

Amplitude for production

$$\gamma p \rightarrow \gamma^* p$$

Cross section for

$$\gamma p \rightarrow l^+ l^- p$$

Exclusive photoproduction of lepton pairs at high energy

Gabriela Ślipek

Institute of Nuclear Physics PAN, Cracow

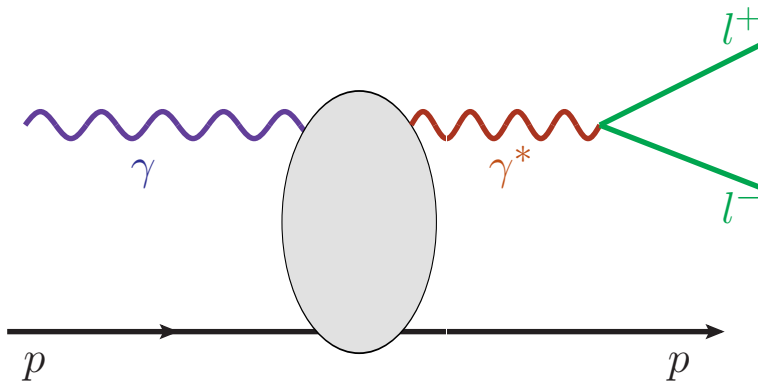
Excited QCD 2010,1-5 February, Slovakia

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in collaboration with W.Schäfer & A.Szczurek

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- $VM(\text{vector meson}) \rightarrow e^+ e^-$ decays is not the only source of leptons

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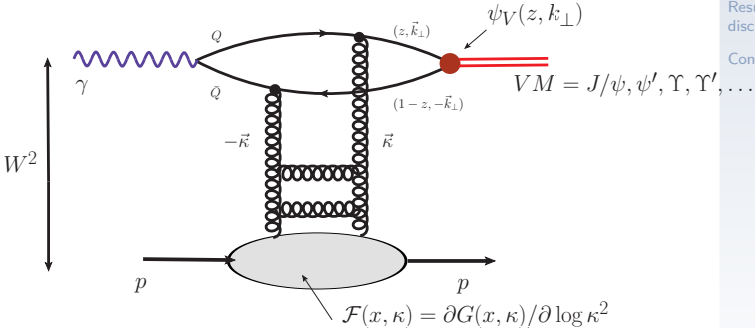
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- $VM(\text{vector meson}) \rightarrow e^+ e^-$ decays is not the only source of leptons



- formalism of VM photoproduction can be found



I.P. Ivanov, N.N. Nikolaev, A.A. Savin Phys. Part. Nucl. **37**:1-85,(2006)



A. Rybarska, W. Schäfer and A. Szczurek, Phys. Lett. B **668** (2008) 126.

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- very much the same formalism applies to $\gamma p \rightarrow l^+l^-p$ via timelike Compton scattering
- The QED process $pp \rightarrow pl^+l^-p$ is important for measuring the luminosity at LHC. It is therefore important to estimate non-QED contributions to exclusive l^+l^- production.

Hadroproduction

Exclusive
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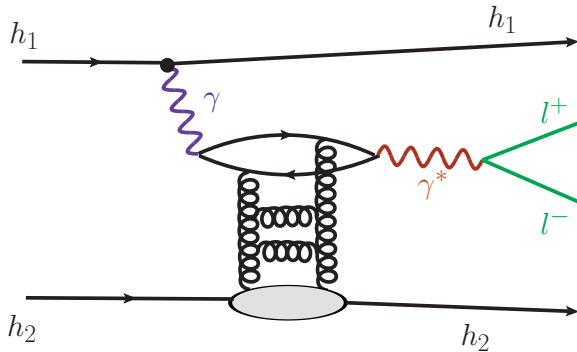
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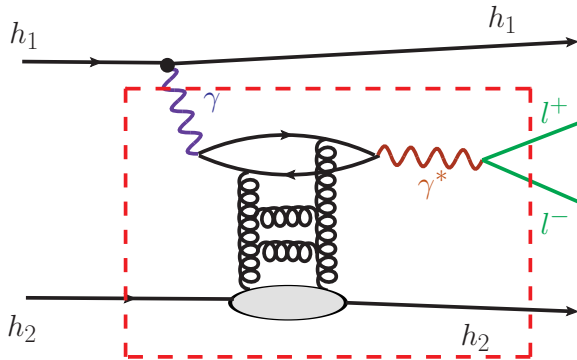
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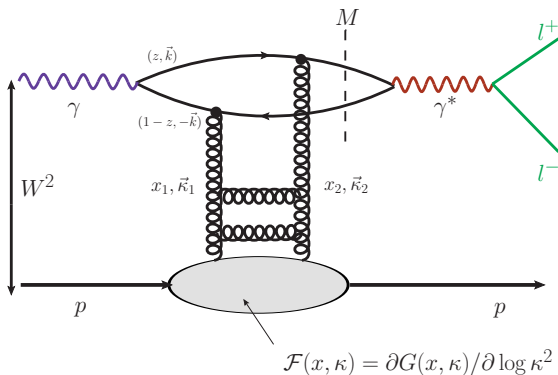
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 A. Cisek, W. Schafer, and A. Szczurek Phys. Rev. **D80** (2009) 074013.

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$$\mathcal{M} = \sum_f \mathcal{M}_f$$

$$\mathcal{M}_f = W^2 c_f 2 g_{em}^2 \int_0^1 \frac{dz}{z(1-z)} \int_0^\infty \pi dk^2 \frac{\mathcal{A}_f(z, k^2)}{\left[\frac{k^2 + m_f^2}{z(1-z)} - q^2 - i\epsilon \right]}$$

where

- $\mathcal{A}_f(z, k^2)$ - glue convolution \rightarrow see next page

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where

- $\mathcal{A}_f(z, k^2)$ - glue convolution \rightarrow see next page
- z -longitudinal momentum fraction of quark in the photon

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- z -longitudinal momentum fraction of quark in the photon
- k -transverse momentum of the quark

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where

- $\mathcal{A}_f(z, k^2)$ - glue convolution \rightarrow see next page
- z -longitudinal momentum fraction of quark in the photon
- k -transverse momentum of the quark
- $M^2 = \frac{k^2 + m_f^2}{z(1-z)}$ -invariant mass of $q\bar{q}$ pair

\mathcal{A}_f - function

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$$\mathcal{A}_f(z, k^2) = \pi \int_0^\infty \frac{\pi d\kappa^2}{\kappa^4} \alpha_s(q^2) \mathcal{F}(x, \kappa^2)$$

$$\left[A_{0f}(z, k^2) W_{0f}(k^2, \kappa^2) + A_{1f}(z, k^2) W_{1f}(k^2, \kappa^2) \right]$$

$$A_{0f}(z, k^2) = m_f^2$$

$$A_{1f}(z, k^2) = [z^2 + (1-z)^2] \frac{k^2}{k^2 + m_f^2}$$

$$W_{0f}(k^2, \kappa^2) = \frac{1}{k^2 + m_f^2} - \frac{1}{\sqrt{(k^2 - m_f^2 - \kappa^2)^2 + 4m_f^2 k^2}}$$

$$W_{1f}(k^2, \kappa^2) = 1 - \frac{k^2 + m_f^2}{2k^2} \left(1 + \frac{k^2 - m_f^2 - \kappa^2}{\sqrt{(k^2 - m_f^2 - \kappa^2)^2 + 4m_f^2 k^2}} \right)$$

Spectral density

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$$M^2 = \frac{k^2 + m_f^2}{z(1-z)}$$

$$\mathcal{M}_f = W^2 c_f^2 \cdot 2\pi \int_{4m_f^2}^{\infty} \frac{a_f(W^2, M^2)}{M^2 - q^2 - i\epsilon} dM^2$$

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$$\mathcal{M}_f = W^2 c_f 2 \cdot 2\pi \int_{4m_f^2}^{\infty} \frac{a_f(W^2, M^2)}{M^2 - q^2 - i\epsilon} dM^2$$

spectral density

$$a_f(W^2, M^2) = \int_0^{\frac{1}{4}M^2 - m_f^2} \frac{dk^2}{J_f} \mathcal{A}_f(M^2, k^2, W^2)$$

Spectral density

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$$\mathcal{M}_f = W^2 c_f 2 \cdot 2\pi \int_{4m_f^2}^{\infty} \frac{a_f(W^2, M^2)}{M^2 - q^2 - i\epsilon} dM^2$$

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$$J_f = \sqrt{1 - 4 \left(\frac{k^2 + m_f^2}{M^2} \right)}$$

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cross section

$$\frac{d\sigma}{dq^2}(\gamma p \rightarrow l^+ l^- p) = \frac{\alpha_{em}}{3\pi q^2} \sigma(\gamma p \rightarrow \gamma^* p)$$

Cross section for $\gamma p \rightarrow l^+ l^- p$

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cross section

$$\frac{d\sigma}{dq^2}(\gamma p \rightarrow l^+ l^- p) = \frac{\alpha_{em}}{3\pi q^2} \sigma(\gamma p \rightarrow \gamma^* p)$$

$$\sigma(\gamma p \rightarrow \gamma^* p) = \frac{(\Re e \frac{\mathcal{M}}{W^2})^2 + (\Im m \frac{\mathcal{M}}{W^2})^2}{16\pi \mathcal{B}}$$

Cross section for $\gamma p \rightarrow l^+ l^- p$

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cross section

$$\frac{d\sigma}{dq^2}(\gamma p \rightarrow l^+ l^- p) = \frac{\alpha_{em}}{3\pi q^2} \sigma(\gamma p \rightarrow \gamma^* p)$$

$$\sigma(\gamma p \rightarrow \gamma^* p) = \frac{(\Re \frac{\mathcal{M}}{W^2})^2 + (\Im \frac{\mathcal{M}}{W^2})^2}{16\pi \mathcal{B}}$$

- $\mathcal{M} = \sum_f \mathcal{M}_f$
- $\mathcal{B} = 4\text{GeV}^{-2}$ - slope parameter
- q^2 - invariant mass of $l^+ l^-$

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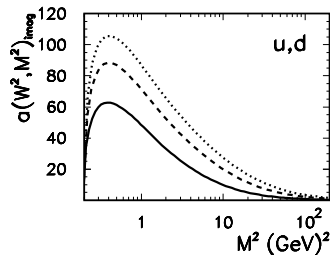
Cross section for

$\gamma p \rightarrow l^+ l^- p$

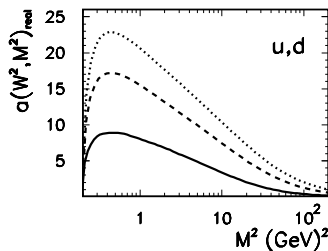
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imag part:



real part:



— 100 GeV - - - 500 GeV 1000 GeV

$$m_u, m_d = 0.22 \text{ GeV}$$

Spectral density for quark s

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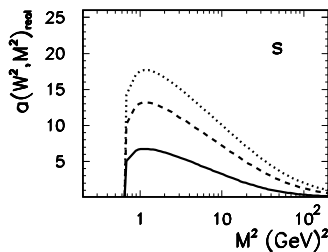
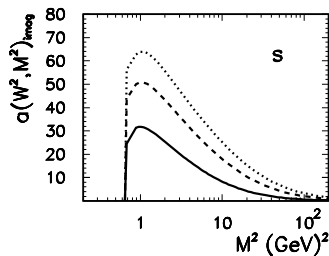
$\gamma p \rightarrow l^+ l^- p$

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imag part:

real part:



— 100 GeV

- - - 500 GeV

..... 1000 GeV

$$m_s = 0.37 \text{ GeV}$$

Spectral density for quark c

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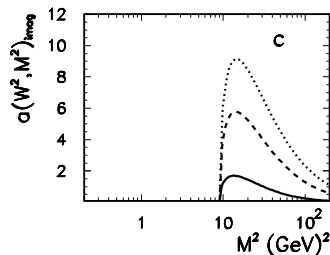
Cross section for

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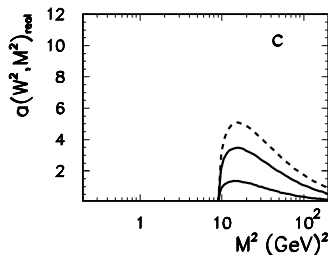
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imag part:



real part:



— 100 GeV

- - - 500 GeV

..... 1000 GeV

$$m_c = 1.5 \text{ GeV}$$

$d\sigma/dq^2$ as a function of γp cm energy

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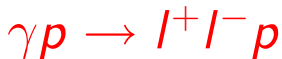
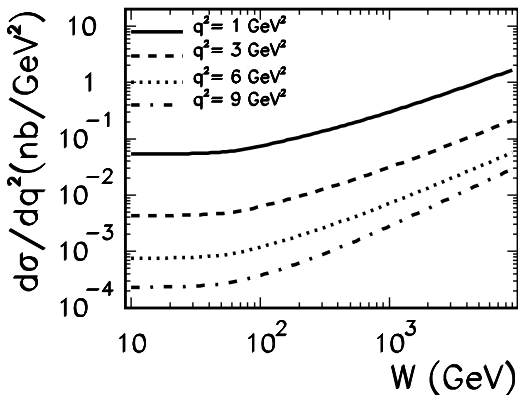
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Dependence on dilepton invariant mass

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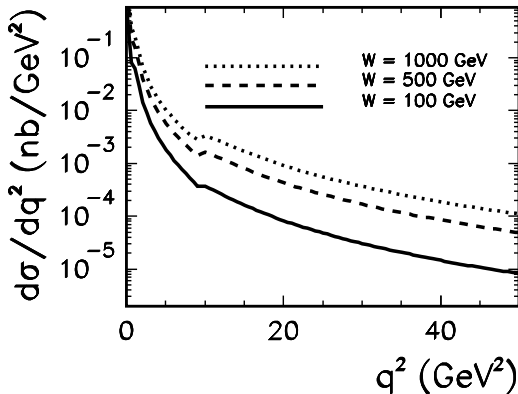
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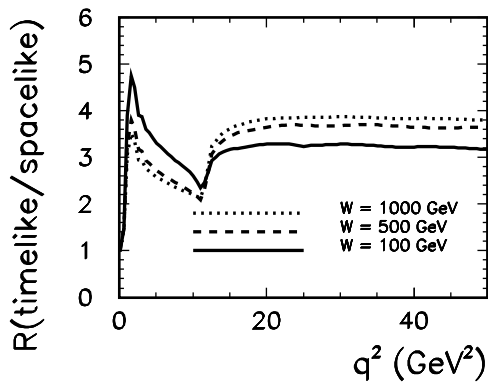


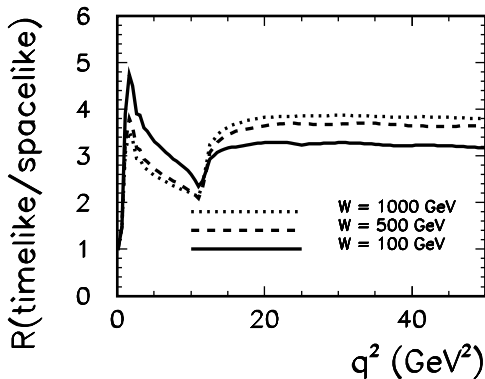
- q^2 - invariant mass of l^+l^-

Amplitude for production

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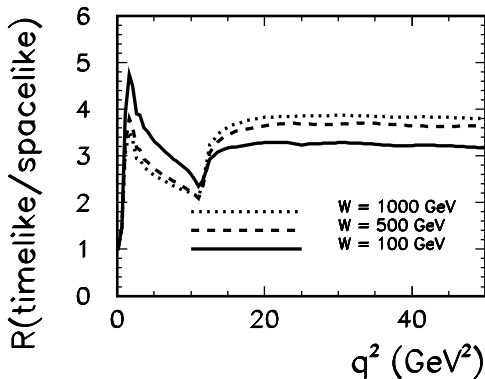
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timelike

$$a(M^2) = \int_{4m_f^2}^{\infty} \frac{a(M^2)}{M^2 - q^2 - i\epsilon} dM^2$$



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timelike

$$a(M^2) = \int_{4m_f^2}^{\infty} \frac{a(M^2)}{M^2 - q^2 - i\epsilon} dM^2$$

spacelike

$$a(M^2) = \int_{4m_f^2}^{\infty} \frac{a(M^2)}{M^2 + q^2} dM^2$$

- Amplitude for the $\gamma p \rightarrow \gamma^* p$ has been derived in the k_t - factorization approach.
- I have presented results for exclusive photoproduction of lepton pairs (energy and $q (= M_{l+l-})$ dependences).
- I have done calculations assuming time-like photons instead of space-like photons as wrongly done in the literature. This leads to an enhancement compared to earlier estimates.
- Future: go to hadroproduction and see what is background to the QED process (luminosity monitor).

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