
HYDJET++:
INTERPLAY BETWEEN
SOFT AND HARD PHYSICS

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OUTLINE

I. HYDJET++ model (hydro + jets)

II. Model results for the ratio
 $v_4/(v_2)^2$ at RHIC and LHC

III. NCQ-scaling at RHIC and LHC

**I. HYDJET++ =
FASTMC + HYDJET**

HYDJET++ event generator

I.Lokhtin, L.Malinina, S.Petrushanko, A.Snigirev, I.Arsene, K.Tywoniuk,
Comp. Phys. Commun.180 (2009) 779-799 (arXiv:0809.2708[hep-ph])

The soft part of HYDJET++ event represents the "thermal" hadronic state.

- ✓ multiplicities are determined assuming thermal equilibrium
- ✓ hadrons are produced on the hypersurface represented by a parameterization of relativistic hydrodynamics with given freeze-out conditions
- ✓ chemical and kinetic freeze-outs are separated
- ✓ decays of hadronic resonances are taken into account (360 particles from SHARE data table) with "home-made" decayer

*the model reproduces soft hadroproduction features at RHIC
(particle spectra, elliptic flow, HBT)*

The hard, multi-partonic part of HYDJET++ event is identical to the hard part of Fortran written HYDJET (PYTHIA6.4xx + PYQUEN1.5) => **now PYTHIA Perugia 2011 tune!!** PYQUEN event generator is used for simulation of rescattering, radiative and collisional energy loss of hard partons in expanding quark-gluon plasma created in ultrarelativistic heavy ion AA collisions. HYDJET++ includes nuclear shadowing correction for parton distributions (important at LHC!) Impact-parameter dependent parameterization of *nuclear shadowing* (K.Tywoniuk, I.Arsene, L.Bravina, A.Kaidalov and E.Zabrodin, Phys. Lett. B 657 (2007) 170)

Model parameters.

1. Thermodynamic parameters at chemical freeze-out: T_{ch} , $\{\mu_B, \mu_S, \mu_Q\}$
2. If thermal freeze-out is considered: T_{th} , $\mu\pi$ -normalisation constant
3. Volume parameters: $T, \Delta T, R$
1. ρ_{max} -maximal transverse flow rapidity for Bjorken-like parametrization
5. η_{max}^u -maximal space-time longitudinal rapidity which determines the rapidity interval $[-\eta_{max}, \eta_{max}]$ in the collision center-of-mass system.
6. Impact parameter range: minimal b_{min} and maximal b_{max} impact parameters
7. Flow anisotropy parameters $\delta(b), \epsilon(b)$

PYTHIA+PYQUEN obligatory parameters

9. Beam and target nuclear atomic weight A
10. $\sqrt{s_{NN}}$ -c.m.s. energy per nucleon pair (PYTHIA initialization at given energy)
11. **ptmin** – minimal pt of parton-parton scattering in PYTHIA event (ckin(3) in /pysubs/)
12. **nhsel** flag to include jet production in hydro-type event:

0 - jet production off (pure FASTMC event),
1 - jet production on, jet quenching off (FASTMC+njet*PYTHIA events),
2 - jet production & jet quenching on (FASTMC+njet*PYQUEN events),
3 - jet production on, jet quenching off, FASTMC off (njet*PYTHIA events),
4 - jet production & jet quenching on, FASTMC off (njet*PYQUEN events);

13. **ishad** flag to switch on/off nuclear shadowing

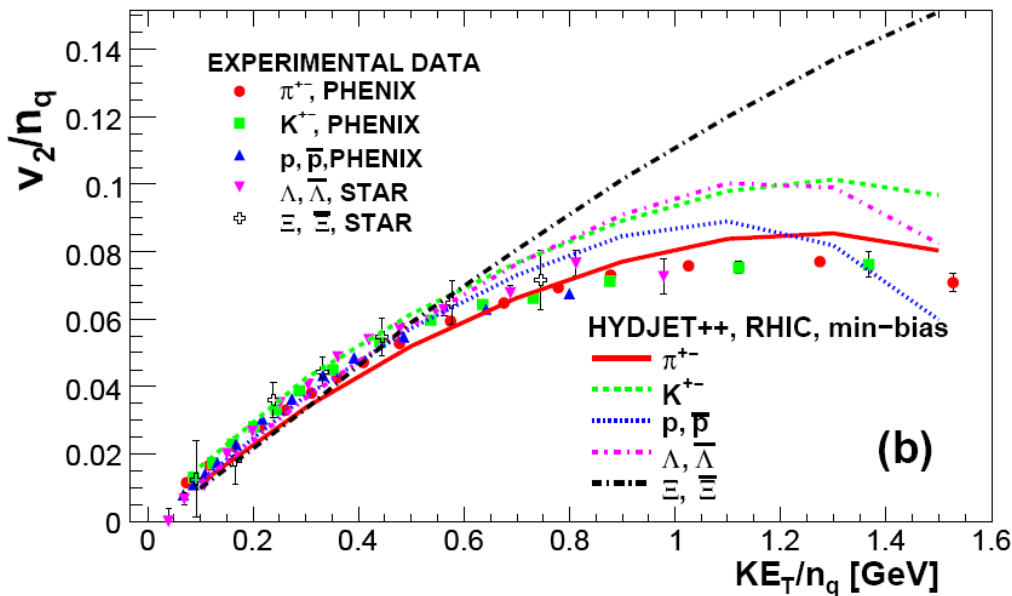
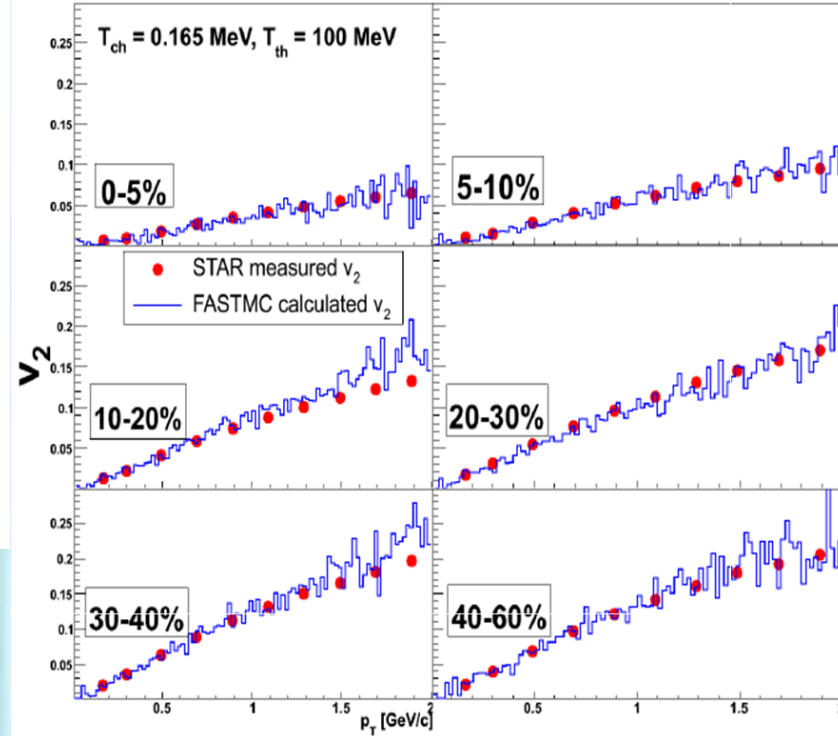
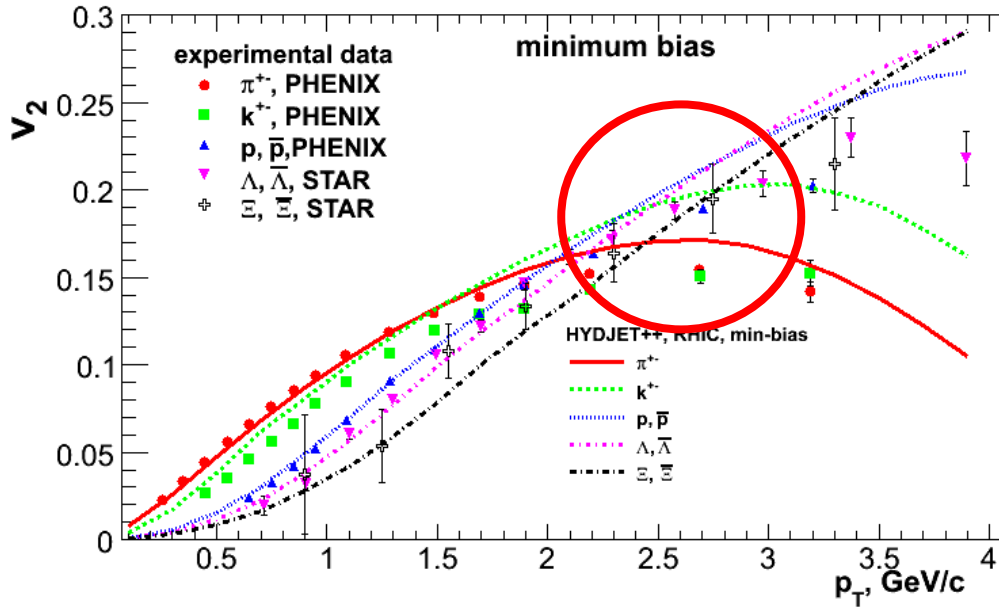
Parameters of energy loss model in PYQUEN

(default, but can be changed from the default values by the user)

1. **T0** - initial temperature of quark-gluon plasma for central Pb+Pb collisions at mid-rapidity (initial temperature for other centralities and atomic numbers will be calculated automatically) at LHC: **T0=1 GeV**, at RHIC(200 AGeV) **T0=0.300 GeV**
2. **tau0** - proper time of quark-gluon plasma formation at LHC: **tau0=0.1 fm/c**, at RHIC(200 AGeV) **tau0=0.4 fm/c**
3. **nf** - number of active quark flavours in quark-gluon plasma (nf=0, 1, 2 or 3) at LHC: **nf=0**, at RHIC(200 AGeV) **nf=2**
4. **ienglu** - flag to fix type of medium-induced partonic energy loss (ienglu=0 - radiative and collisional loss, ienglu=1 - radiative loss only, ienglu=2 - collisional loss only, default value is ienglu=0);
ianglu - flag to fix type of angular distribution of emitted gluons (ianglu=0 - small-angular, ianglu=1 - wide-angular, ianglu=2 - collinear, default value is ianglu=0).
ienglu=0

RHIC DATA VS. HYDJET++ MODEL

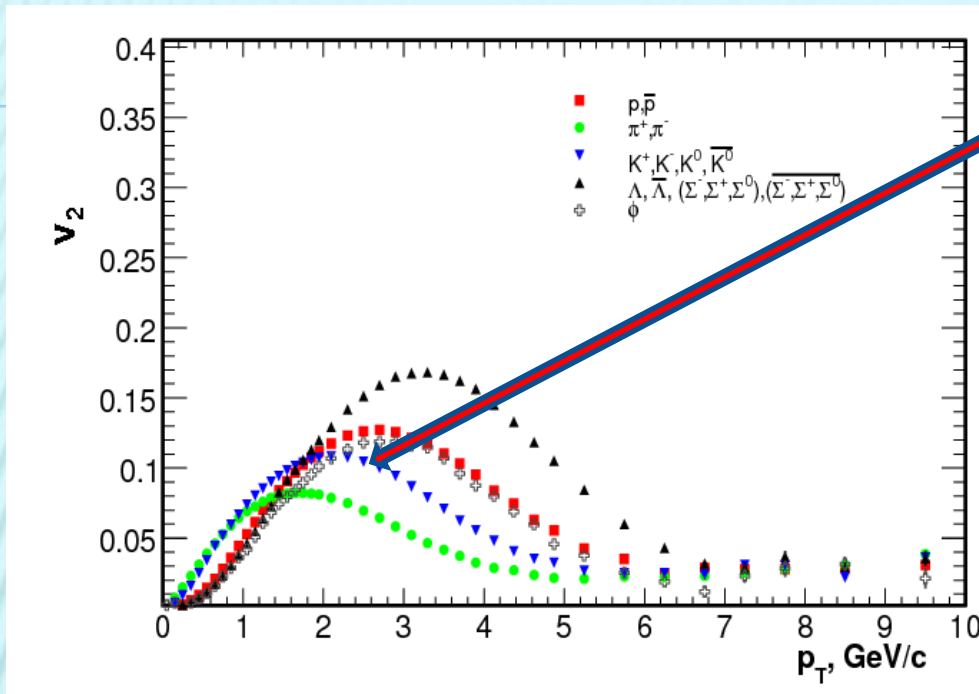
Au+Au @ 200 AGeV



Elliptic flow

G. Eyyubova et al., PRC 80 (2009) 064907;
N.S. Amelin et al., PRC 77 (2008) 014903

V_2 in HYDJET++ for different particles (centrality 30%)



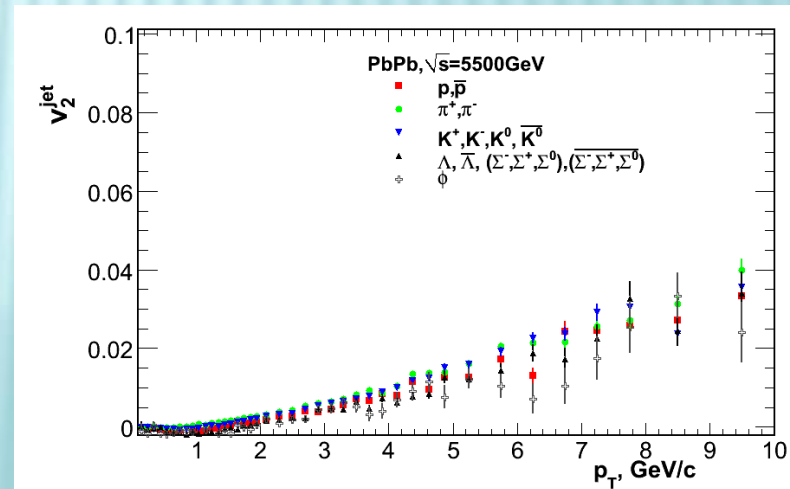
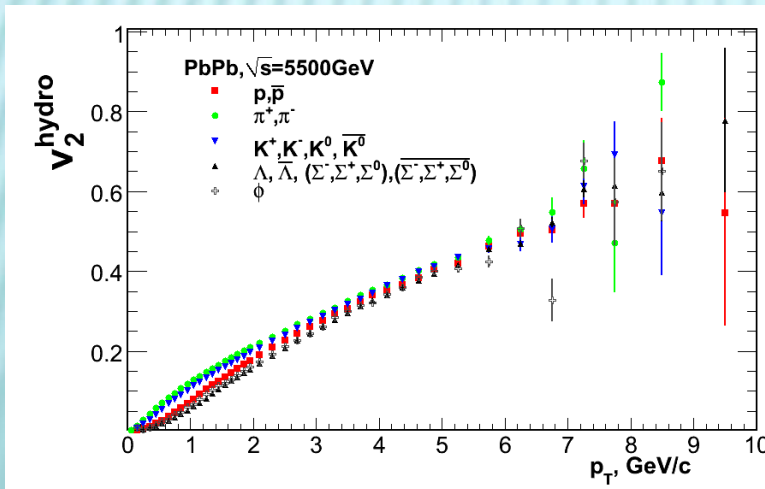
Mass ordering in soft p_T regions then breaks.

Why?

Hydrodynamics gives **mass ordering** of v_2 .
 The model possesses **crossing** of baryon and meson branches.

Hydrodynamics

Jet part + quenching



Interplay between hydrodynamics and jets

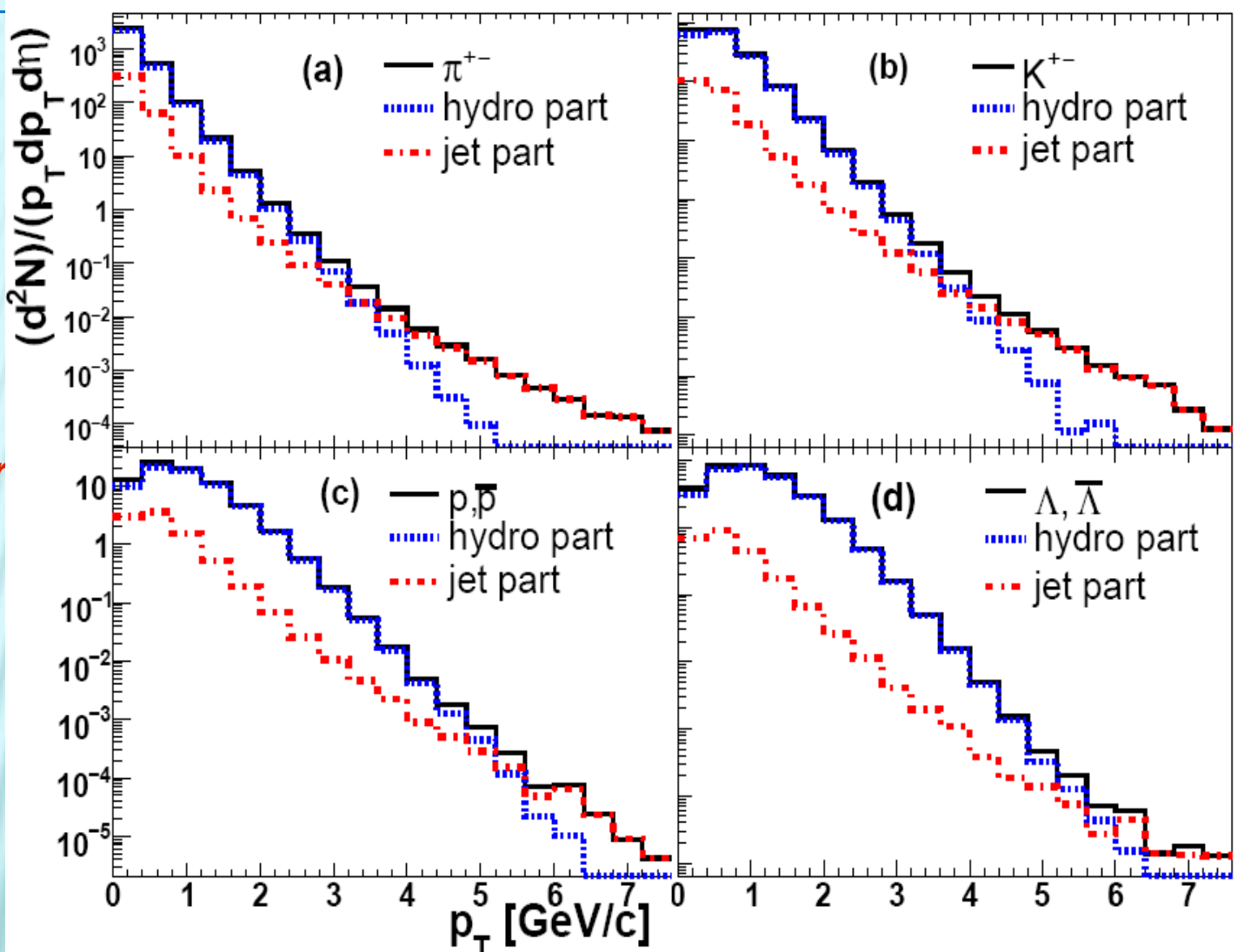
The p_T spectra of π, K, p, Λ with HYDJET++ model, $\sqrt{s}=200\text{GeV}$

The slope for the hydro part depends strongly on mass:

- the heavier the particle -- the harder the spectrum



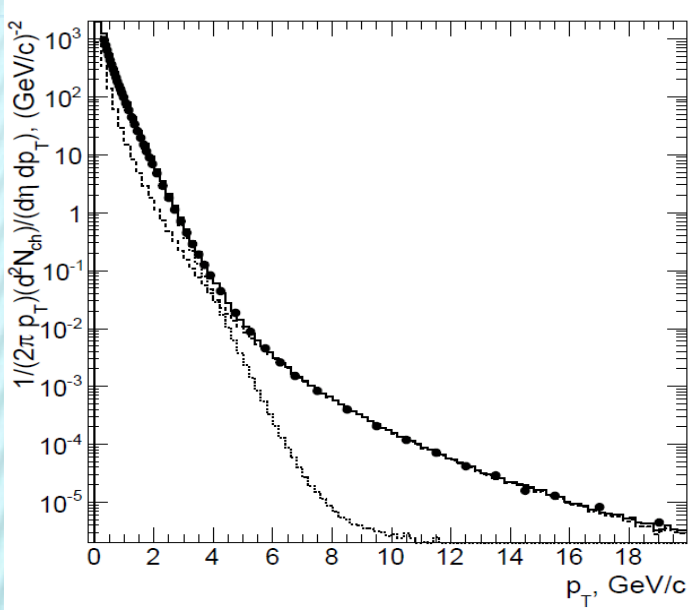
The hydro part dies out earlier for light particles than for heavy ones



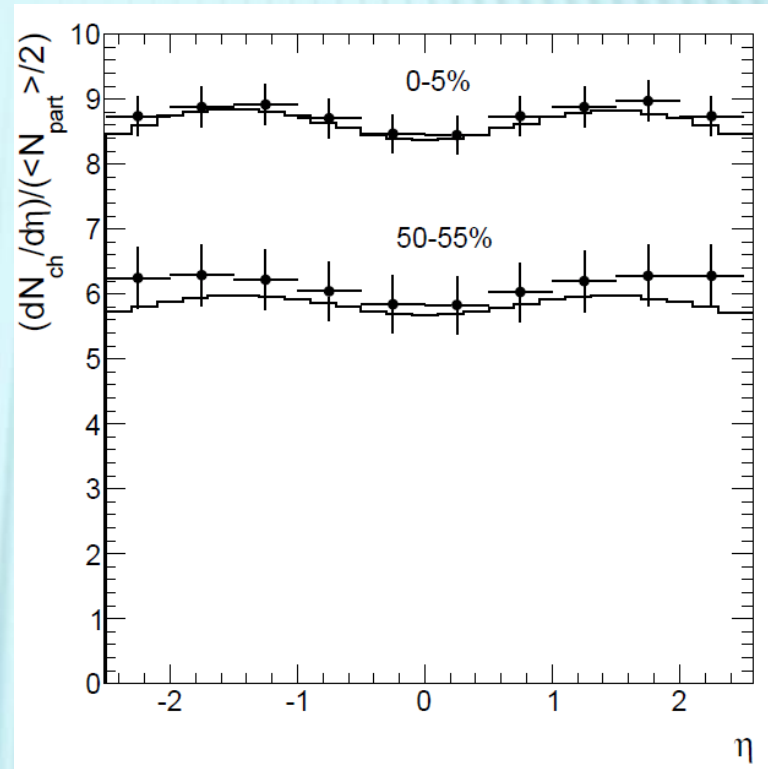
LHC DATA VS. HYDJET++ MODEL

Transverse momentum

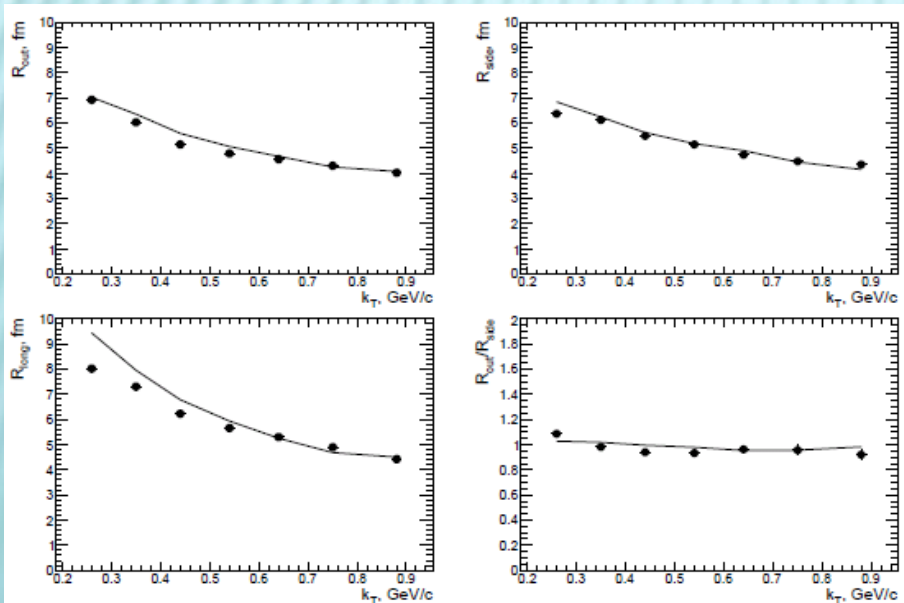
Pb+Pb @ 2.76 ATeV



Rapidity



I. Lokhtin et al., arXiv:1204.4820

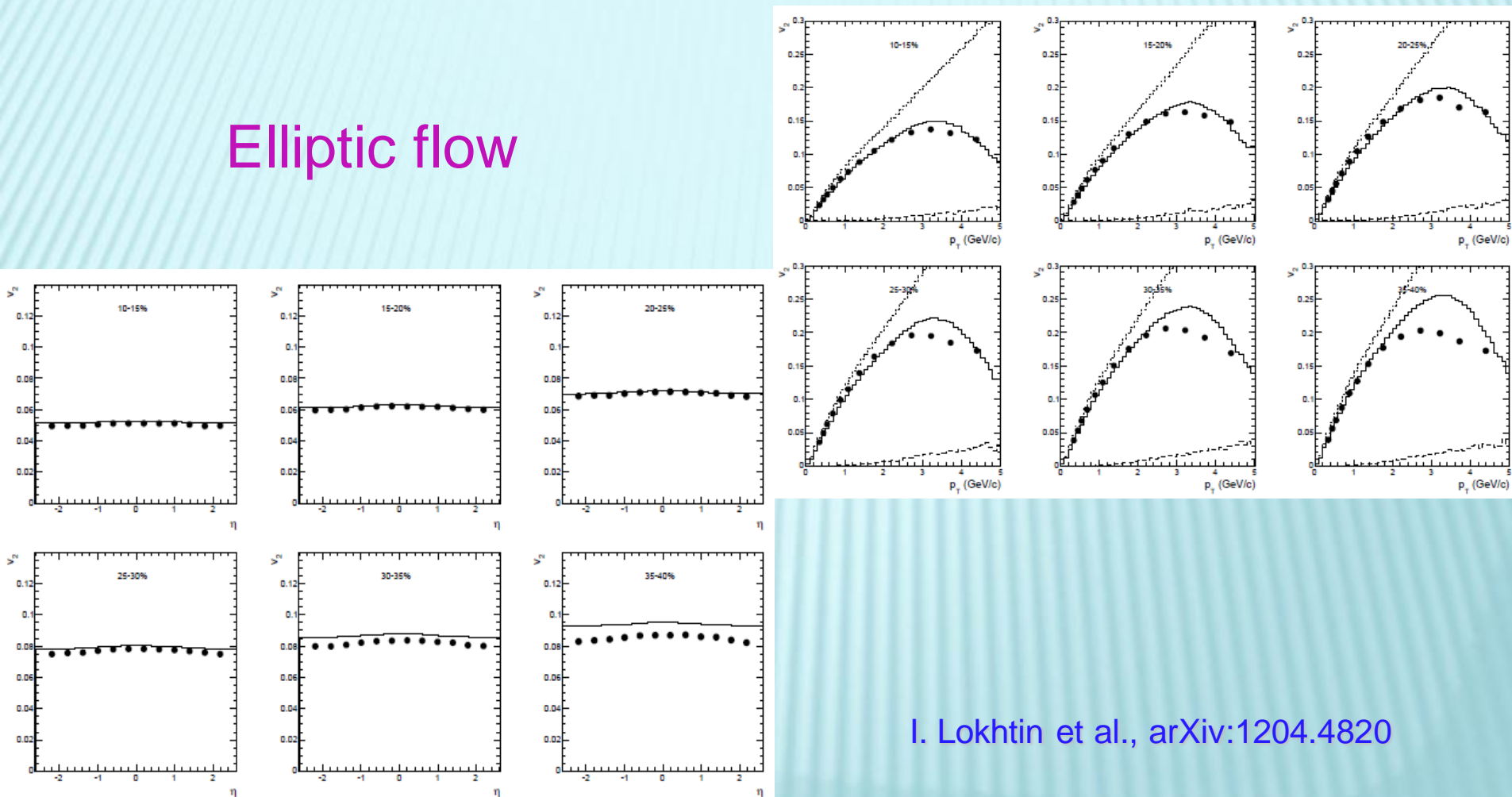


Correlation radii (femtoscropy)

LHC DATA VS. HYDJET++ MODEL

Pb+Pb @ 2.76 ATeV

Elliptic flow



I. Lokhtin et al., arXiv:1204.4820

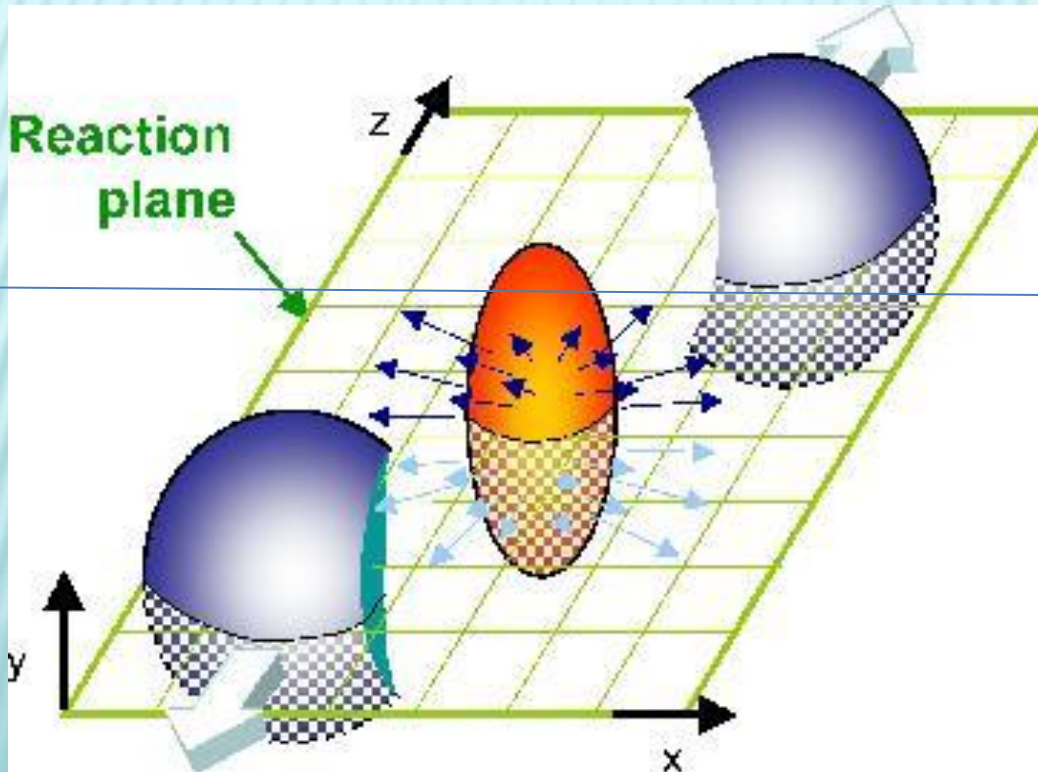
Model gives a fair description of various observables at both RHIC and LHC

II. $V_4/(V_2*V_2)$ RATIO

II. $v_4/(v_2)^2$ ratio

Anisotropic flow

$$\frac{dN}{d\varphi} = \frac{1}{2\pi} \left(1 + \sum_{n=1}^{\infty} 2v_n(p_t) \cos[n(\varphi - \psi_r)] \right)$$



Predictions

N. Borghini, J.-Y. Ollitrault, PLB 642 (2006) 227

- Within the *approximation* that the particle momentum \mathbf{p} and the fluid velocity \mathbf{v} are parallel (valid for *large momentum* p_{\perp} and *low freeze-out temperature* T)

$$dN/d\varphi = \exp(2\varepsilon p_{\perp} \cos(2\varphi)/T)$$

- Expanding to order ε , the $\cos(2\varphi)$ term is

$$v_2 = \varepsilon p_{\perp}/T$$

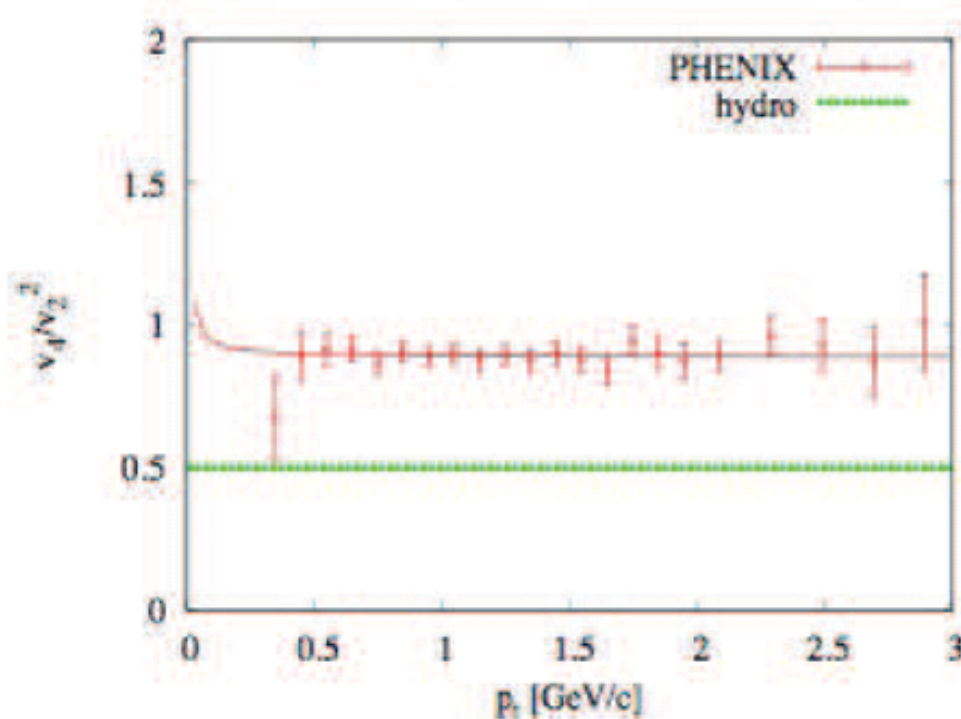
- Expanding to order ε^2 , the $\cos(4\varphi)$ term is

$$v_4 = \frac{1}{2} (v_2)^2$$

Hydrodynamics has a universal prediction for $v_4/(v_2)^2$!

Should be independent of equation of state, initial conditions, centrality, rapidity, particle type

Comparison with data



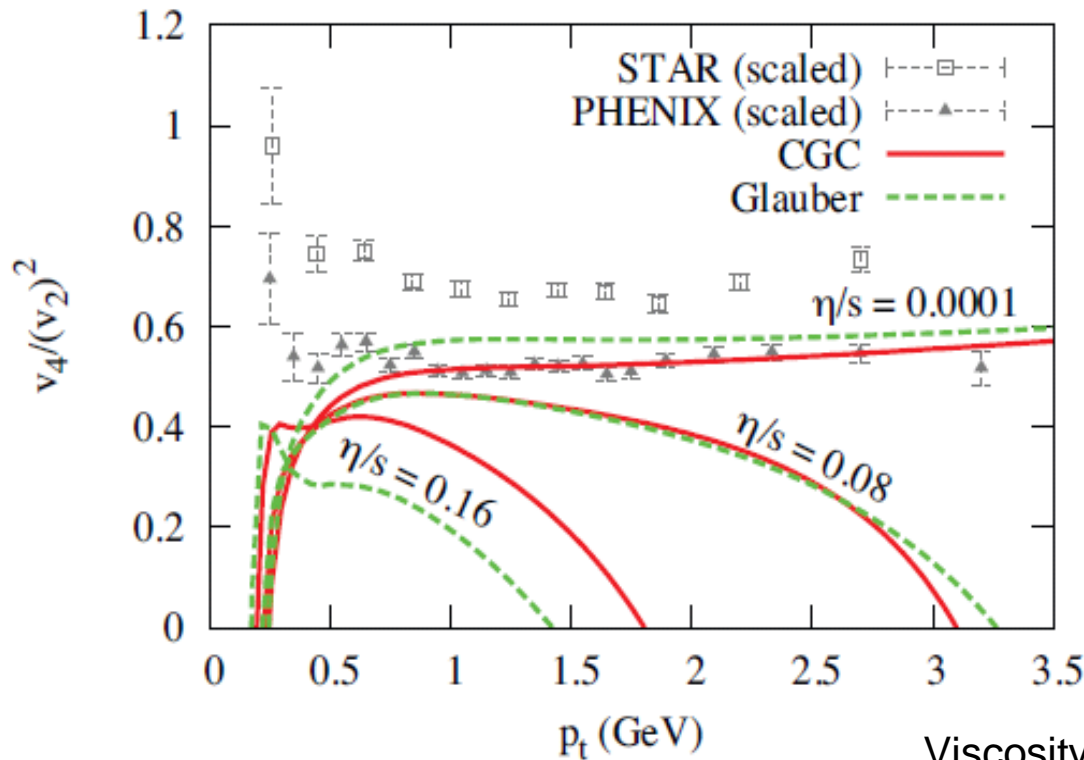
PHENIX data for charged pions

Au-Au collisions at 100+100 GeV

20-60% most central

The ratio is significantly larger than 0.5.
Can this be explained by viscous corrections?

Effects of initial profile and viscosity

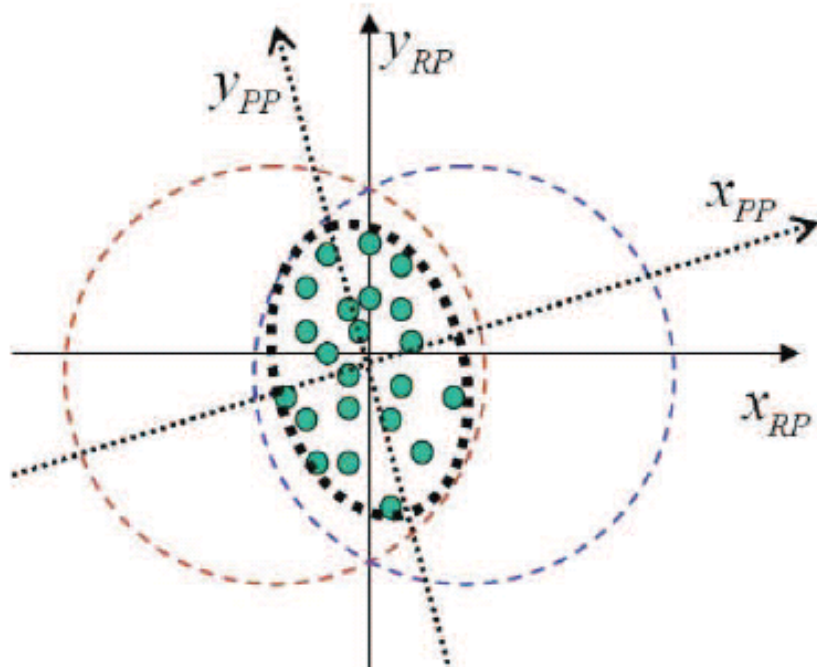


Initial profile has little effect although eccentricities differ.

results strongly depend on viscosity

Viscosity lowers $v_4/(v_2)^2$ for a realistic T_f

Eccentricity fluctuations



Depending on where the participant nucleons are located within the nucleus at the time of the collision, the actual shape of the overlap area may vary: the **orientation and eccentricity** of the ellipse defined by participants fluctuates.

Assuming that v_2 scales like the eccentricity, **eccentricity fluctuations** translate into **v_2 fluctuations**

Eccentricity fluctuation can be computed in MC Glauber model or derived from experiment by comparing different methods for flow calculation.

Why ε fluctuations change v_4/v_2^2

Experimentally, no direct measure of v_2 and v_4

v_2 and v_4 are measured via azimuthal correlations

$$v_2 \text{ from } \langle \cos(2\phi_1 - 2\phi_2) \rangle = \langle (v_2)^2 \rangle$$

$$v_4 \text{ from } \langle \cos(4\phi_1 - 2\phi_2 - 2\phi_3) \rangle = \langle v_4 (v_2)^2 \rangle$$

Experimentally measured

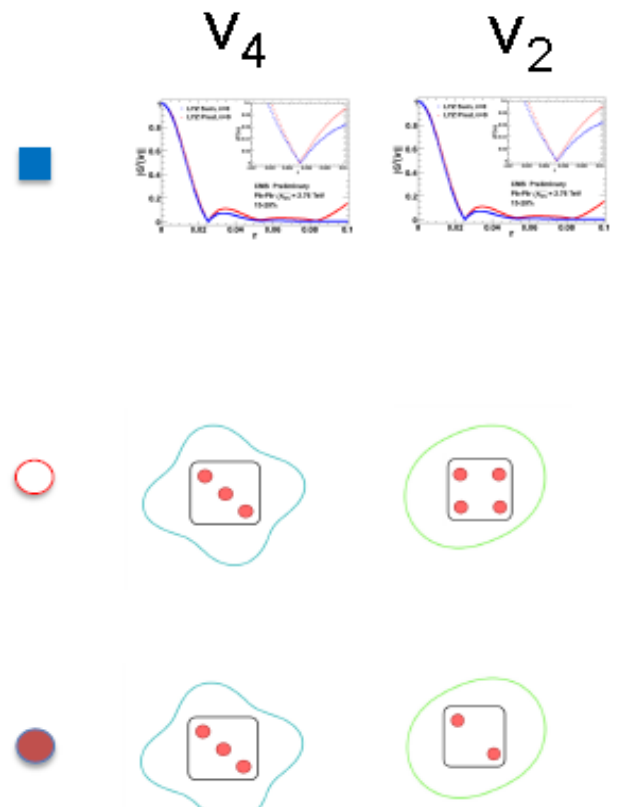
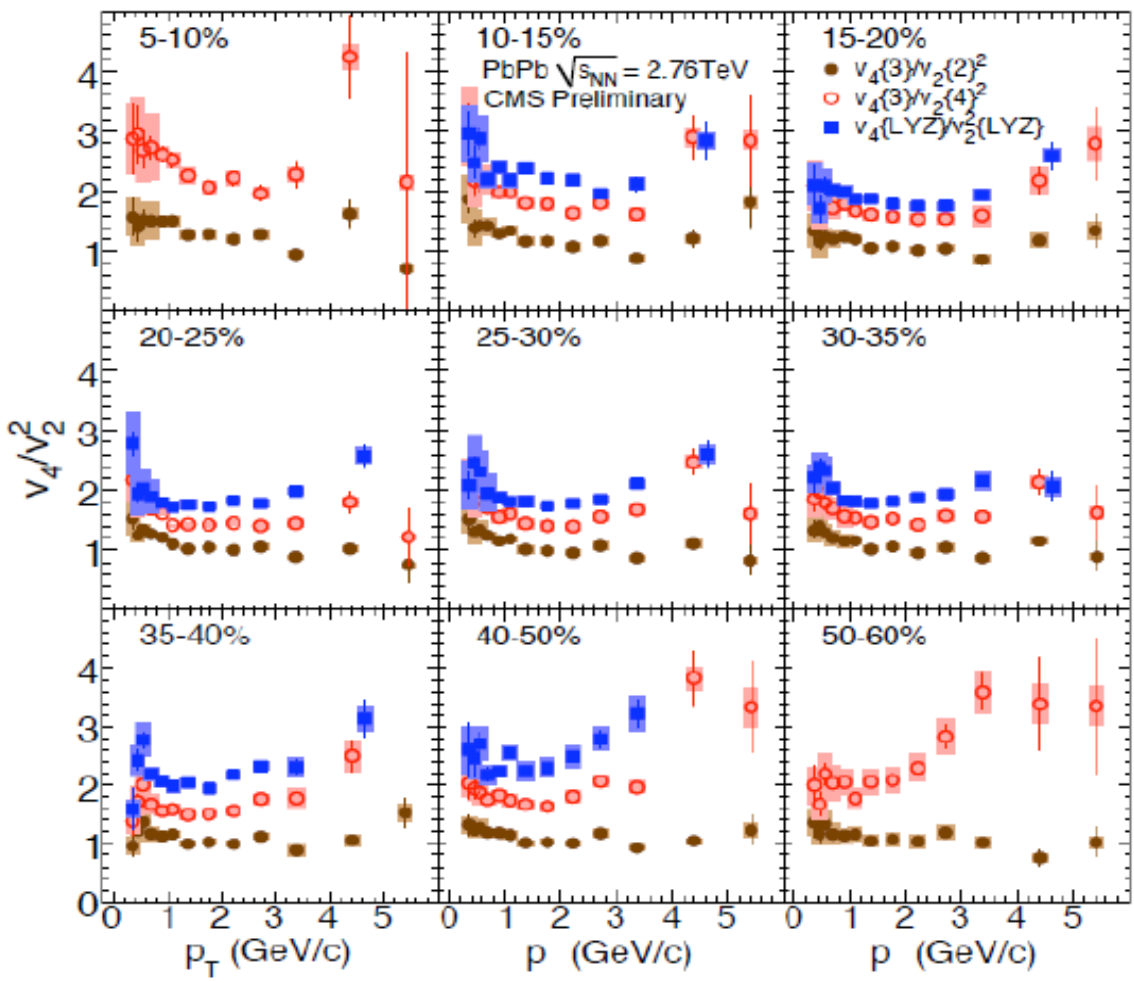
$$\frac{v_4}{v_2^2} = \frac{\langle v_4 (v_2)^2 \rangle}{\langle (v_2)^2 \rangle^2} = \frac{1}{2} \frac{\langle (v_2)^4 \rangle}{\langle (v_2)^2 \rangle^2} > \frac{1}{2}$$

fluctuations

hydro

Similar results obtained using Event Plane method

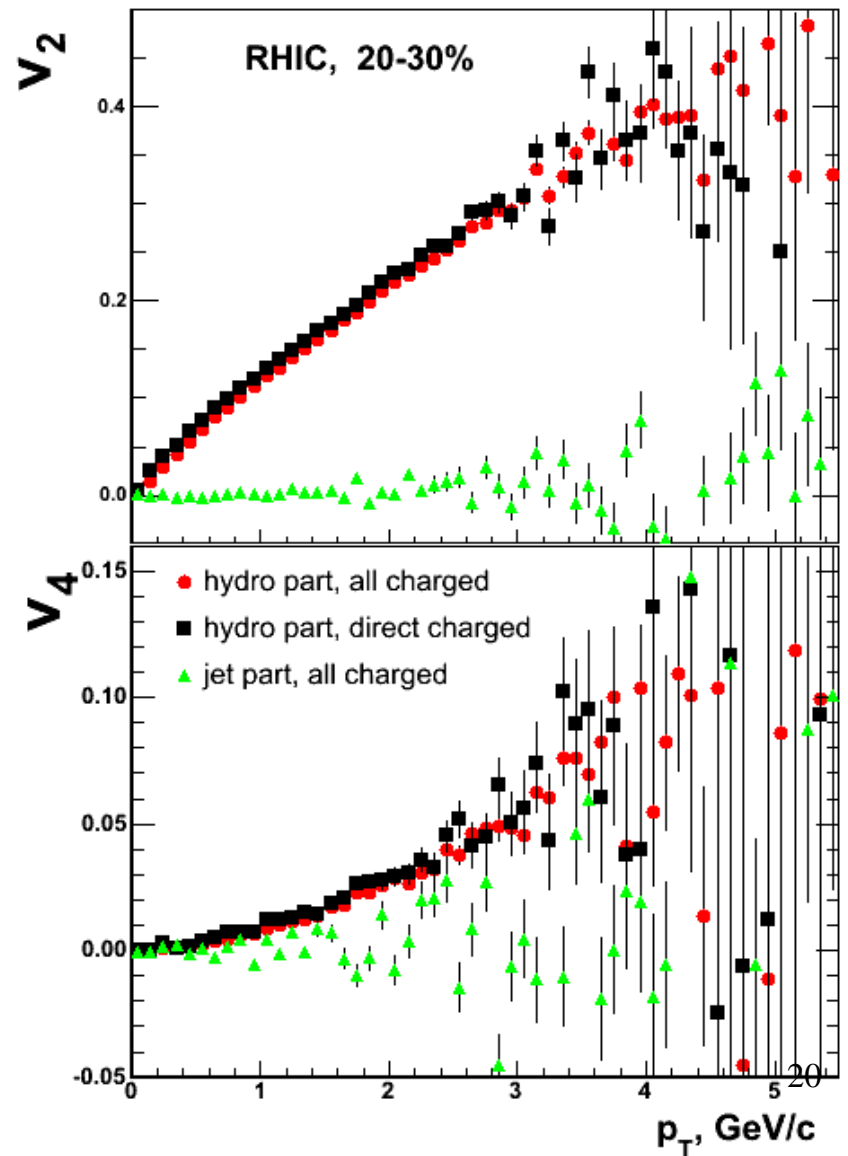
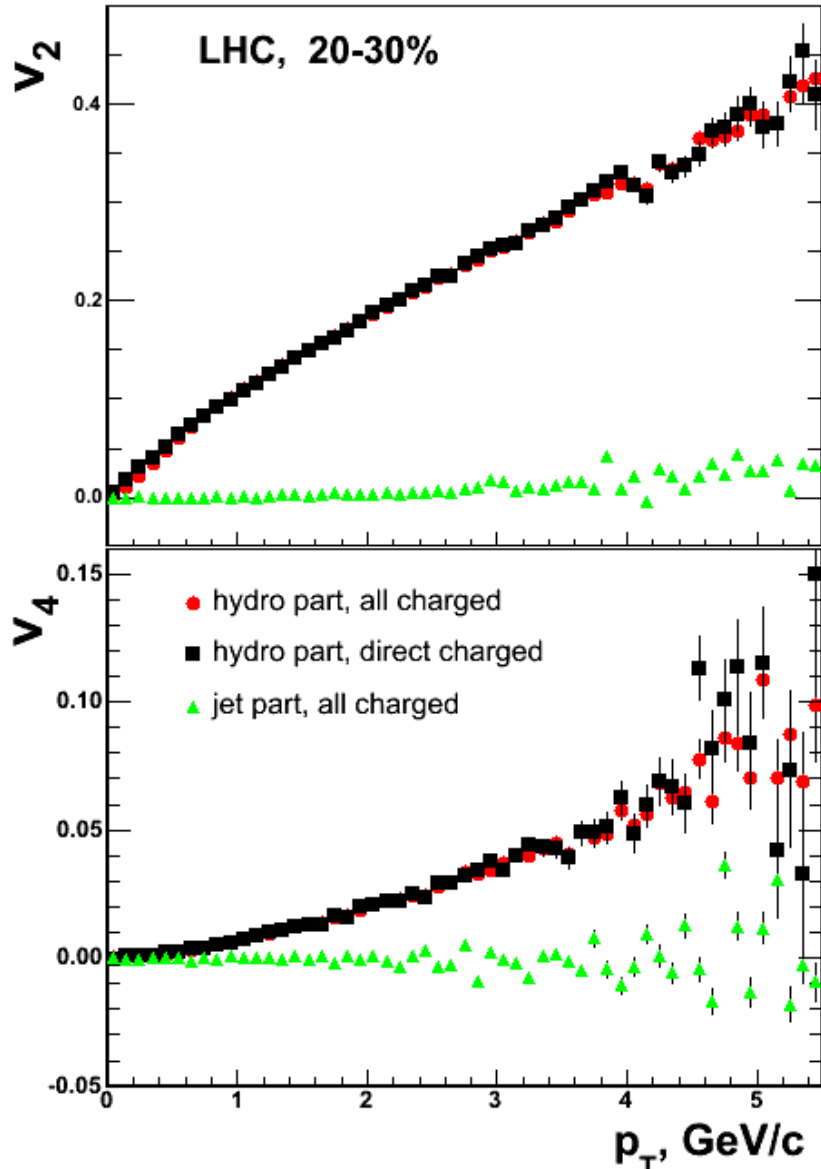
$v_4 / v_2^2(p_T)$ at mid-rapidity $|\eta| < 0.8$



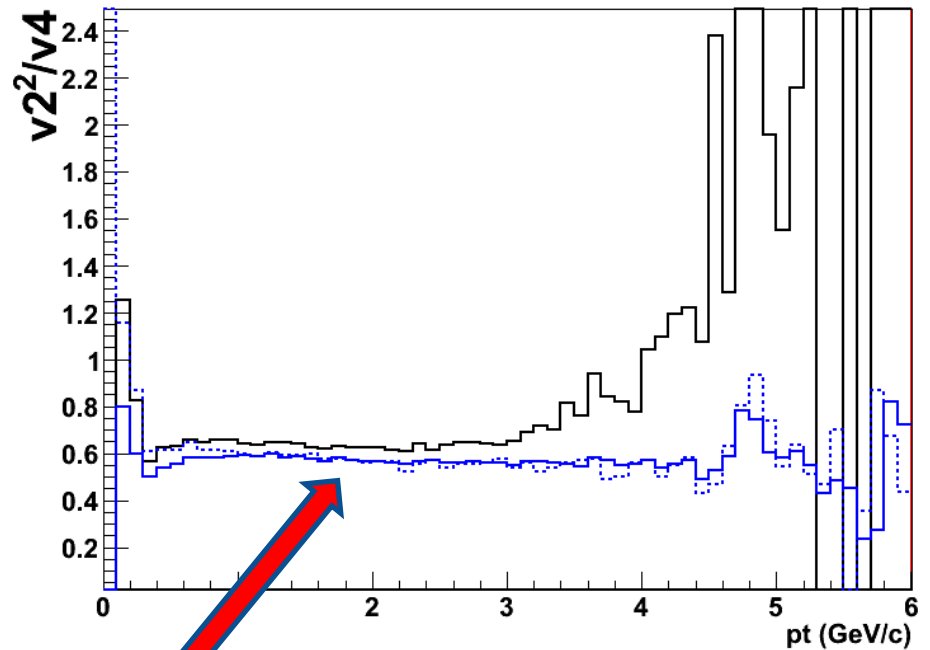
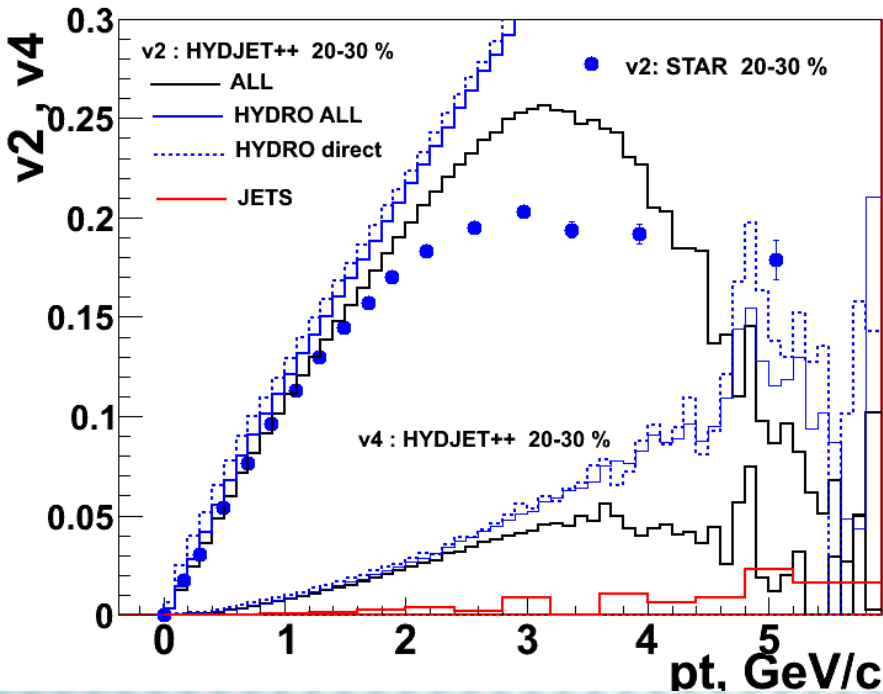
Significantly higher than RHIC: experimental method dependent

HYDJET++

Effects to be studied: resonance decay and hard part influence



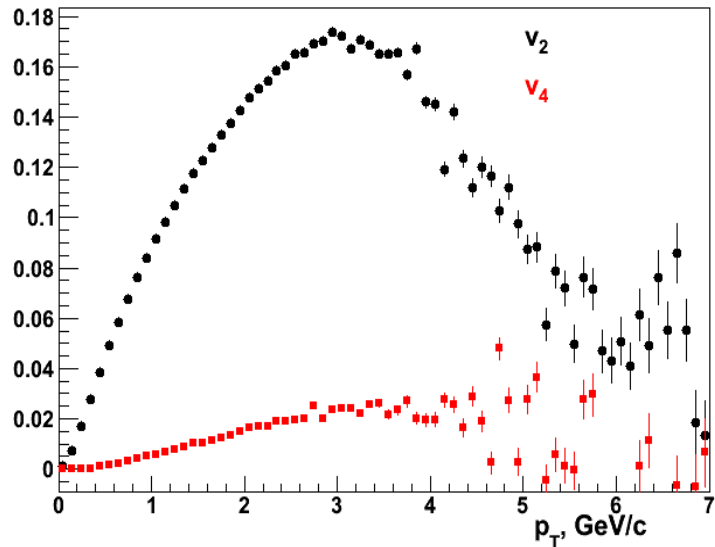
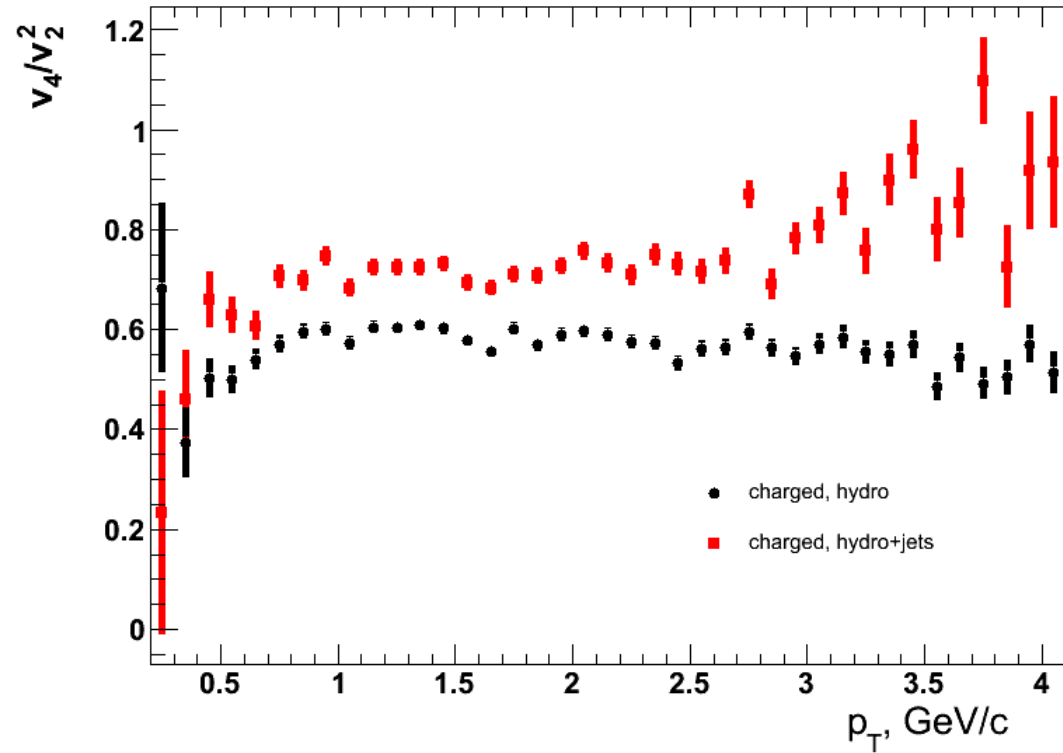
HYDJET++ RESULTS FOR RHIC



Jets increase the ratio

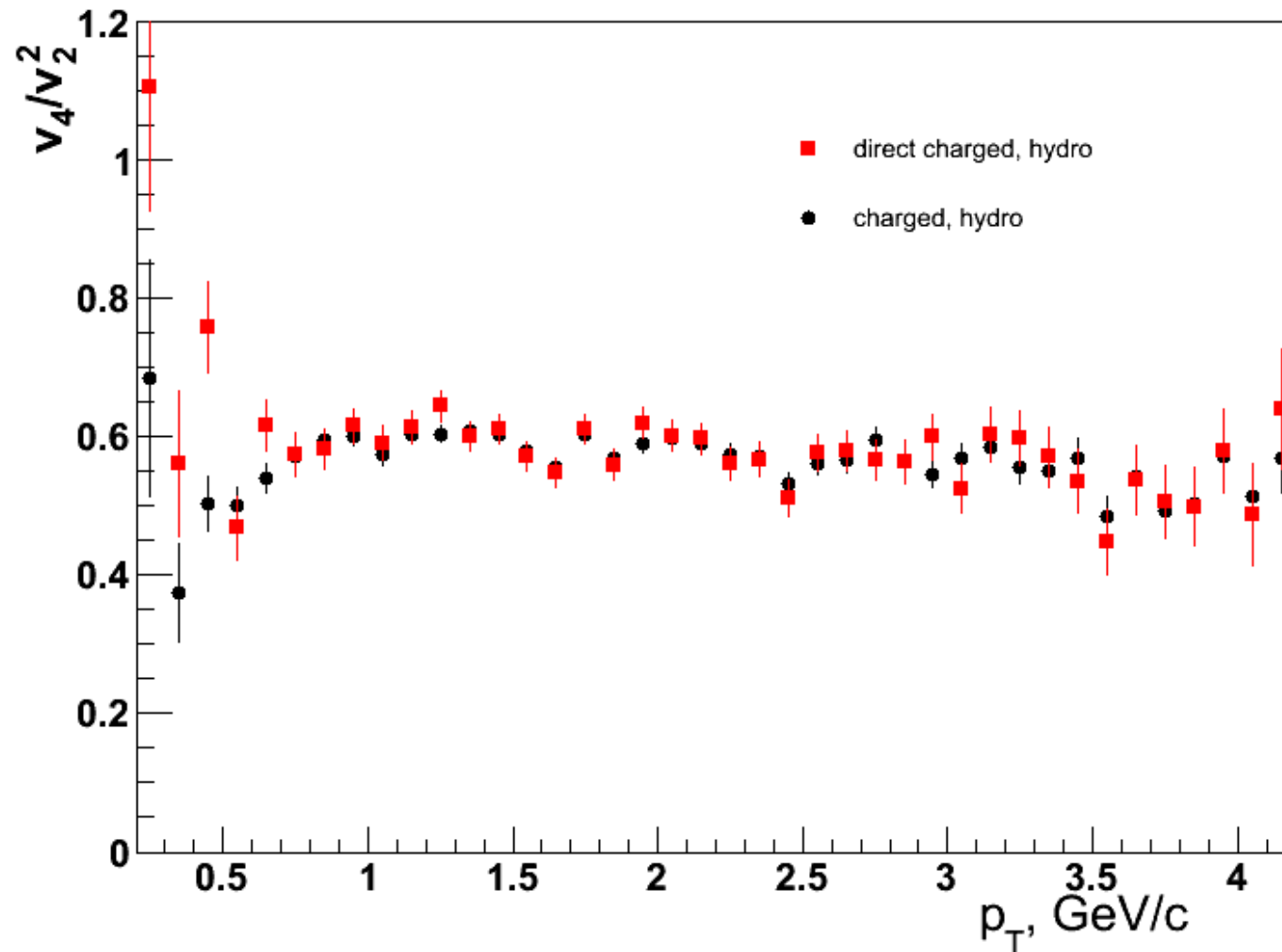
HYDJET++ RESULTS FOR LHC

The same tendency is observed in Pb+Pb at LHC



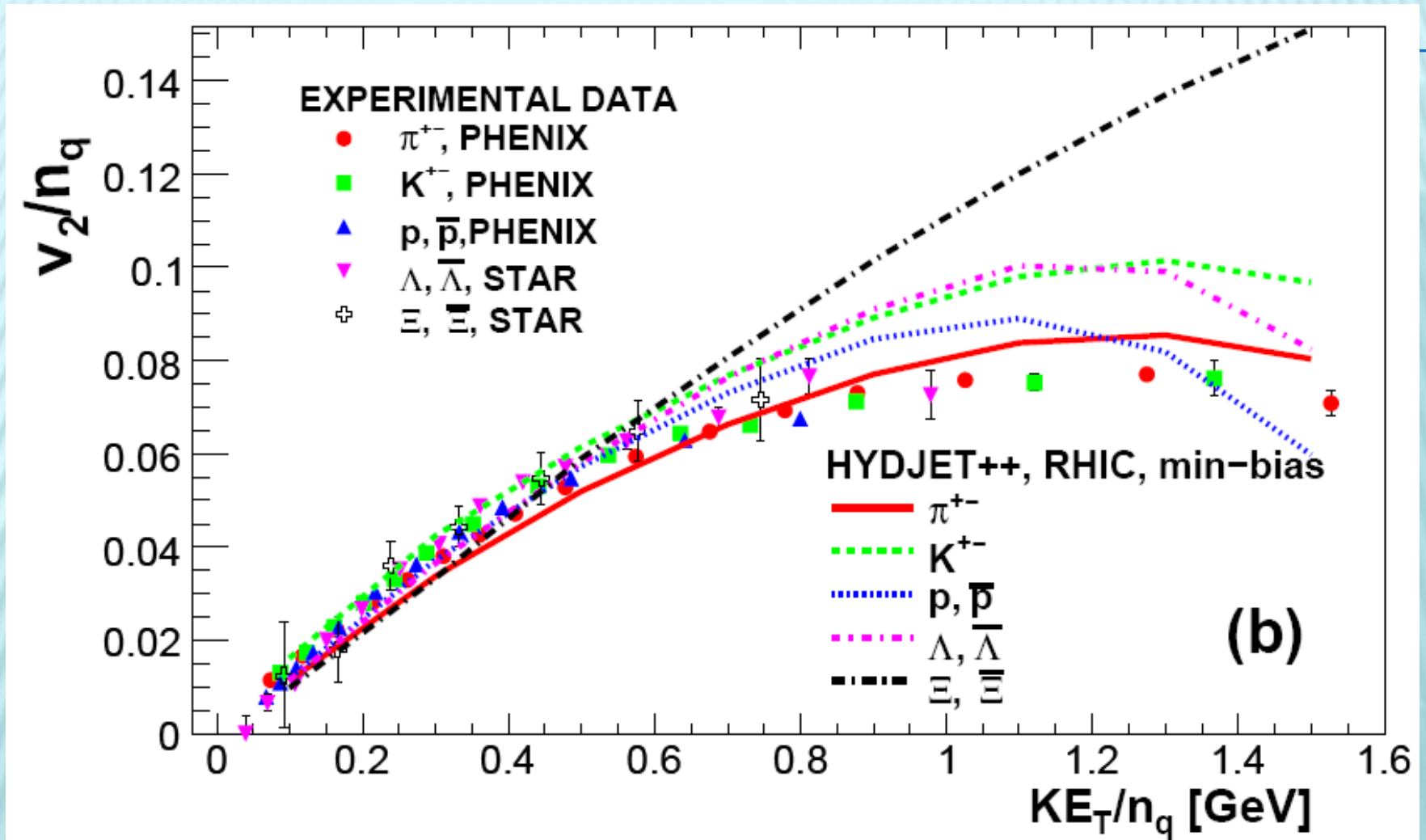
Still, the ratio is below 1

DECAYS OF RESONANCES PLAY MINOR ROLE



III. Number-of- constituent- quark (NCQ) scaling

COMPARISON WITH RHIC DATA

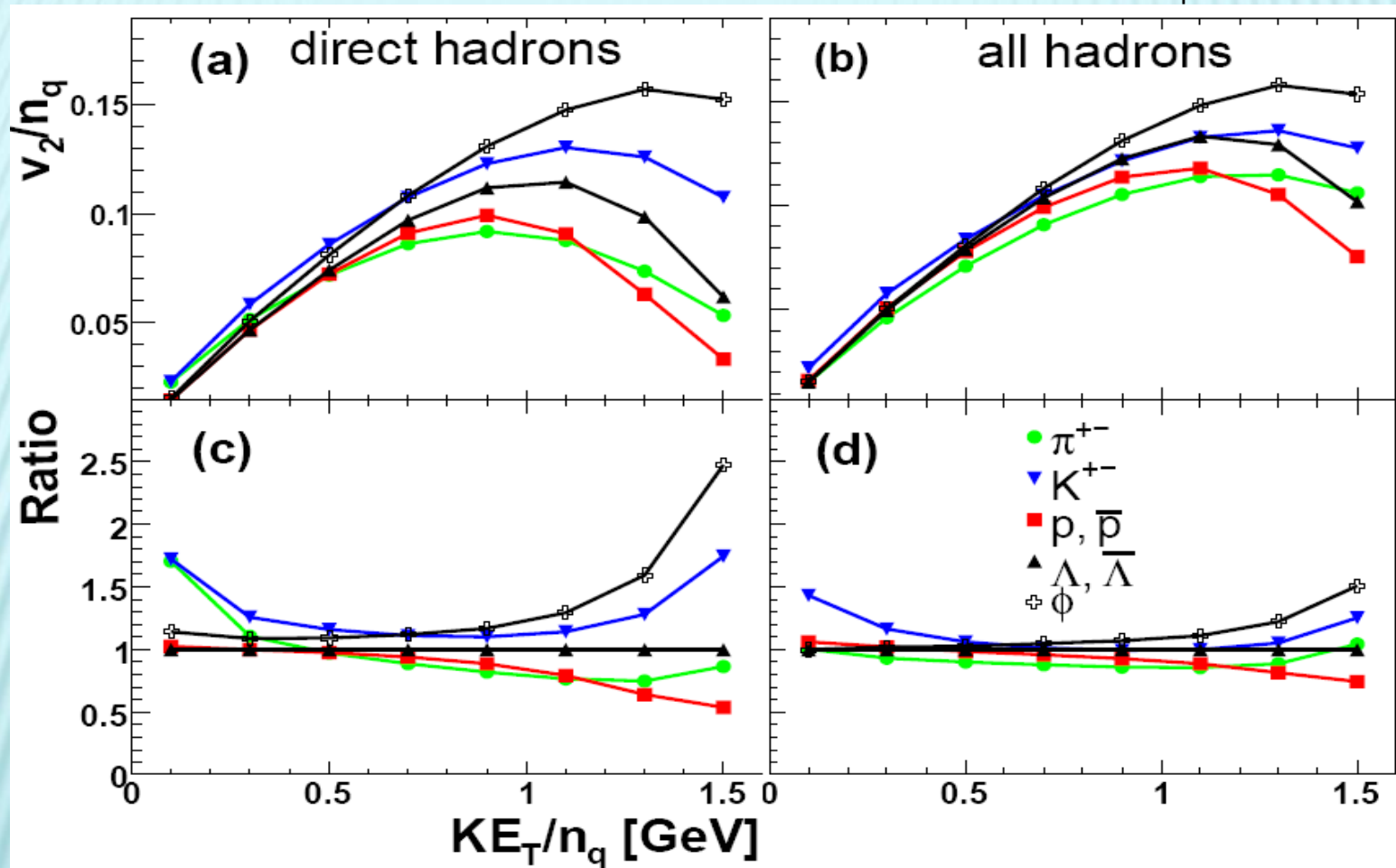


The agreement seems to be good at $KE_T/n_q < 0.7$ GeV

Number-of-constituent-quark scaling at RHIC

Direct particles: scaling is not good.

All particles: KE_T/n_q scaling

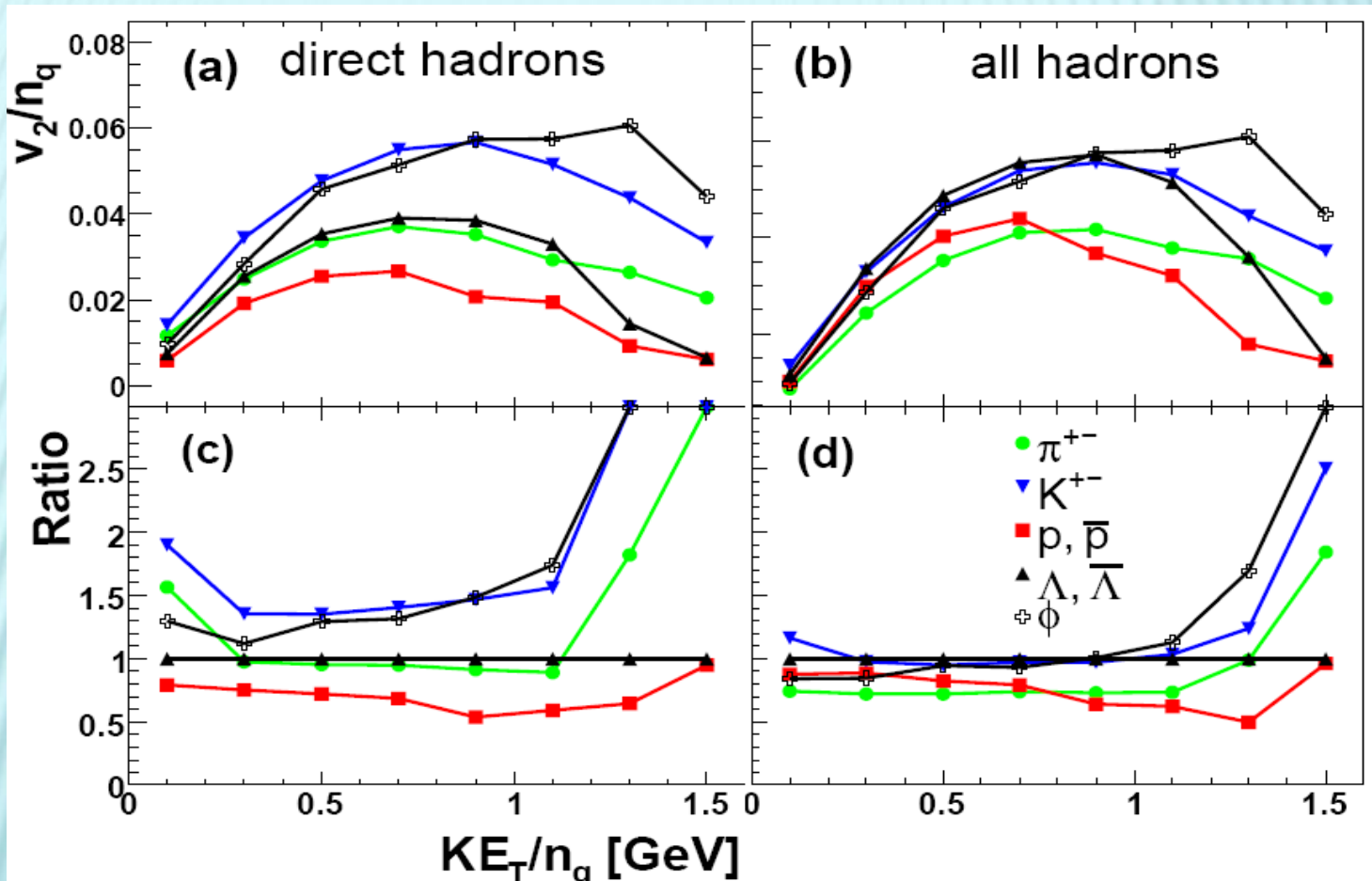


One of the explanations of KE_T/n_q scaling is partonic origin of the elliptic flow. *However, final state effects (such as resonance decays and jets) may also lead to appearance of the scaling*

NCQ scaling at LHC

No scaling for
direct particles

Appearance of the approximate
scaling for all particles



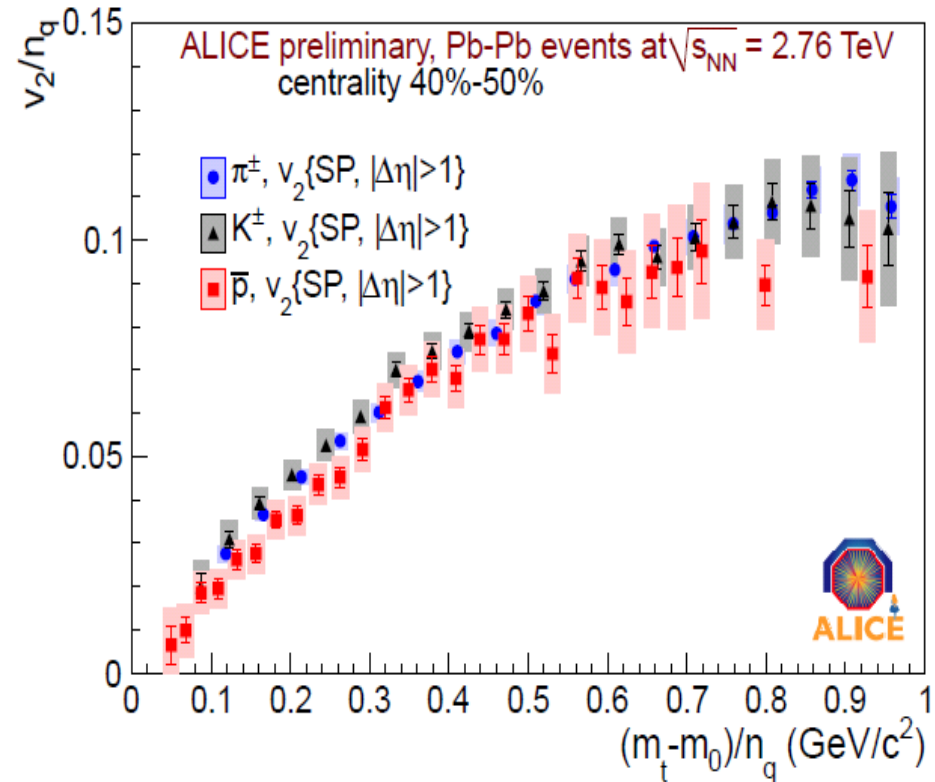
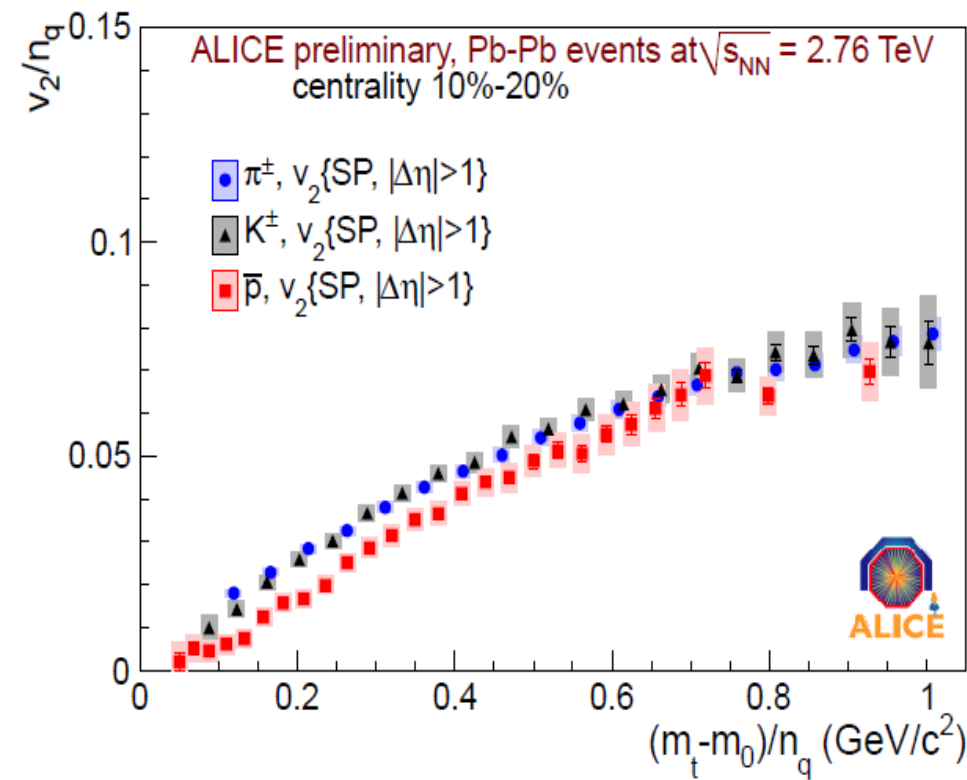
LHC: NCQ scaling will be only approximate (prediction, 2009)

Experimental results (LHC)

ALICE collab., M. Krzewicki et al., JPG 38 (2011) 124047

Semi-central collisions

Semi-peripheral collisions



The NCQ scaling is indeed only approximate (2011)

CONCLUSIONS

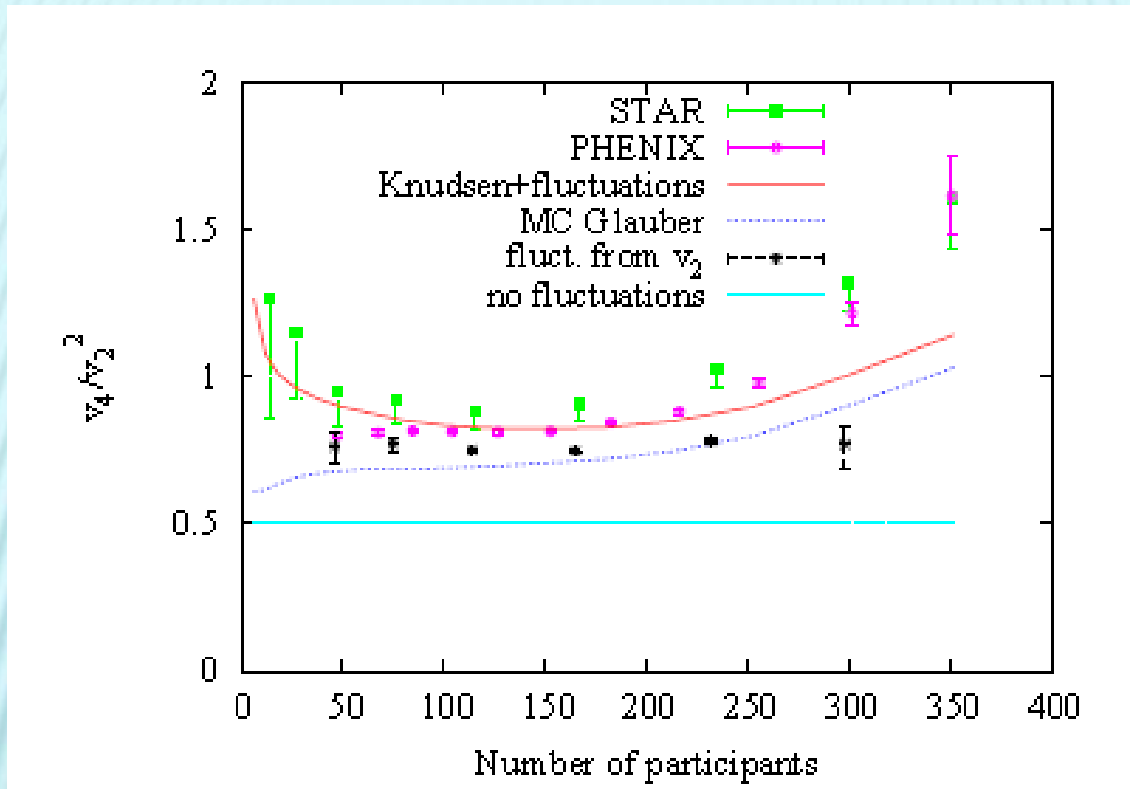
The HYDJET++ model allows to investigate flow of hydro and jet parts separately, to look at reconstruction of pure hydro flow and its modification due to jet part.

- *Jets result to increase by 25% - 30% of the ratio $v_4/(v_2*v_2)$*
- *Eccentricity fluctuations can increase the ratio by factor 1.5*
- *Jets + eccentricity fluctuations are enough to explain RHIC data*
- *For LHC we can explain 75% of the signal. Other effects are needed*
- *The predicted violation of the NCQ scaling at LHC is observed*

Back-up Slides

Effects of flow fluctuations and partial thermalization

M. Luzum, C. Gombeaud, J.-Y. Ollitrault, Phys.Rev.C81:054910,2010.

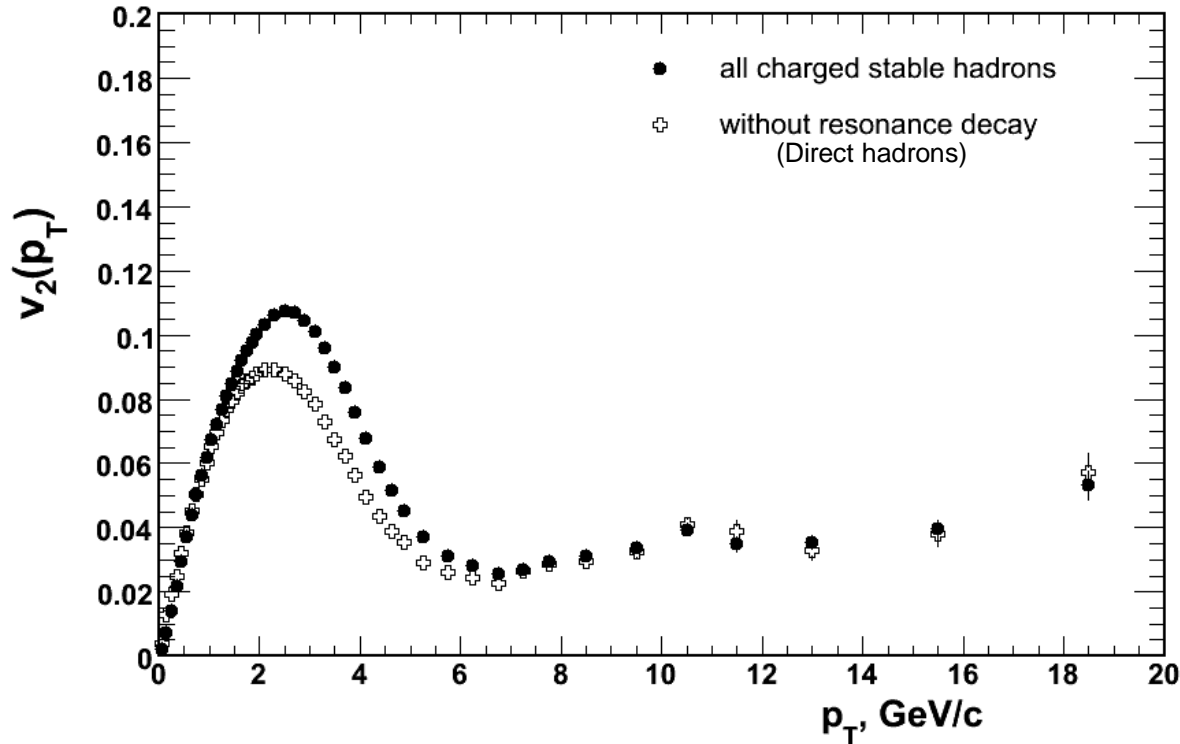


Stars: with fluctuations inferred from the difference between $v_2\{2\}$ and $v_2\{LYZ\}$.

Dotted line: eccentricity fluctuations from a Monte-Carlo Glauber

III . INFLUENCE OF RESONANCE DECAYS

Influence of resonance decay on v_2 value



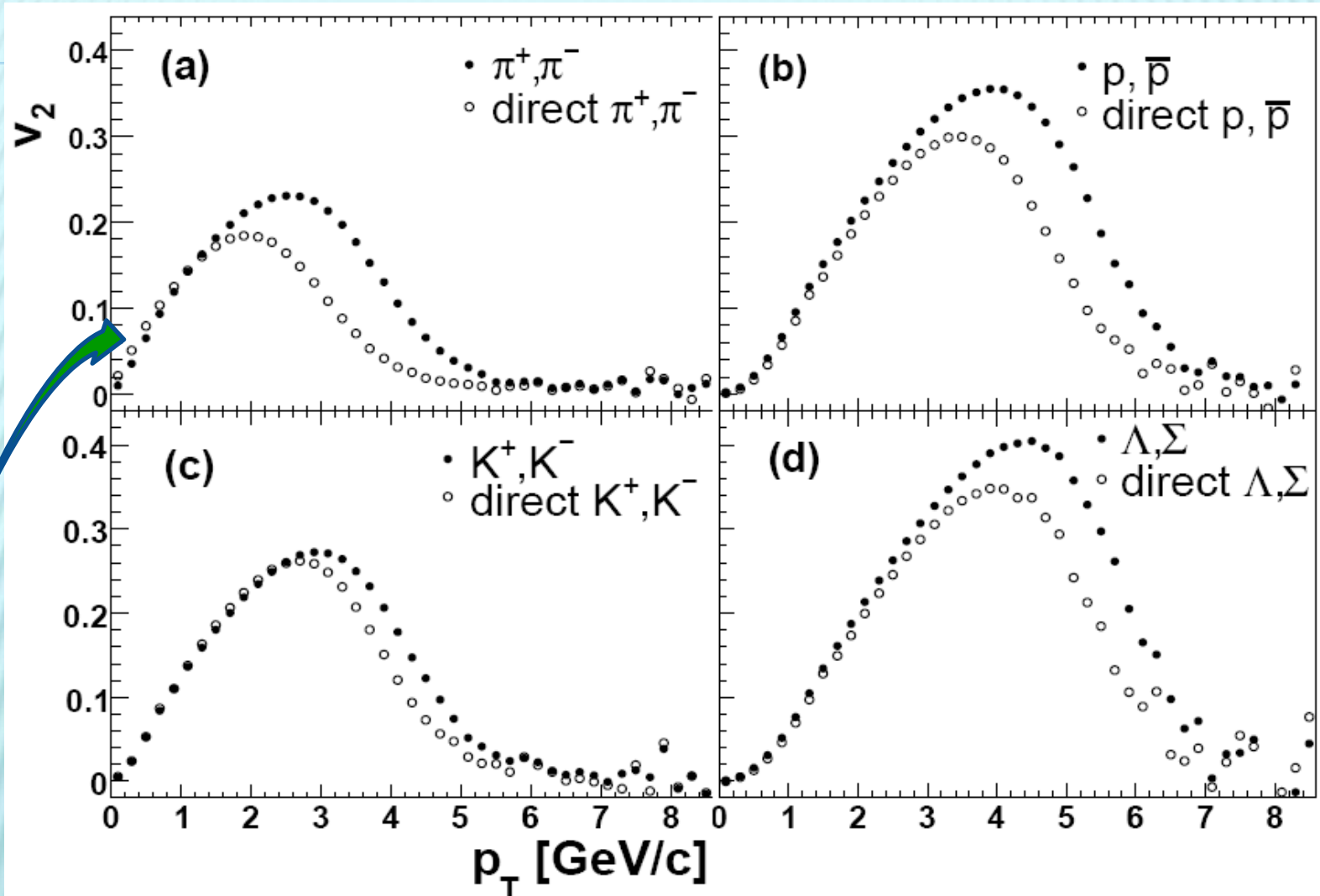
PbPb collisions, $c=30\%$

The elliptic flow of directly produced particles is smaller than that for all particles.

TABLE I: Yields of the particles produced directly and with resonance decays, $5.6 \cdot 10^6$ events, $c=42\%$, midrapidity

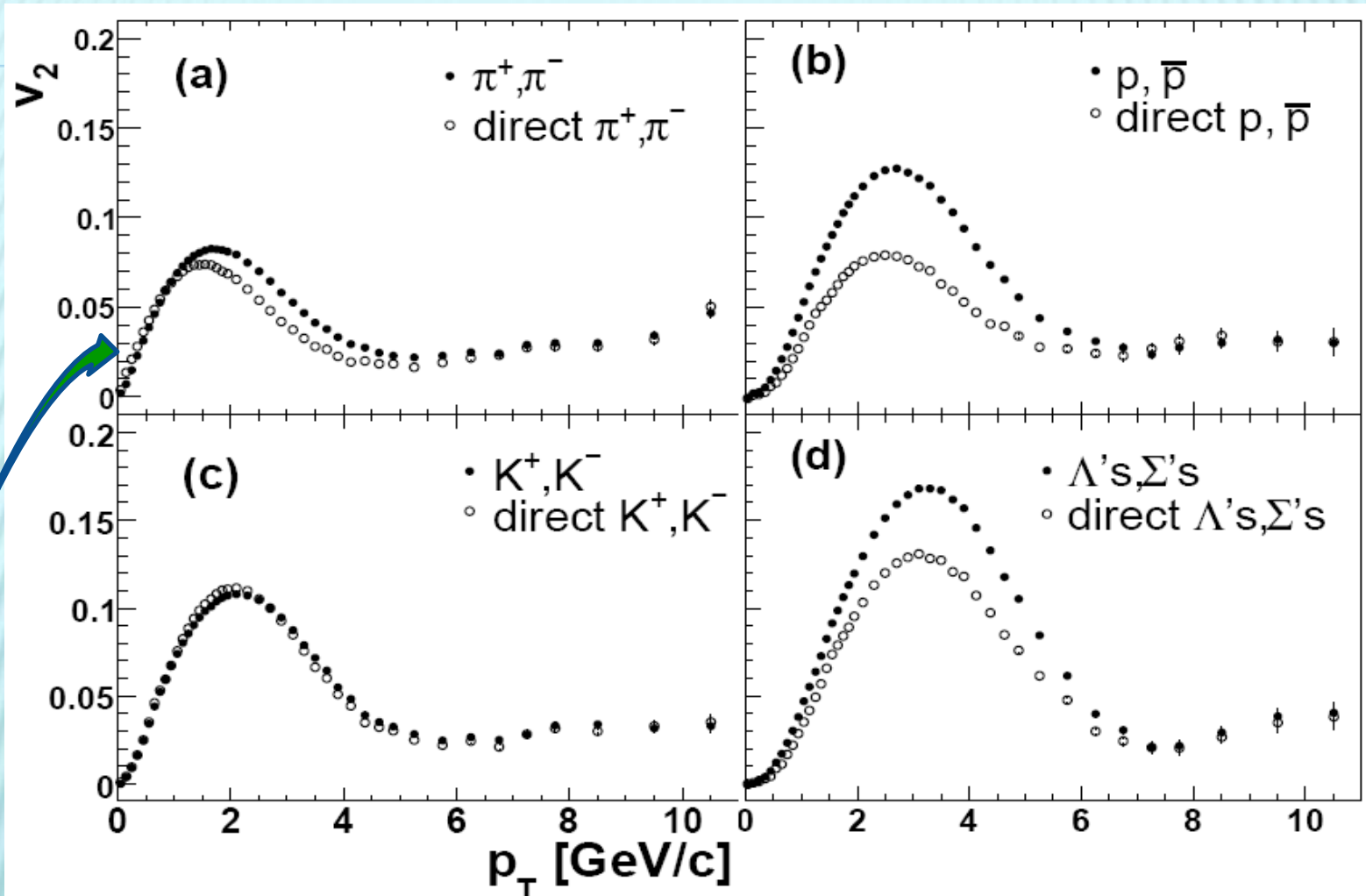
| | π^\pm | $K + \bar{K}$ | $p + \bar{p}$ | $\Lambda + \bar{\Lambda} + \Sigma + \bar{\Sigma}$ | ϕ |
|----------|-----------|---------------|---------------|---------------------------------------------------|--------|
| all | 860 | 185 | 63.8 | 42.3 | 6.55 |
| direct | 169 | 81.4 | 18.6 | 14.2 | 6.5 |
| direct % | 20 % | 44 % | 30 % | 39 % | 99 % |

Influence of resonance decays for different type of particles at RHIC



Pions and kaons: the resulting flow is weaker at low- p_T and larger at high- p_T
Baryons: the resulting flow is stronger than the flow of direct particles

Influence of resonance decays for different type of particles at LHC

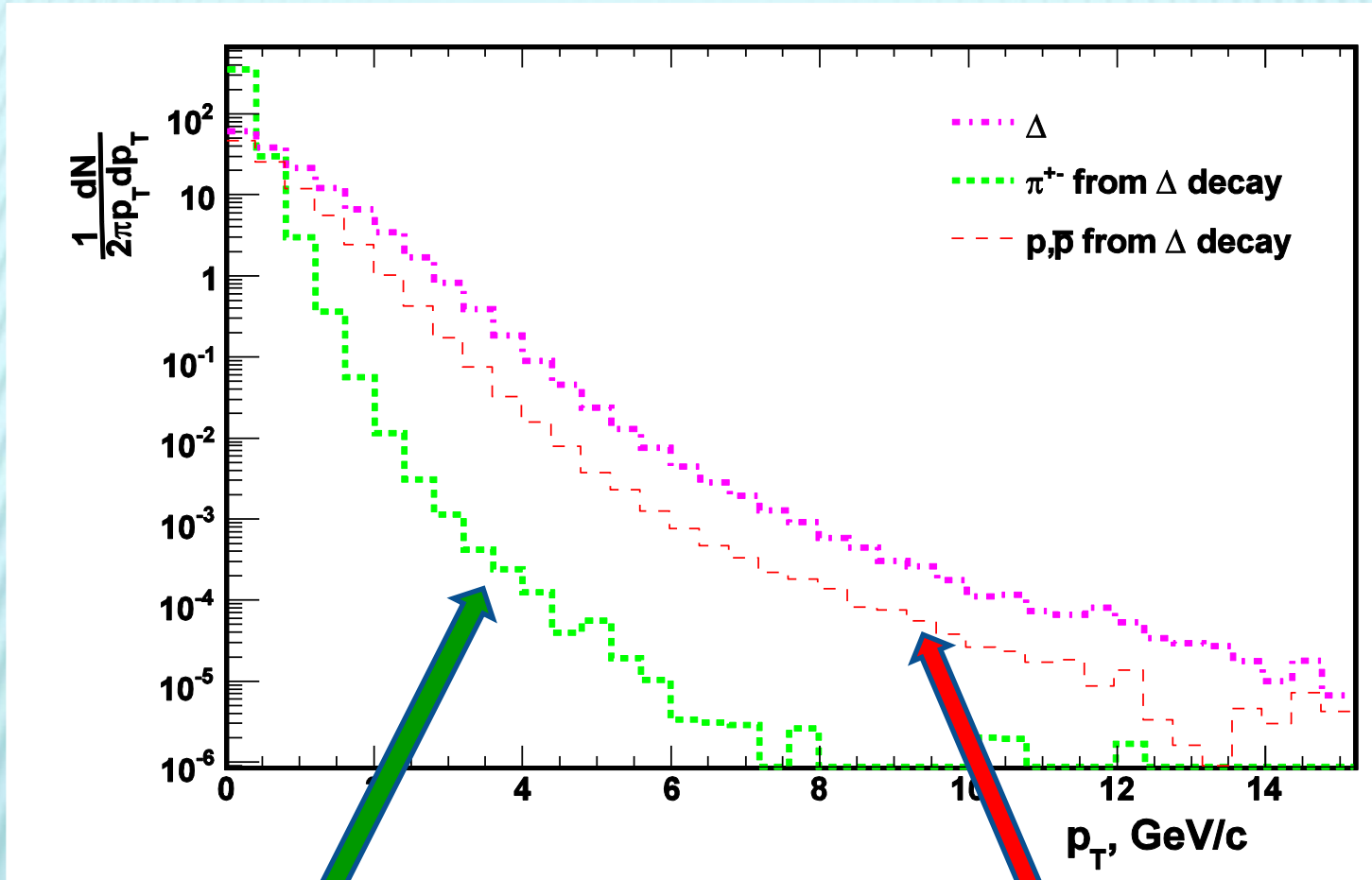


Pions: the resulting flow is weaker at low-pt and larger at high-pt

Kaons: both flows almost coincide

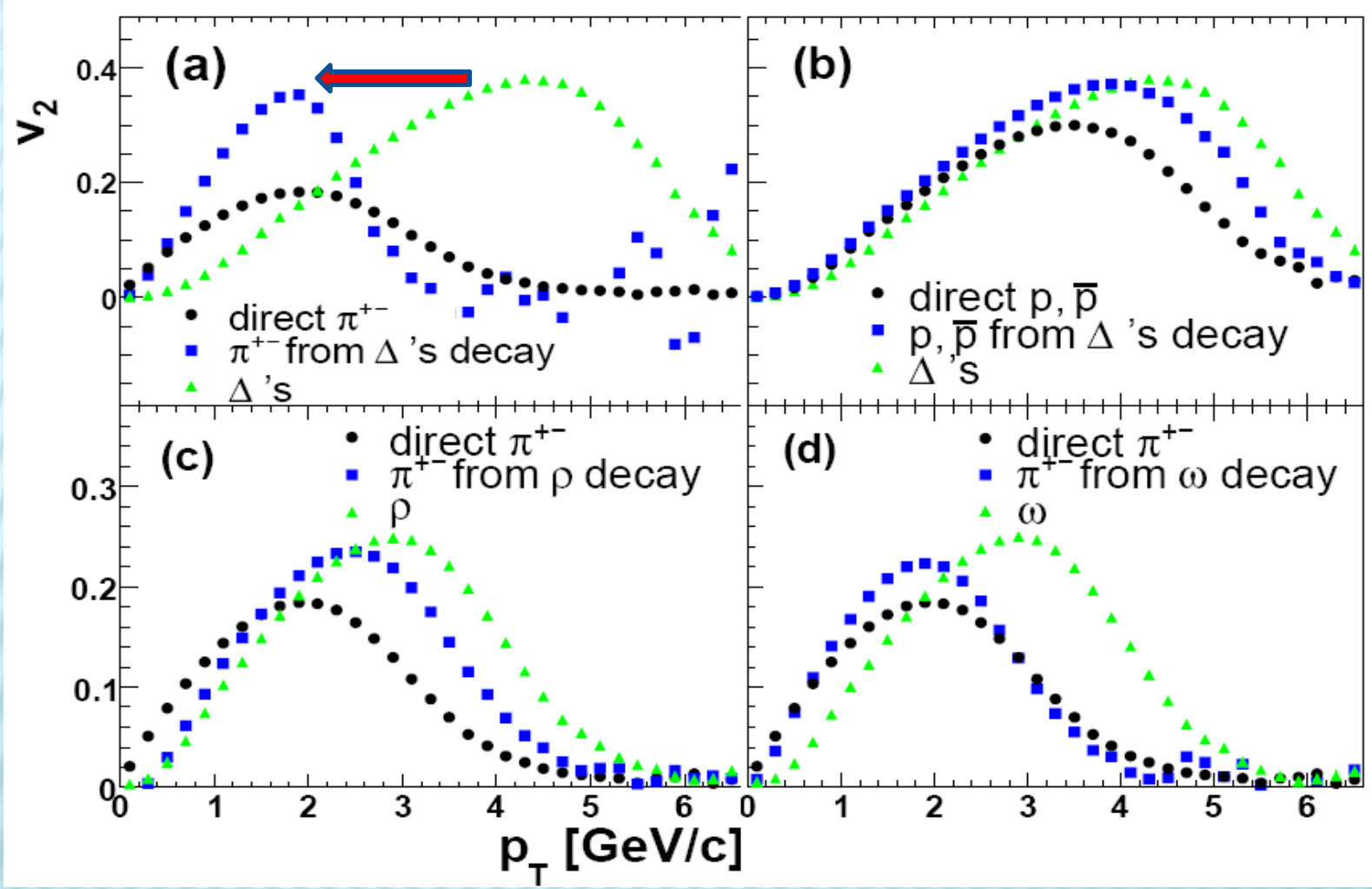
Baryons: the resulting flow is stronger than the flow of direct particles

TRANSVERSE MOMENTUM OF SECONDARY PARTICLES



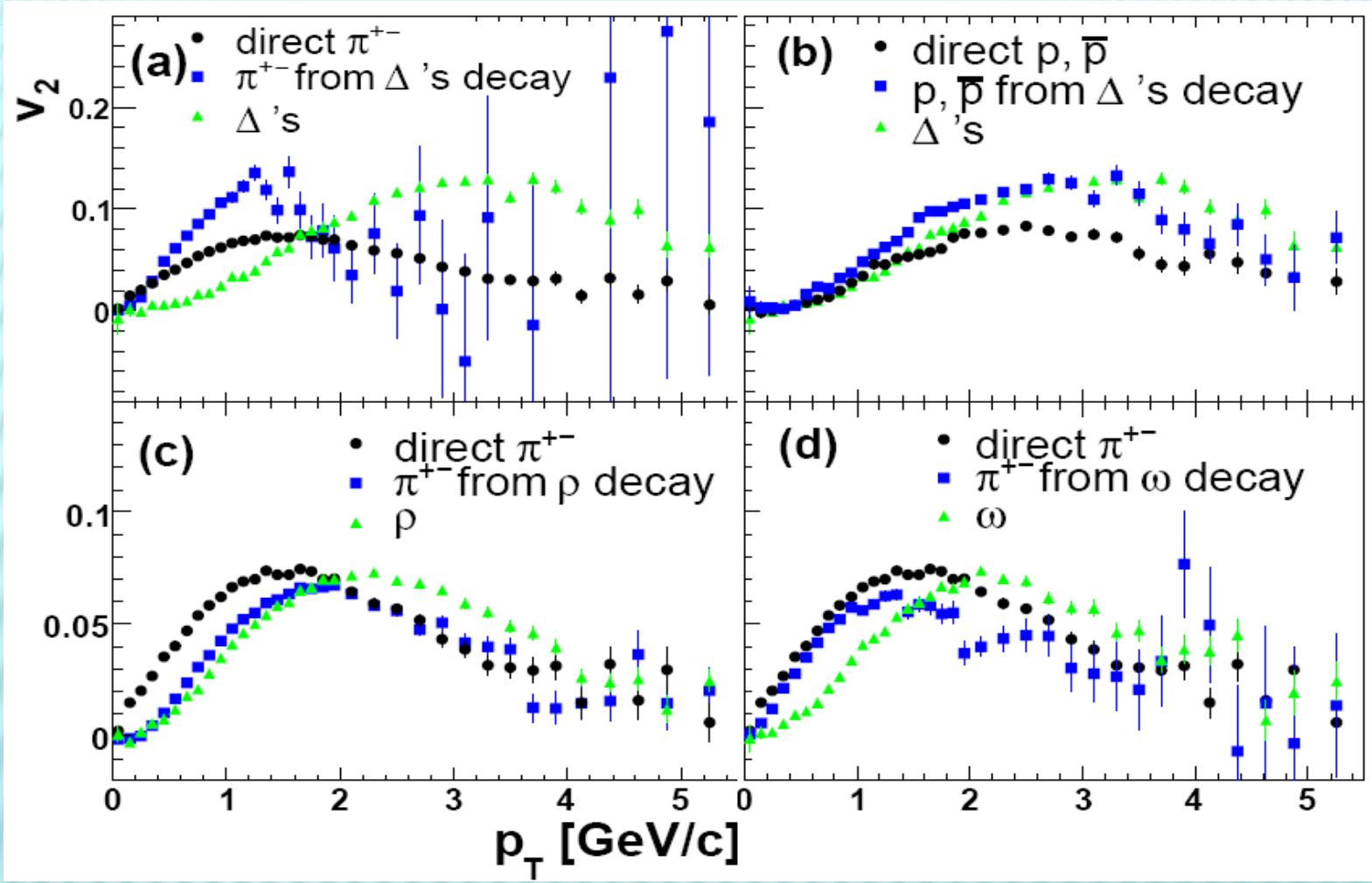
The secondary pion spectrum is much softer than proton spectrum

ELLIPTIC FLOW OF DIRECT AND SECONDARY PARTICLES AT RHIC



The heavier resonances have larger v_2 at high transverse momenta
The decay kinematics keeps this high v_2 for products of resonance decays

ELLIPTIC FLOW OF DIRECT AND SECONDARY PARTICLES AT LHC



At low transverse momenta: pions from baryon resonances enhance the flow; pions from meson resonances reduce it

V . PARAMETERS OF THE MODEL

Methods for v_2 calculation

(1) Event plane method

$$v_2^{obs} \{EP\} = \langle \cos 2(\varphi_i - \Psi_2) \rangle$$

Ψ_2 is the calculated reaction plane angle: $\tan n \psi_n = \frac{\sum_i \omega_i \sin n \varphi_i}{\sum_i \omega_i \cos n \varphi_i}$, $n \geq 1$, $0 \leq \psi_n < 2\pi/n$

$$v_2 \{EP\} = \frac{v_2^{obs} \{EP\}}{R} = \frac{v_2^{obs} \{EP\}}{\langle \cos 2(\Psi_2 - \Psi_R) \rangle}$$

(2) Two particle correlation method

$$v_2 \{2\} = \sqrt{\langle \cos 2(\varphi_i - \varphi_j) \rangle}$$

(3) Lee-Yang zero method

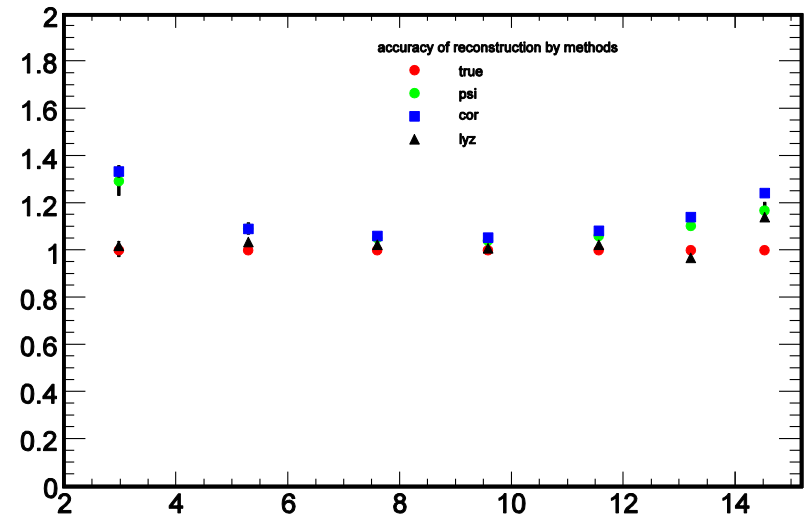
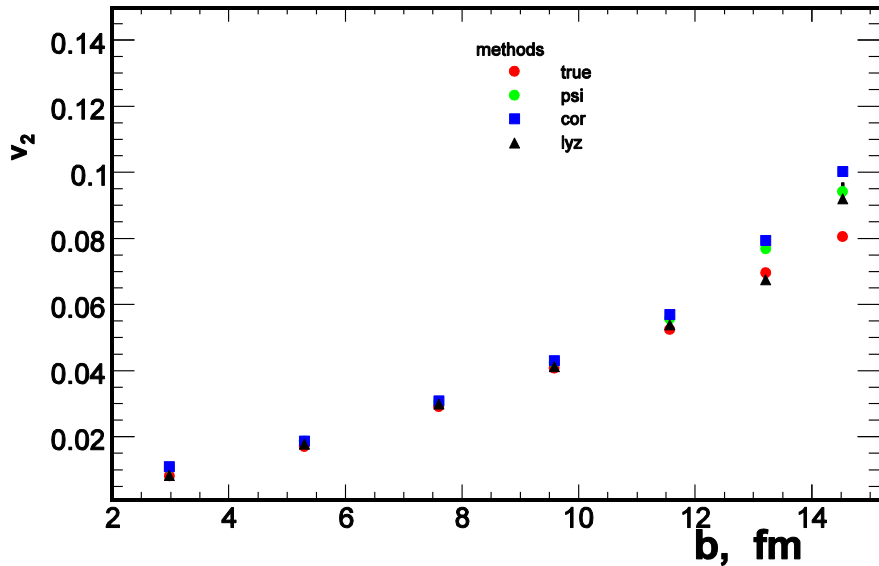
$$G(ir) = \langle e^{irQ} \rangle, Q = \sum \cos(2\varphi)$$

Integral v_2 is connected with the first minimum r_0 of the module of the $G(ir)$:

$$v_2 = \frac{j_0}{Nr_0}$$

Differential flow is calculated by the formula:
$$\frac{v_2(p_T)}{Nv_2} = \text{Re} \left(\frac{\langle \cos(2\varphi) e^{ir_0 Q} \rangle}{\langle Q e^{ir_0 Q} \rangle} \right)$$

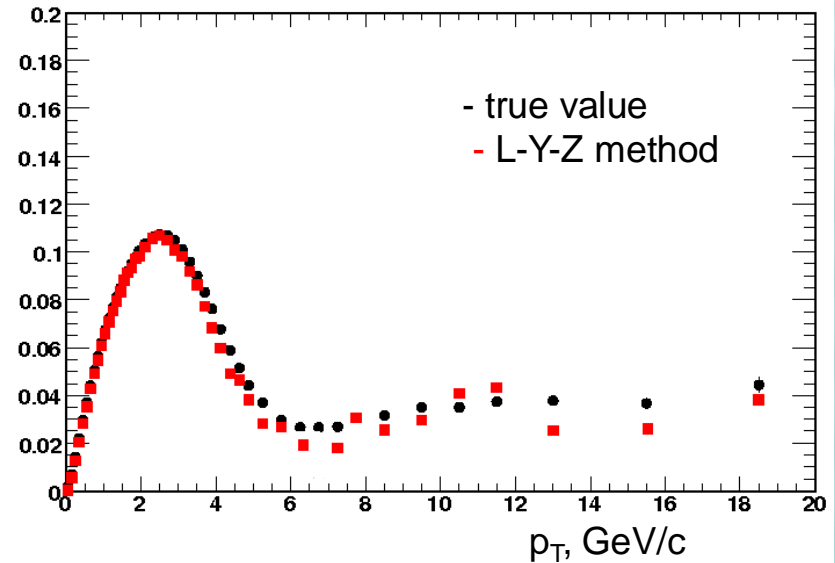
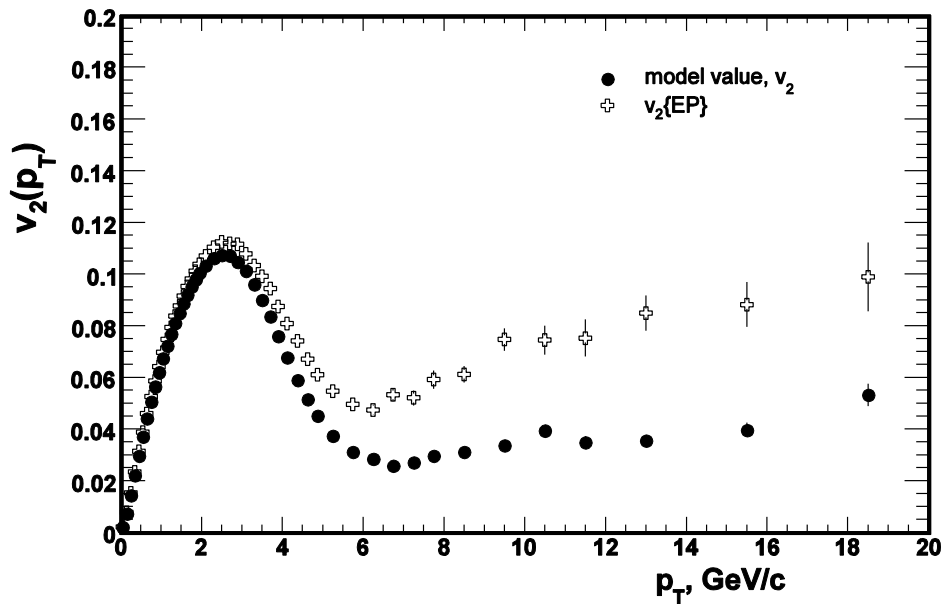
RECONSTRUCTION OF INTEGRAL VALUE OF V_2 BY THE METHODS



The better reconstruction is achieved in midcentral collision for the methods, while Lee-Yang zero method tends to reconstruct true value at more central and more peripheral collision.

Comparison of Event Plane and Lee-Yang zeroes methods ($c=30\%$)

EventPlane method



Lee-Yang zeroes Method

Event Plane method overestimates v_2 at high p_t due to non-flow correlation (mostly because of jets).