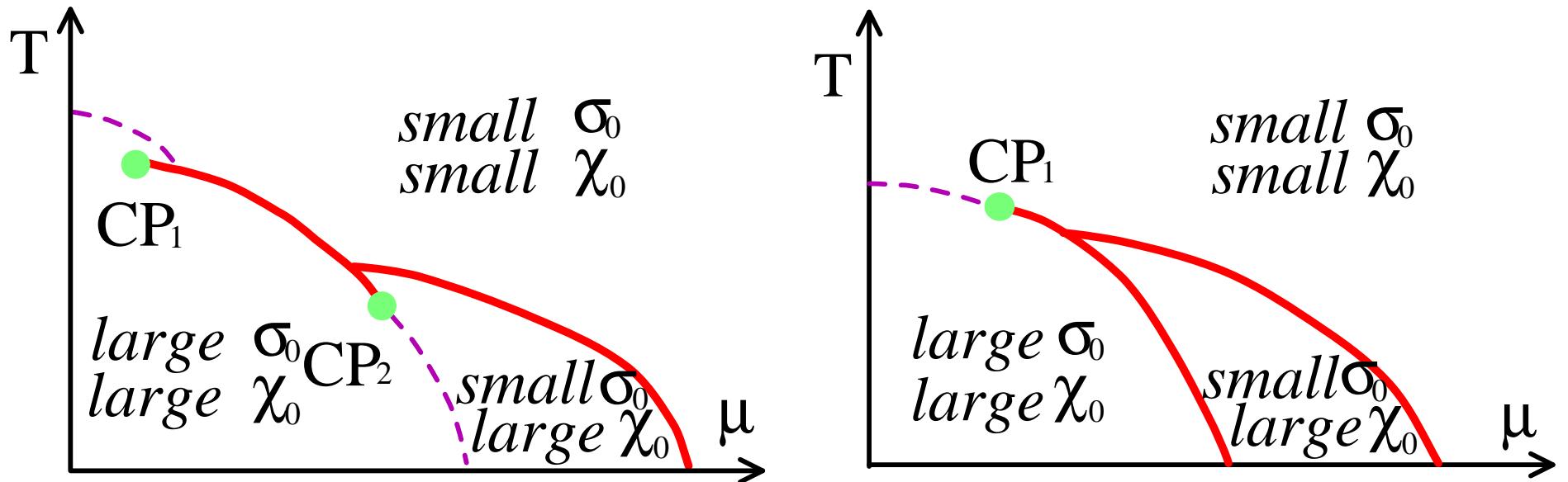
- 2 phases with broken symmetry: distinguished by n_B

[Harada-CS-Takemoto (09)]



- symmetry breaking: $SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V \times Z_{N_f} \rightarrow SU(N_f)_V$
- order parameters: 2-quark state $\sigma \sim \bar{q}q$ and 4-quark state $\chi \sim (\bar{q}q)^2 + \bar{q}\bar{q}-qq$
- 3 phases from a Ginzburg-Landau potential
- I-II: χ_B max. ($\sigma \rightarrow 0$)
- II-III: χ_B no much change (no Yukawa term $\bar{N}N\chi$ in phase II)
- χ_B max. along Z_2 restoration line
can be interpreted as “quarkyonic transition”: baryons more activated

- hypothetical phase diagram in T - μ plane (w/ explicit breaking)



- 2 order parameters: σ (2-quark) and χ (4-quark)
 \Rightarrow 2 phase transitions: restoration of Z_2 center and chiral symmetries
- multiple critical points:
 CP_1 and CP_2 belong to the same universality class
 \Leftrightarrow different universality from anomaly induced CP [Hatsuda et al. (06-07)]
 $\therefore U(1)_B$ is broken in CFL phase.

- hadron mass spectra

phase I: $\sigma_0 \neq 0, \chi_0 \neq 0$	phase II: $\sigma_0 = 0, \chi_0 \neq 0$
$SU(2)_V$	$SU(2)_V \times (Z_2)_A$
$m_S \neq 0, m_P = 0$ $m_V \neq m_A$	$m_S \neq m_P \neq 0$ $m_V \neq m_A$
$m_{N^+} \neq 0$	(i) naive: $\begin{cases} m_{N^+} = 0 \text{ (ground state)} \\ m_{N'^+} = m_{N'^-} \neq 0 \\ \text{(excited states)} \end{cases}$ (ii) mirror: $\begin{cases} m_{N^+} = m_{N^-} \neq 0 \\ \text{(all states)} \end{cases}$

phase I: $\sigma_0 \neq 0, \chi_0 \neq 0$	phase II: $\sigma_0 = 0, \chi_0 \neq 0$
$SU(3)_V$	$SU(3)_V \times (Z_3)_A$
$m_S \neq 0, m_P = 0$ $m_V \neq m_A$	$m_S = m_P \neq 0$ $m_V \neq m_A$
$m_{N^+} \neq 0$	(i) naive: $m_{N^+} \neq 0$ (ii) mirror: $m_{N^+} = m_{N^-} \neq 0$

$$N_f = 2 + 1: (\text{u,d sector}) \quad m_S \neq m_P \quad (\text{s sector}) \quad m_S \simeq m_P$$

Summary and prospects

- **dense nuclear matter and its modeling**
 - saturation properties \Rightarrow parity doublet model
 - meson-baryon “transition”
 - $SU(N_f)_L \times SU(N_f)_R \times Z_{N_f}$ in dense matter
 \Rightarrow a model for 2- and 4-quark states
 - enhancement of χ_B associated with Z_{N_f} symmetry restoration
“quarkyonic transition”: baryons are more activated
 - 2 domains in χ -broken phase?
- **origin of hadron masses?**
 - trace anomaly and hadron mass generation?
cf. $\langle G_{\mu\nu}G^{\mu\nu}\rangle_{T_\chi} \neq 0$
 - naive vs. mirror? sign of axial-couplings