

# *Renormalization of Conserving Selfconsistent Dyson equations*

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## *Motivation*

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- Thermodynamics of strongly interacting systems
- Conservation laws, detailed balance, thermodynamical consistency
- Finite width effects (resonance, damping, ...)

## *Concepts*

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- The  $\Phi$ -derivable scheme
- Renormalization (example  $\phi^4$ )
- Conclusions and outlook

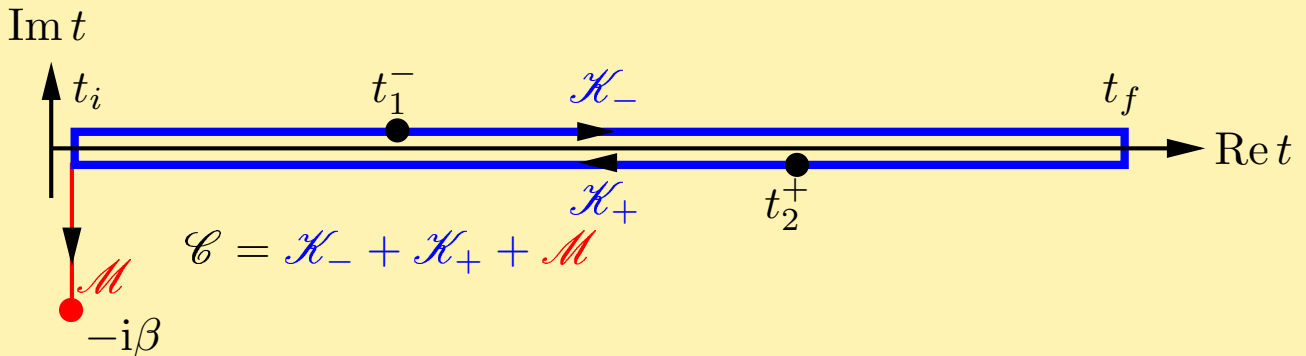
# Schwinger-Keldysh Formalism

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- In equilibrium:  $\rho = \exp(-\beta\mathbf{H})/Z$  with  $Z = \text{Tr} \exp(-\beta\mathbf{H})$

- “Imaginary time” evolution

☞ Imaginary part of the time contour



- Correlation functions with real times:  $iG_{\mathcal{C}}(x_1^-, x_2^+)$
- Fields periodic (bosons) or anti-periodic (fermions)
- Introduce **local** and **bilocal** auxiliary sources:

$$Z[J, K] = N \int D\phi \exp \left[ iS[\phi] + i \{J_1 \phi_1\}_1 + \left\{ \frac{i}{2} K_{12} \phi_1 \phi_2 \right\}_{12} \right]$$

- Generating functional for **connected diagrams**

$$Z[J, K] = \exp(iW[J, K])$$

- The **mean field** and the **connected Green's** function

$$\underbrace{\varphi_1 = \frac{\delta W}{\delta J_1}, G_{12} = -\frac{\delta^2 W}{\delta J_1 \delta J_2}}_{\text{standard quantum field theory}} \Rightarrow \frac{\delta W}{\delta K_{12}} = \frac{1}{2} [\varphi_1 \varphi_2 + iG_{12}]$$

- Legendre transformation for  $\varphi$  and  $G$ :

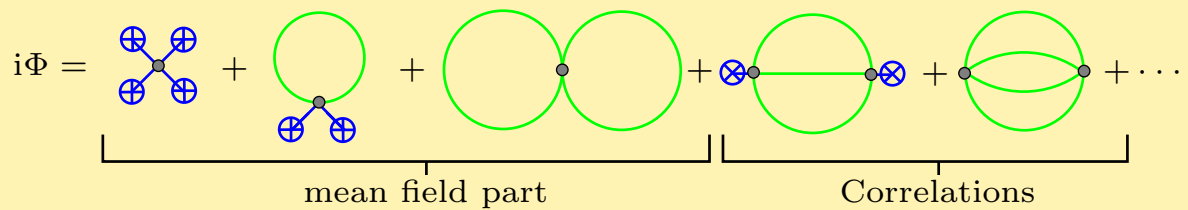
$$\begin{aligned} \Gamma[\varphi, G] &= W[J, K] - \{\varphi_1 J_1\}_1 - \frac{1}{2} \{(\varphi_1 \varphi_2 + iG_{12}) K_{12}\}_{12} \\ &= S_0[\varphi] + \frac{i}{2} \text{Tr} \ln(-iG^{-1}) + \frac{i}{2} \{D_{12}^{-1}(G_{12} - D_{12})\}_{12} \\ &\quad + \Phi[\varphi, G] \Leftarrow \text{all closed 2PI interaction diagrams} \end{aligned}$$

# “Diagrammar”

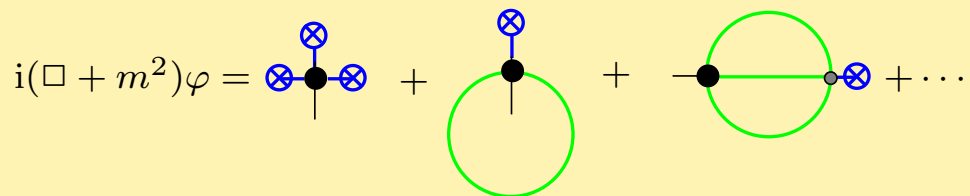
- $\phi^4$ -theory

$$\mathcal{L} = \frac{1}{2}(\partial_\mu\phi)(\partial^\mu\phi) - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

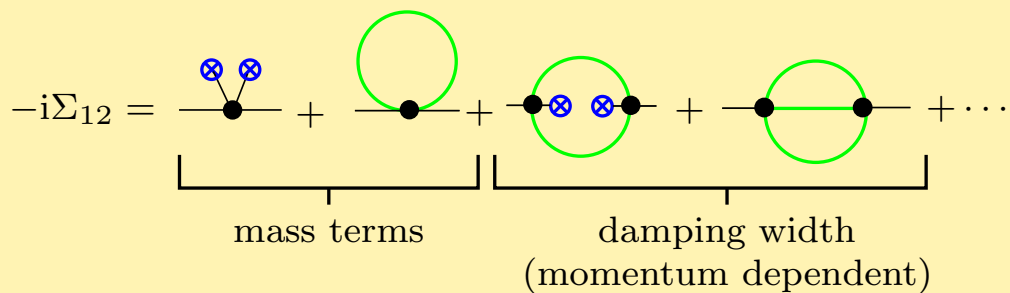
- 2PI Generating Functional



- Mean field equation of motion



- Self-energy



# Properties of the $\Phi$ -derivable Approximations

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## Why using the $\Phi$ -functional?

- Truncation of the Series of diagrams for  $\Phi$
- ☞ Expectation values for currents are conserved  
⇒ “Conserving Approximations”
- In equilibrium  $i\Gamma[\varphi, G] = \ln Z(\beta)$   
(thermodynamical potential)
- consistent treatment of **Dynamical quantities** (real time formalism) and **thermodynamical bulk properties** (imaginary time formalism) like **energy, pressure, entropy**
- Real- and Imaginary-Time quantities “glued” together by **Analytic properties** from (anti-)periodicity conditions of the fields (**KMS-condition**)
- Self-consistent set of equations for self-energies and mean fields

# Problem of Renormalization

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## Why renormalization?

- ☞ Diagrams UV-divergent
- ☞ Control the physical parameters in vacuum
- ☞ Temperature dependence from theory alone

## How to renormalize self-consistent diagrams?

- ☞ In terms of perturbation theory: Resummation of all self-energy insertions in propagators
- ☞ Self-consistent diagrams with explicit nested and overlapping sub-divergences
- ☞ “Hidden” sub-divergences from self-consistency

## How to manage it numerically?

- ☞ Power counting (Weinberg) valid for self-consistent diagrams
- ☞ At finite temperatures:  
Self-consistent scheme rendered finite with local counterterms independent of temperature
- ☞ Analytical properties  $\Rightarrow$  subtracted dispersion relations
- ☞ BPHZ-renormalization  $\Rightarrow$  Subtracting the integrands
- ☞ Advantage: Clear scheme how to subtract temperature independent sub-divergences
- ☞  $\Phi$ -functional  $\Rightarrow$  consistency of counterterms

# Self-consistent Renormalization

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## First step: Vacuum

- Power-counting for **self-consistent propagators** as in perturbation theory:  $\delta = 4 - E$
- Usual **BPHZ-renormalization** for **wave function, mass and coupling constant renormalization**
- In practice: Use Lehmann-representation and dimensional regularization
- ✓ **Closed self-consistent finite** Dyson-equations of motion
- ✓ **Numerically treatable**

## Second step: Finite Temperature

- Split propagator in **vacuum** and **T-dependent** part

$$\text{---} = \text{---} + \text{---}$$

$$iG = iG^{(\text{vac})} + iG^{(\text{T})}$$

- Expand self-energy around vacuum part

$$\text{---} + \text{---} + \text{---}$$

$$-i\Sigma^{(\text{vac})} \quad -i\Sigma^{(0)} \quad -i\Sigma^{(\text{r})}$$

- Need further splitting of propagator

$$\text{---} = \text{---} + \text{---}$$

$$iG^{(\text{T})} = iG^{(\text{vac})} + iG^{(\text{r})}$$

# Self-consistent Renormalization

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## Third step: 4-point vertex renormalization

- $\Sigma^0$  linear in  $G^{(r)}$   $\Rightarrow$

$$\text{---} \bigcirc_{\Gamma^{(4)}} \text{---} = \square_{\Lambda}$$

- Equation of motion  $\Rightarrow$

$$\square_{\Lambda} = \bigcirc_{\Gamma^{(4)}} + \begin{array}{c} \bigcirc_{\Gamma^{(4)}} \\ | \\ \square_{\Lambda} \end{array}$$

☞ s-channel Bethe-Salpeter equation

$$\text{---} \bigcirc_{\Gamma^{(4)}} \text{---} \quad \text{cuts more than three lines!}$$

$\Rightarrow$  “BPHZ Boxes” in ladder-diagrams **do not cut inside**  $\Gamma^{(4)}$ .

$\Rightarrow$  Asymptotics + BPHZ-formalism:

$$\Gamma^{(4)}(l, p) - \Gamma^{(4)}(l, 0) \cong O(l^{-\alpha}) \text{ with } \alpha > 0$$

$\Rightarrow$  Renormalized eq. of motion for  $\Lambda$ :

$$\begin{aligned} \Lambda(p, q) = & \Lambda(0, 0) + \Gamma^{(4)}(p, q) - \Gamma^{(4)}(0, 0) \\ & + i \int \frac{d^4 l}{(2\pi)^4} [\Gamma^{(4)}(p, l) - \Gamma^{(4)}(0, l)] [G^{\text{vac}}]^2(l) \Lambda(l, q) \\ & + i \int \frac{d^4 l}{(2\pi)^4} \Lambda(0, l) [G^{\text{vac}}]^2(l) [\Gamma^{(4)}(l, q) - \Gamma^{(4)}(l, 0)] \end{aligned}$$

✓ Self-energy finite with **vacuum counter terms**

# Example: Tadpole+Sunset

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## Approximation of the $\Phi$ -functional

$$\begin{aligned}
 i\Phi &= \text{[Two circles joined at a point]} + \text{[Circle with two internal arcs]} \\
 -i\Sigma &= \text{[Circle on a line]} + \text{[Circle with a horizontal line through it]} \\
 -i\Gamma^{(4)} &= \text{[Four lines meeting at a point]} + \text{[Oval with four external lines]}
 \end{aligned}$$

## Renormalization (*vacuum*)

$$\begin{aligned}
 -i\Sigma &= \text{[Circle on a line]} + \text{[Circle on a line, dashed box]} \\
 &+ \text{[Circle with horizontal line, dashed box]} + \text{[Circle with horizontal line, dashed box]} \\
 &+ \text{[Circle with horizontal line, dashed box]}
 \end{aligned}$$

+overall

- In practice: Use dispersionrelations for propagators
- ☞ Kernels, can be calculated analytically with standard formulae of dimensional regularization
- ✓ Finite Self-consistent integral equations of motion  $\Rightarrow$  **Solved iteratively**
- Calculate also  $\Gamma^{(4)}$  and  $\Lambda(0, q)$



# Example: Tadpole+Sunset

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## Renormalization (*Finite Temperature*)

$$-i\Sigma^{(T)}(p) =$$

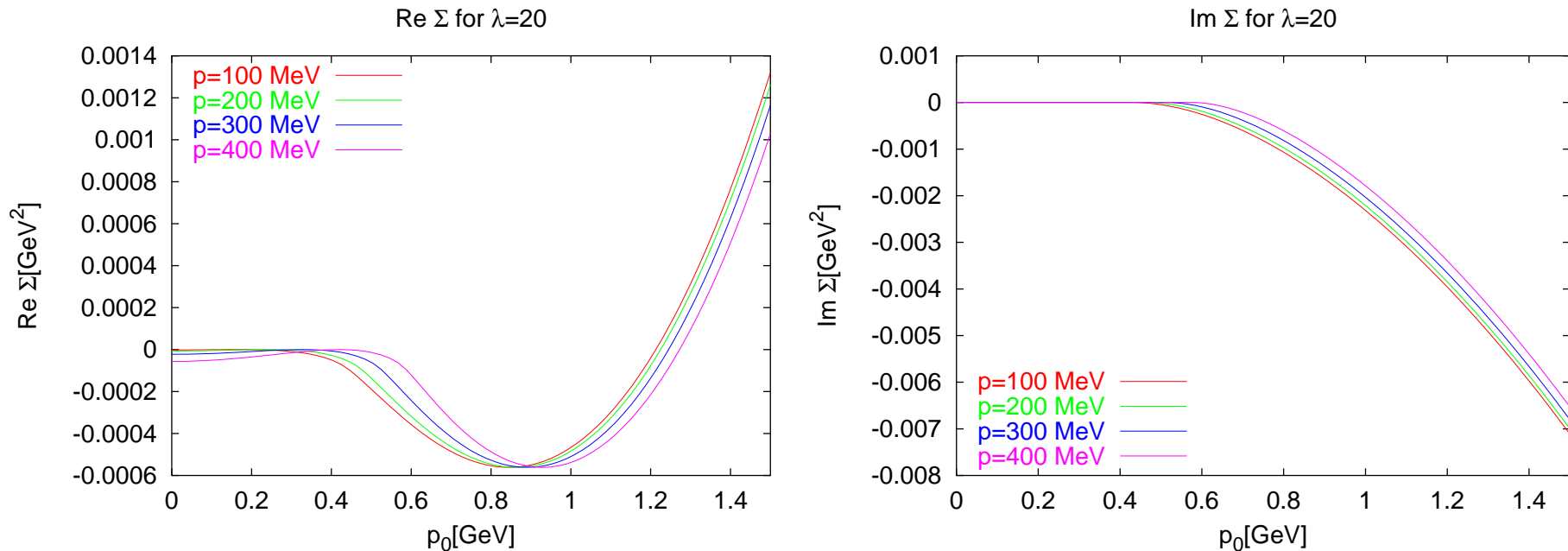
The equation shows the renormalization of the tadpole and sunset diagrams at finite temperature. The left-hand side is  $-i\Sigma^{(T)}(p)$ . The right-hand side consists of five diagrams:

- Diagram 1: A tadpole diagram with a blue loop and a red loop, both attached to a vertex with momentum  $p$ . The external momentum is  $p$ .
- Diagram 2: A sunset diagram with a blue loop and a red loop, both attached to a vertex with momentum  $0$ . The external momentum is  $0$ .
- Diagram 3: A tadpole diagram with a green loop and a blue square vertex labeled  $\Lambda$ , both attached to a vertex with momentum  $0$ . The external momentum is  $0$ .
- Diagram 4: A sunset diagram with a blue loop and a red loop, both attached to a vertex with momentum  $0$ . The external momentum is  $0$ .
- Diagram 5: A sunset diagram with a red loop and a red loop, both attached to a vertex with momentum  $0$ . The external momentum is  $0$ .

- Only finite integrals
- ✓ Numerics for three-dim integrals on a lattice in  $p_0$  and  $|\vec{p}|$
- ✓ Equations of motion solved iteratively

# Results for “Sunset + Tadpole” at $T = 0$

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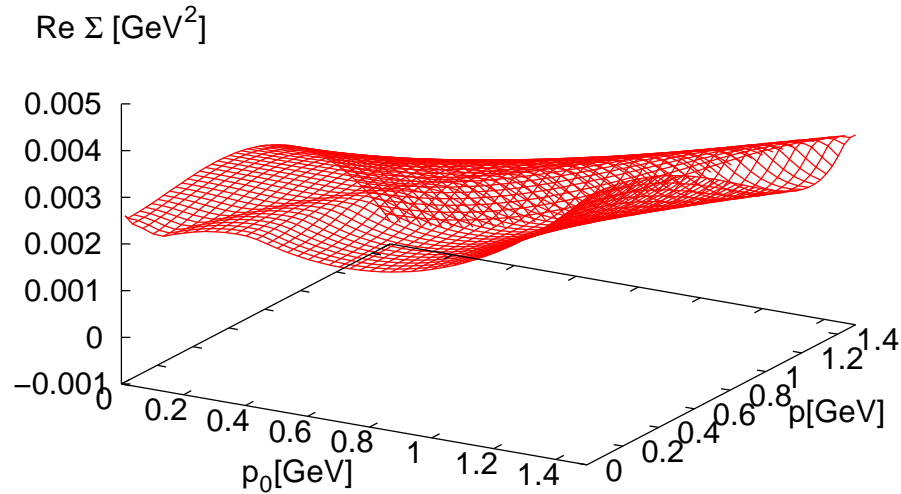


- Difference between perturbative and self-consistent calculation invisible!
- ☞ Tadpole contribution “renormalized away”  $\Rightarrow$  on-shell renormalization scheme
- ☞ Main contribution from the pole term of the propagator
- ☞ Threshold for continues part of the spectral function  $\sqrt{s} = 3m!$

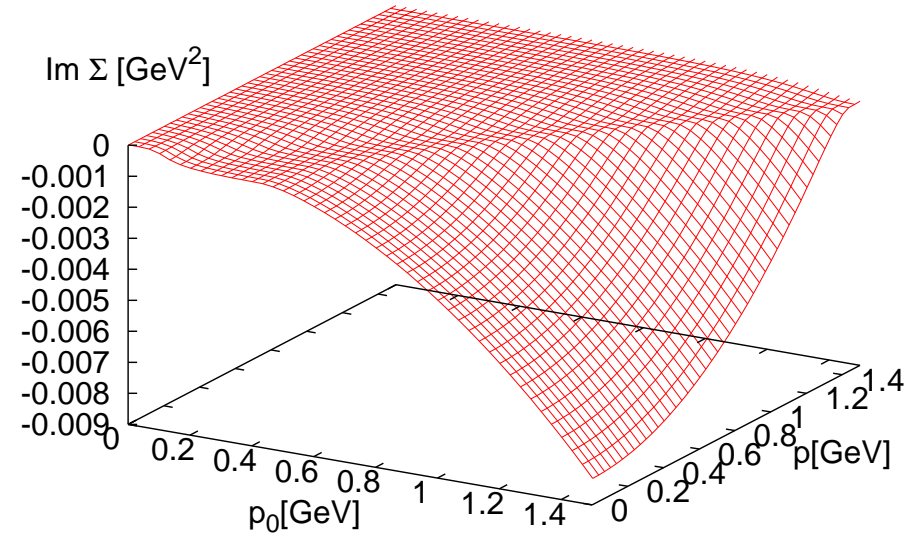
# Results for “Sunset + Tadpole” at $T > 0$

#11

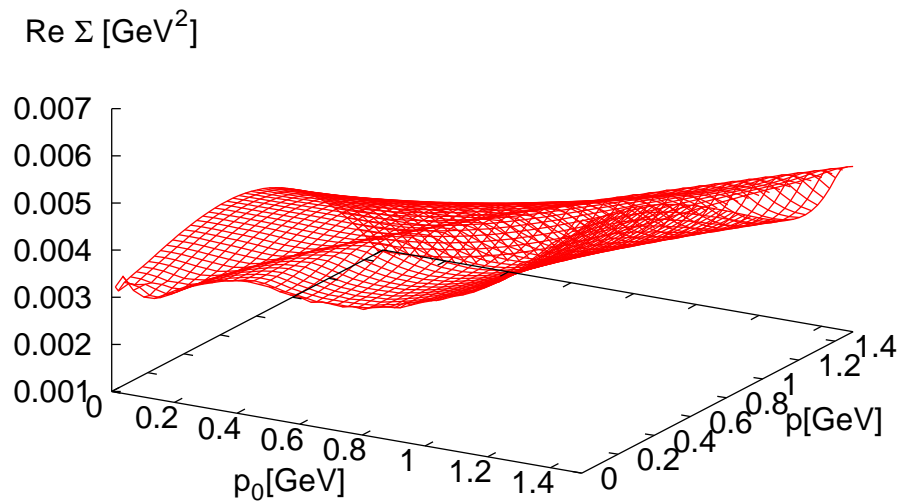
Pert. Re  $\Sigma$  for  $T=100\text{MeV}$ ,  $\lambda=20$



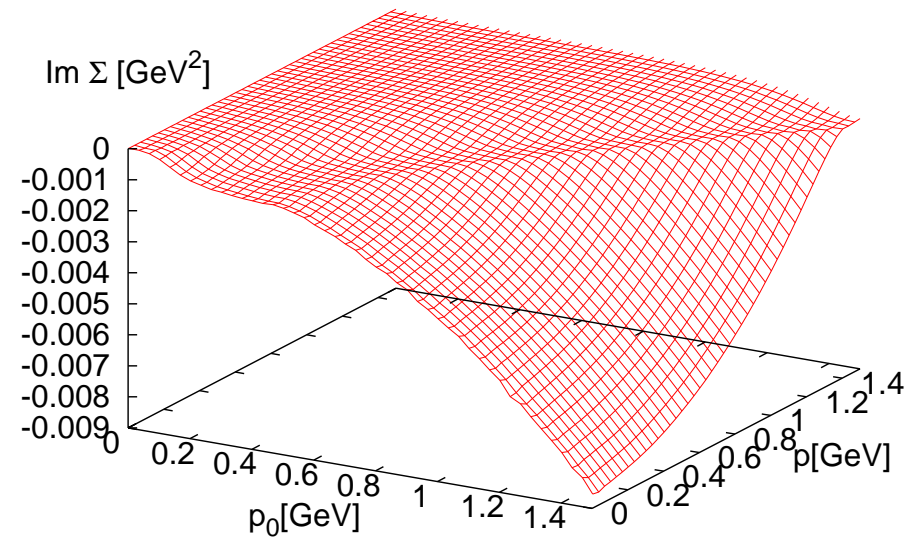
Pert. Im  $\Sigma$  for  $T=100\text{MeV}$ ,  $\lambda=20$



Re  $\Sigma$  for  $T=100\text{MeV}$ ,  $\lambda=20$



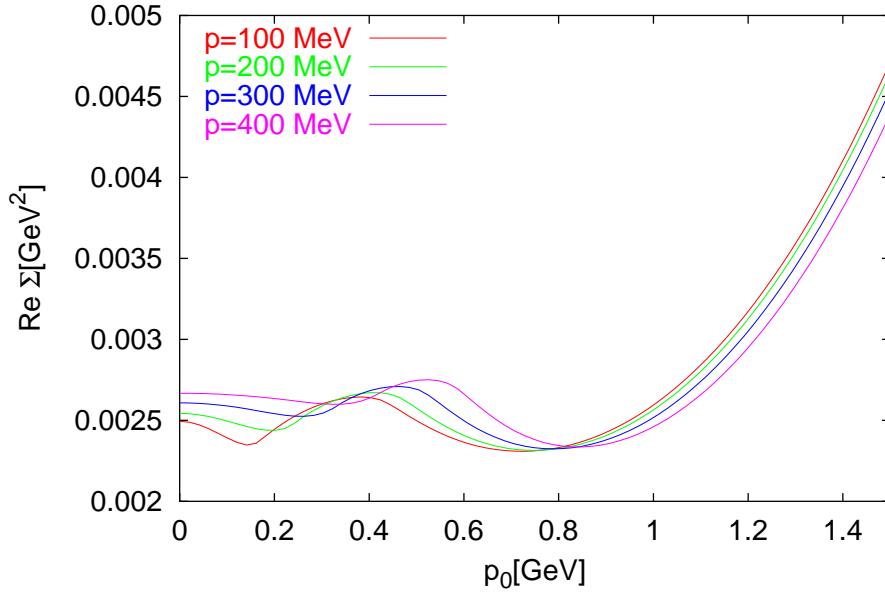
Im  $\Sigma$  for  $T=100\text{MeV}$ ,  $\lambda=20$



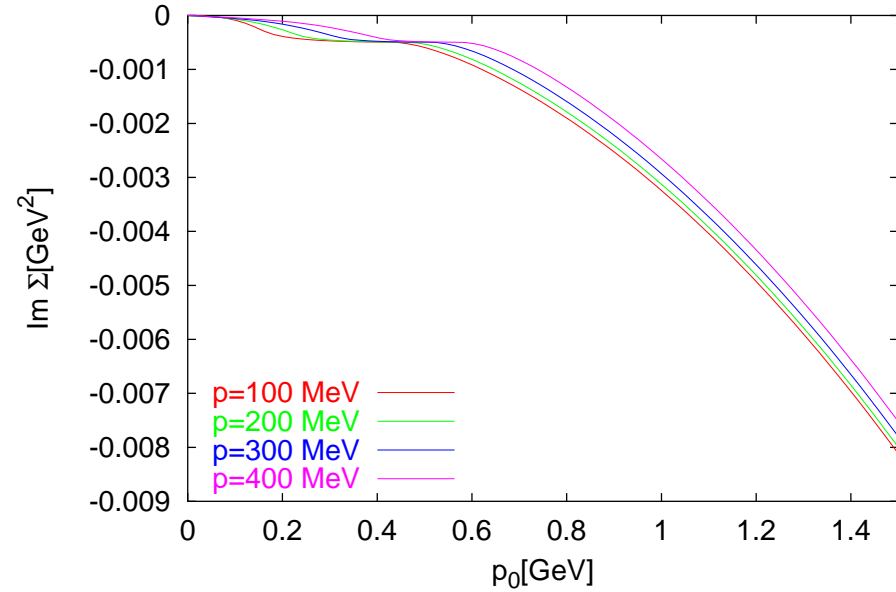
# Results for “Sunset + Tadpole” at $T > 0$

#12

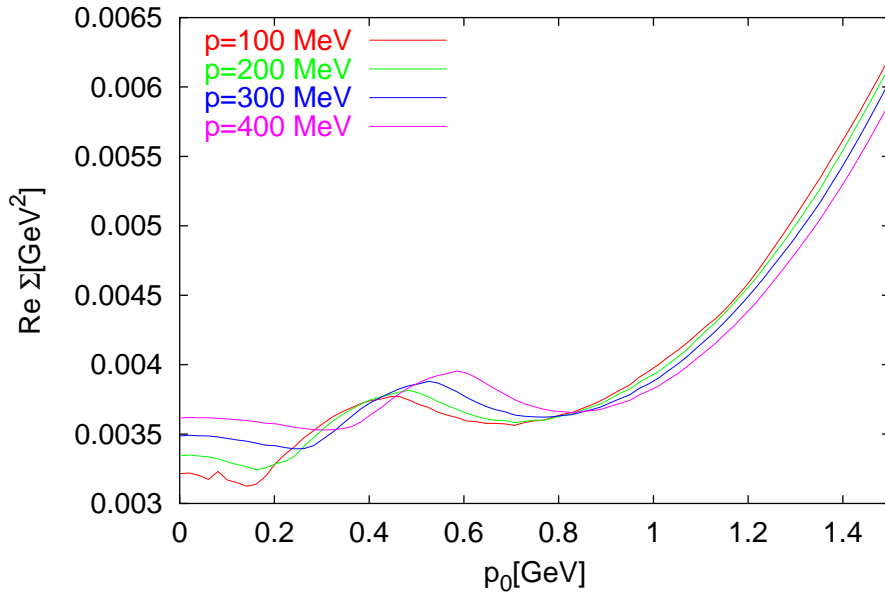
Pert. Re  $\Sigma$  for  $\lambda=20$ ,  $T=100$  MeV



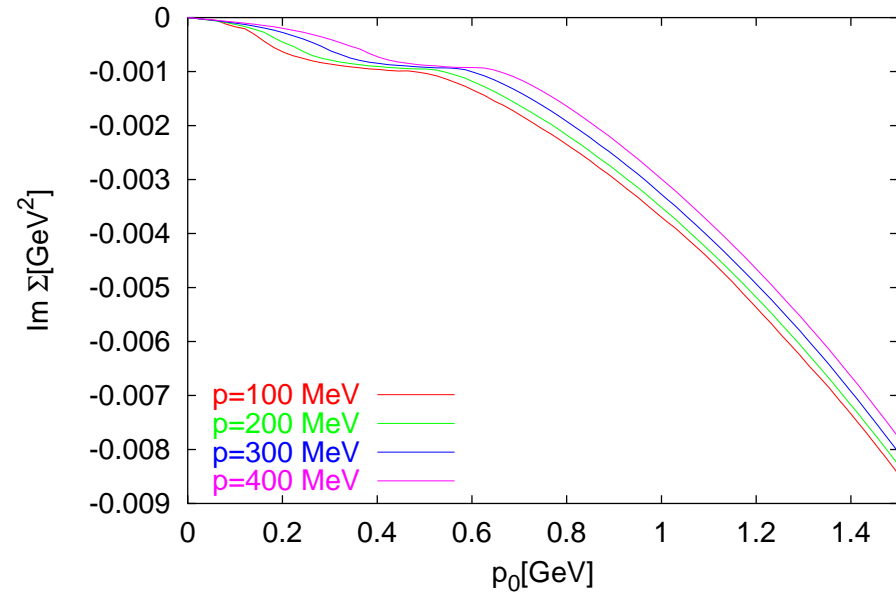
Pert. Im  $\Sigma$  for  $\lambda=20$ ,  $T=100$  MeV



Re  $\Sigma$  for  $\lambda=20$ ,  $T=100$  MeV



Im  $\Sigma$  for  $\lambda=20$ ,  $T=100$  MeV



# Conclusions and Outlook

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## Summary

- ✓ Self-consistent  $\Phi$ -derivable schemes
- ✓ Renormalization:  
<http://arXiv.org/abs/hep-ph/0107200>
- ✓ Numerical treatment

## Outlook

- ✓ “Toolbox” for application to realistic models
  - ✓ Symmetry analysis for  $\Phi$ -derivable approximations:  
PhD-thesis: <http://theory.gsi.de/~vanhees>
  - ✓ Perspectives for self-consistent treatment of vector particles:  
<http://arXiv.org/abs/hep-ph/0002087>
  - ✗ General gauge theories?
  - ✗ QCD e.g. beyond HTL?
  - ✓ Transport equations for particles with finite width
- ☞ See talk given by Jörn Knoll