

*Partons and mesons in a strongly coupled plasma  
from AdS/CFT*

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# Introduction

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  - “strongly coupled quark–gluon plasma (sQGP)”

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- How to perform calculations at strong coupling ?
  - **lattice QCD**: first–principle method ... but only for static quantities (thermodynamics, screening masses)
  - $\mathcal{N} = 4$  **supersymmetric Yang–Mills (SYM) theory**: string theory techniques via the AdS/CFT duality

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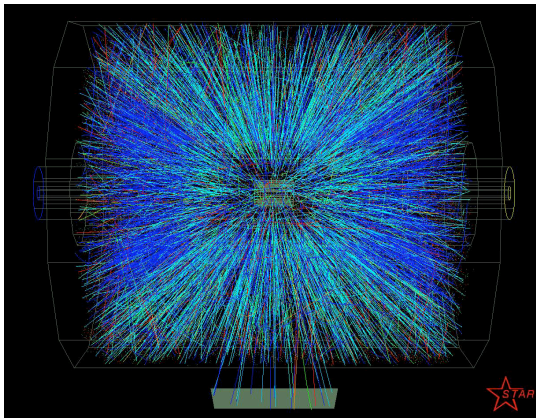
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- How to perform calculations at strong coupling ?
  - **lattice QCD**: first–principle method ... but only for static quantities (thermodynamics, screening masses)
  - $\mathcal{N} = 4$  **supersymmetric Yang–Mills (SYM) theory**: string theory techniques via the AdS/CFT duality
- Can we trust AdS/CFT in relation with QCD ?
  - for **qualitative** understanding at most (still interesting !)
  - for **specific** observables & phenomena

- 1 Hard probes at RHIC
- 2 AdS/CFT
- 3 DIS in AdS/CFT
- 4 Phenomenology
- 5 Conclusions

# Motivations: Heavy Ion Collisions @ RHIC & LHC

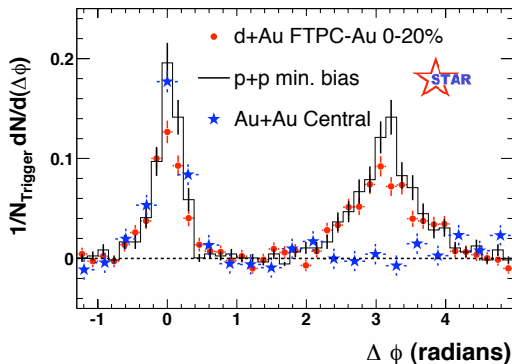
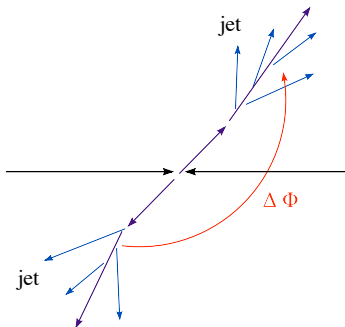


# Hadron production at RHIC



- $\sim 3000$  hadrons in the final state
- Particle correlations are essential to disentangle phenomena

# Jets in proton–proton collisions



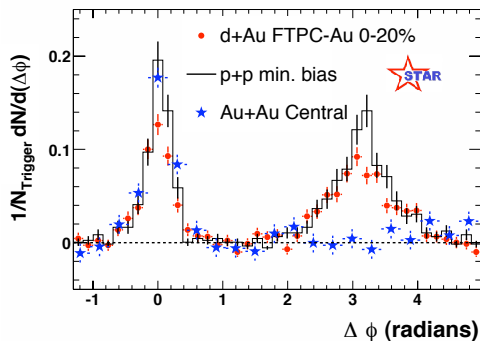
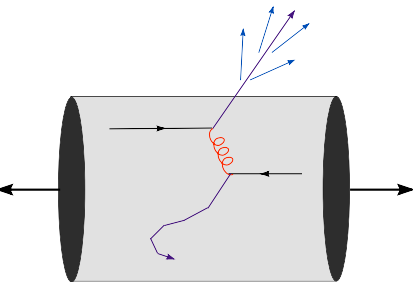
[Nucl.Phys.A783:249-260,2007]

- Azimuthal correlations between the produced jets:

p+p or d+Au : a peak at  $\Delta\Phi = 180^\circ$



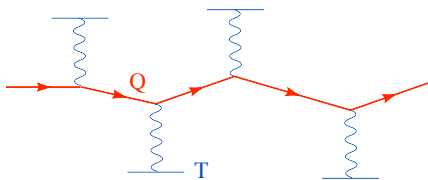
# Jet ‘quenching’ in nucleus–nucleus collisions



- The “away–side” jet has disappeared (for relatively hard momenta:  $Q \sim 2 \div 10 \text{ GeV} \gg T \simeq 400 \text{ MeV}$ )  
 $\implies$  strong interactions in the medium
- Perturbation theory seems unable to explain this suppression

# Jet quenching in perturbative QCD

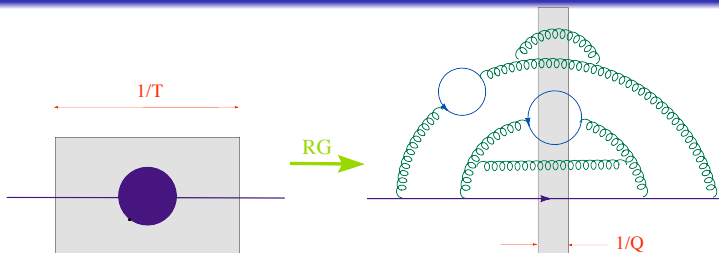
- Medium rescattering  $\implies$  transverse momentum broadening



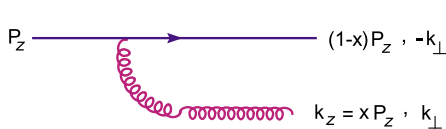
$$\frac{d\langle k_{\perp}^2 \rangle}{dt} \equiv \hat{q} \simeq \alpha_s N_c \mathcal{G}(x, Q^2)$$

- $\mathcal{G}(x, Q^2)$  : gluon distribution per unit volume in the medium on the resolution scales  $Q^2 \sim \langle k_{\perp}^2 \rangle$  and  $1/x \sim k_z T/Q^2$
- If “medium” = QCD plasma at temperature  $T$ :  
we expect quarks and gluons with momenta  $\sim T$
- Jet quenching requires **parton evolution** from  $T$  up to  $Q \gg T$

# Parton evolution



- A 'quasiparticle' (quark or gluon) on the scale  $T$  may reveal itself as highly composite on the harder scale  $Q \gg T$
- Weak coupling: Bremsstrahlung



$$d\mathcal{P}_{\text{Brem}} \sim \alpha_s N_c \frac{d^2 k_{\perp}}{k_{\perp}^2} \frac{dx}{x}$$

soft ( $x \ll 1$ ) & collinear ( $k_{\perp} \rightarrow 0$ )

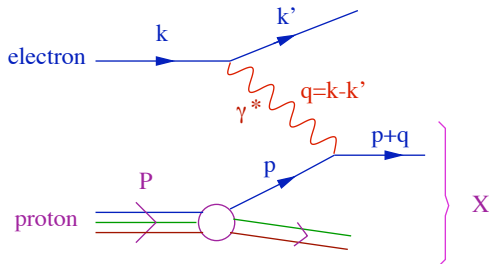
# Deep inelastic scattering

- How to study parton evolution at **strong coupling** ?
- **DIS**: a non-perturbative & observable definition of 'partons'

- 2 independent variables:

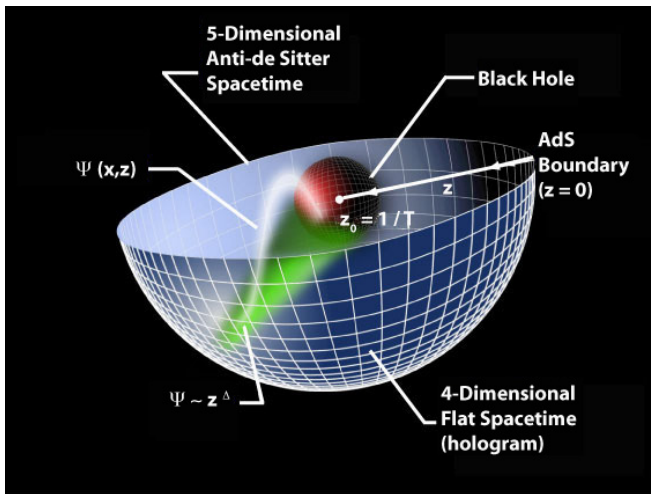
$$Q^2 \equiv -q^\mu q_\mu \geq 0$$

$$x \equiv \frac{Q^2}{2P \cdot q}$$



- **Physical picture**:  $\gamma^*$  absorbed by a quark excitation with
  - transverse size  $\Delta x_\perp \sim 1/Q$
  - and longitudinal momentum  $p_z = xP$
- DIS 'structure function'  $F_2(x, Q^2)$ : quark distribution

# Parton evolution at strong coupling



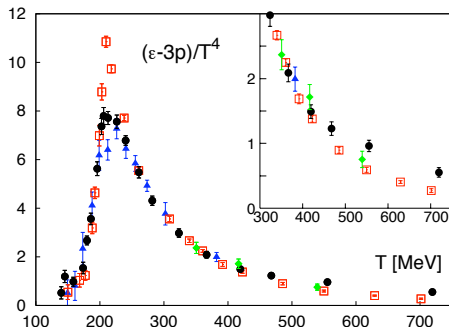
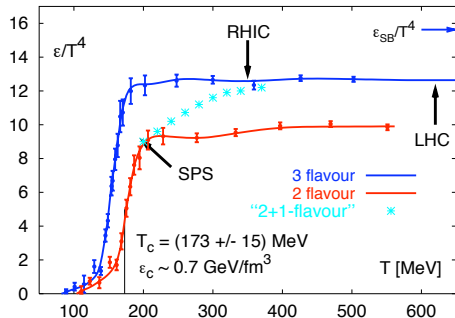
8-2007  
8685A14

*(courtesy of Stan Brodsky and MM)*

# The AdS/CFT correspondance (Maldacena, 1997)

- A 'duality' (equivalence) between two very different theories
- Gauge theory in  $D = 4$ :  $\mathcal{N} = 4$  Supersymmetric Yang–Mills
  - color gauge group  $SU(N_c)$
  - conformal invariance, fixed coupling  $g$ , no confinement
  - strong 't Hooft coupling :  $\lambda \equiv g^2 N_c \gg 1$  &  $g^2 \ll 1$
- Is this a good model for QCD ??
- Perhaps better suited for studies of QCD at finite temperature
  - deconfined phase (quark–gluon plasma)
  - nearly conformal for  $T \gtrsim 2T_c$
  - coupling is relatively strong in the relevant range

# 'Trace anomaly' in lattice QCD



$$\beta(g) \frac{dp}{dg} = \langle T_{\mu}^{\mu} \rangle = \mathcal{E} - 3p$$

- $(\mathcal{E} - 3p)/\mathcal{E}_0 \lesssim 10\%$  for any  $T \gtrsim 2T_c \simeq 400$  MeV
- $g \approx 1.5 \div 2 \implies \lambda \equiv g^2 N_c \simeq 6 \div 10$

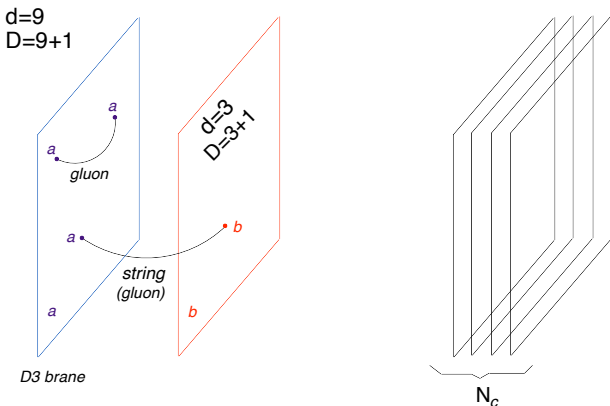
# The AdS/CFT correspondance (*Maldacena, 1997*)

- The  $\mathcal{N} = 4$  SYM gauge theory in  $D = 3 + 1$  is dual to type IIB string theory in  $D = 9 + 1$  ( $AdS_5 \times S^5$  space-time)
- The strong 't Hooft coupling regime of the gauge theory:
  - $\lambda \equiv g^2 N_c \gg 1$  &  $g^2 \ll 1$  (large  $N_c$ )
- ... corresponds to the semi-classical regime of the string theory
  - weak coupling & weak curvature
  - string theory reduces to 'supergravity' :  
classical equations of motion in a curved space-time
- The gauge theory emerges as the low-energy limit of a string theory with  $N_c$  'D3-branes'
- A 'D3-brane': a manifold with  $D=3+1$  immersed into  $D=9+1$



# The AdS/CFT correspondance

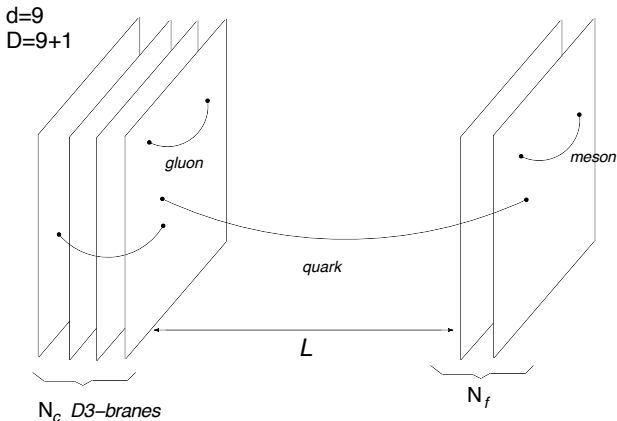
- 'gluons' = strings connecting  $N_c$  overlapping D3-branes



- Gluons are massless since the minimal string length is zero

# The AdS/CFT correspondance

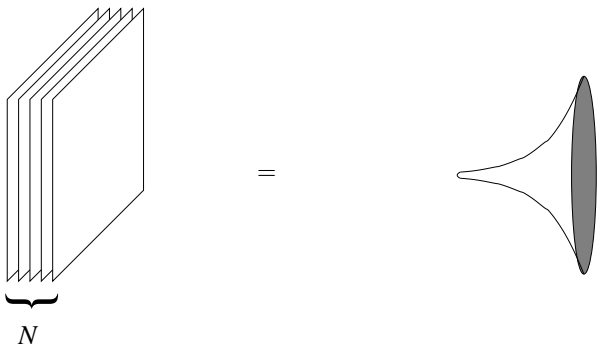
- 'quarks' = strings connecting one of the  $N_f$  'flavor' branes to one of the  $N_c$  'color' branes



- Minimal string length is  $L \implies$  quark mass  $M_q = L \times \mathcal{T}$

# The AdS/CFT correspondance

- The branes are massive  $\implies$  they deform the space-time



- $AdS_5$  : Our Minkowski world  $\times$  a 'radial' dimension  $\chi$ 
  - $0 \leq \chi < \infty$  : 'radial', or '5th', coordinate
  - gauge theory lives at the Minkowski boundary  $\chi = 0$

# Heating AdS<sub>5</sub> : the Black Hole

- $\mathcal{N} = 4$  SYM at finite temperature  $\iff$  Black Hole in AdS<sub>5</sub>  
A Black Hole has **entropy** and **thermal (Hawking) radiation**
- Gauge theory operator  $\mathcal{O}(t, \vec{x})$  on the 'boundary' ( $\chi = 0$ ) :  
source for gravity fields  $\phi(t, \vec{x}, \chi)$  inside the 'bulk' ( $\chi > 0$ )
- Classical EOM for  $\phi(t, \vec{x}, \chi)$  into AdS<sub>5</sub>  $\implies$   
quantum correlations for  $\mathcal{O}(t, \vec{x})$  in the limit  $\lambda \rightarrow \infty$
- Absorbtion of the field  $\phi(t, \vec{x}, \chi)$  by the BH  $\implies$   
dissipative phenomena within the strongly-coupled plasma
- What is the rôle of the 5th dimension ?  
dual to the 'loop' momenta in the usual Feynman graphs

# DIS off the Black Hole (*Hatta, E.I., Mueller, 07*)

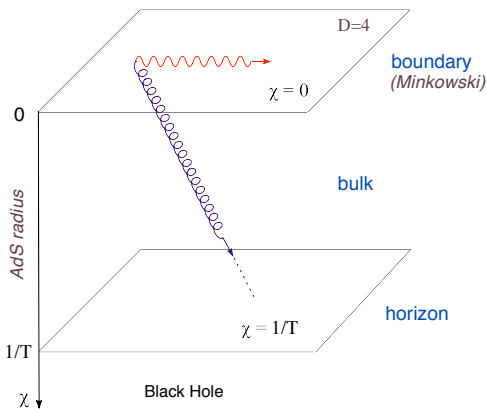
- E.m. current  $J^\mu$  in 4D  $\longleftrightarrow$  Maxwell wave  $A_\mu$  in AdS<sub>5</sub> BH
- DIS cross section  $\longleftrightarrow$  absorption of the wave by BH

- Physical world:  $\chi = 0$   
Black Hole horizon:  $\chi = 1/T$
- Maxwell equations in AdS<sub>5</sub> BH

$$\partial_m (\sqrt{-g} g^{mn} g^{pq} F_{nq}) = 0$$

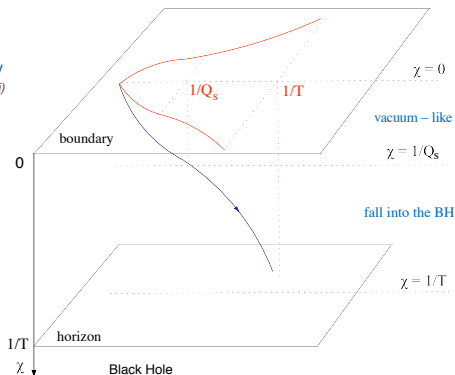
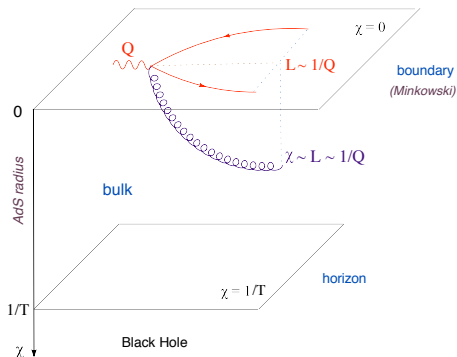
$$F_{mn} = \partial_m A_n - \partial_n A_m$$

- No explicit coupling



# Space-like photon in the plasma

- Interactions are controlled by the kinematics ( $x, Q^2$ )
- For **low energies**, the virtual photon does not 'see' the BH !



- ... while for **large enough energies**, it is **completely absorbed** !

# Saturation line

- Gravitational interactions are proportional to the energy density in the wave ( $\omega$ ) and in the plasma ( $T$ )

$$\text{DIS kinematics : } x \equiv \frac{Q^2}{2\omega T} \text{ and } Q \gg T$$

- Large  $\omega T$  is tantamount to small Bjorken's  $x$
- Critical ('saturation') value  $x_s(Q) \simeq \frac{T}{Q} \ll 1$

- $x > x_s \simeq T/Q$  :  $F_2(x, Q^2) \approx 0$  : no partons !

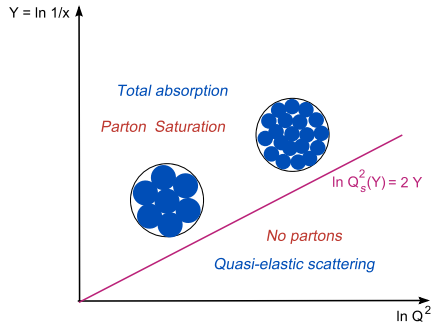
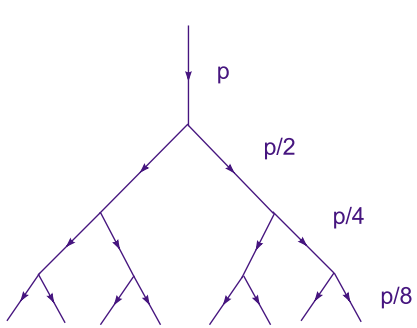
- $x < x_s \simeq T/Q$  :  $F_2(x, Q^2) \sim x N_c^2 Q^2$

$\implies$  Parton saturation with occupation numbers  $\mathcal{O}(1)$

- Where did all the partons go ??

# Physical interpretation: Parton evolution

- The energy of the plasma is carried mostly by the partons along the saturation line:  $x_s \simeq T/Q \ll 1$

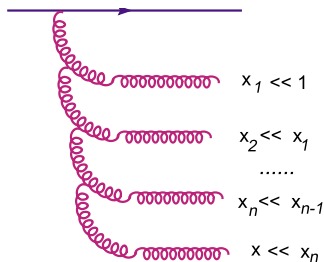


- All partons branch down to the smallest value of  $x$  consistent with energy conservation  $\implies$  no pointlike constituents



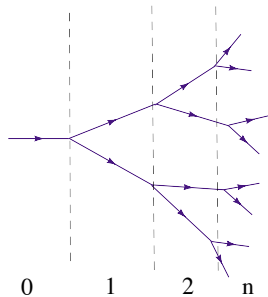
# Parton evolution: weak vs. strong coupling

- Weak coupling



- Bremsstrahlung
- Soft & collinear emissions
- Low multiplicity :  $N \propto \ln E$

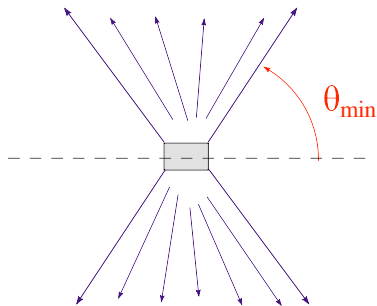
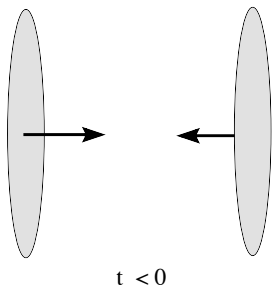
- Strong coupling



- Quasi-democratic branching :  
 $\omega_n \sim \omega_{n-1}/2$
- High multiplicity :  $N \propto E/\Lambda$

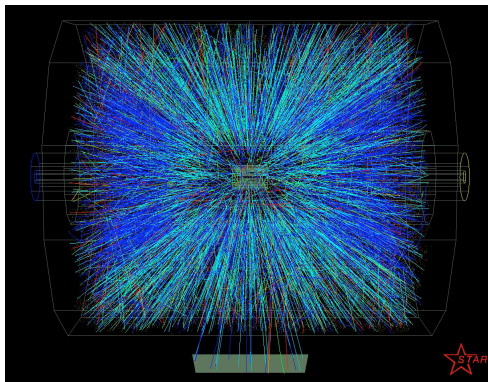
# No forward/backward jets in $AA$ collisions

- No large- $x$  partons  $\implies$  no hard ( $Q \gg \Lambda$ ) particle production at forward/backward rapidities



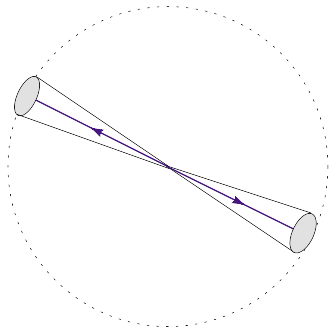
- All the energy is carried out by soft particles with  $p \sim \Lambda$

# Partons at RHIC

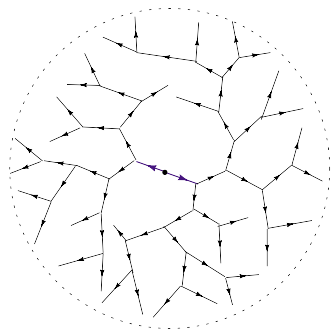


- Partons are actually ‘seen’ (liberated) in the high energy hadron–hadron collisions
  - central rapidity: small- $x$  partons
  - forward/backward rapidities: large- $x$  partons

# No jets in $e^+e^-$ annihilation !



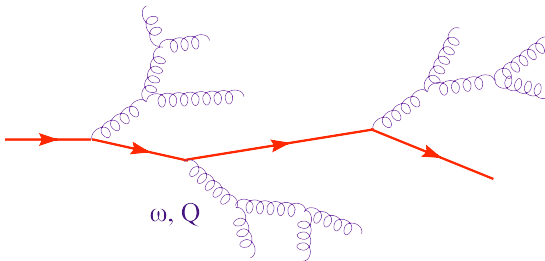
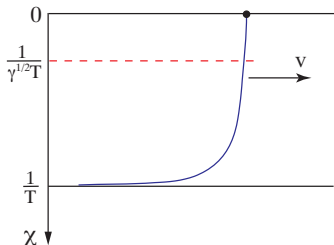
weak coupling



strong coupling

- An isotropic distribution of soft hadrons in the detector  
(similar conclusions by Hofman and Maldacena, 2008)

# Heavy Quark in a strongly-coupled plasma



- Medium-induced radiation

- virtual quanta with  $Q \lesssim Q_s$  are liberated into the plasma
- energy loss, momentum broadening
- quantum stochasticity (random emissions)

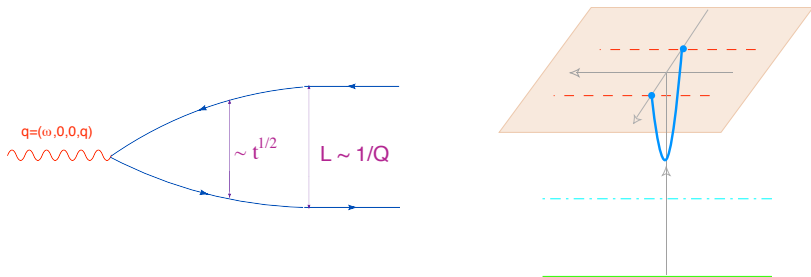
*Casalderrey-Solana and Teaney; Herzog, Karch, Kovtun, Kozcaz, and Yaffe; Gubser, 2006;*

*G. Giacold, E.I., and A. Mueller 2009*

- Different mechanism than in pQCD: radiation vs. rescattering

# Meson dissociation in the plasma

- $q\bar{q}$  fluctuation  $\approx$  a 'meson' with transverse size  $L \sim 1/Q$



- The 'meson' melts when  $L \geq L_s(v)$

$$L_s(v) \sim \frac{1}{Q_s(v)} \sim \frac{(1-v^2)^{1/4}}{T}$$

- Velocity-dependent screening length:  
the faster mesons are easier destroyed in the medium !

*Peeters et al; Liu, Rajagopal, and Wiedemann (2006)*

# AdS/CFT applications to heavy ion collisions

- A very active area of research, with some serious limitations
  - no asymptotic freedom
  - so far, no systematic way to improve
  - high-energy physics looks very different from real-life QCD
- ... but which is still interesting
  - new perspectives on old problems: QGP = Black Hole
  - the potential to solve longstanding puzzles at RHIC (early thermalization, elliptic flow, jet quenching)
  - most likely, the coupling is moderately strong
- It teaches us the unity of physics
  - quantum field theory, nuclear physics, statistical physics, gravity, hydrodynamics ...