



Tracing the QCD pressure

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Why this talk may be relevant – or not ;)

- idea relevant for many lattice-QCD calculations at $T>0$
- understand QCD where we think it is simple
- revisit first non-perturbative coefficient

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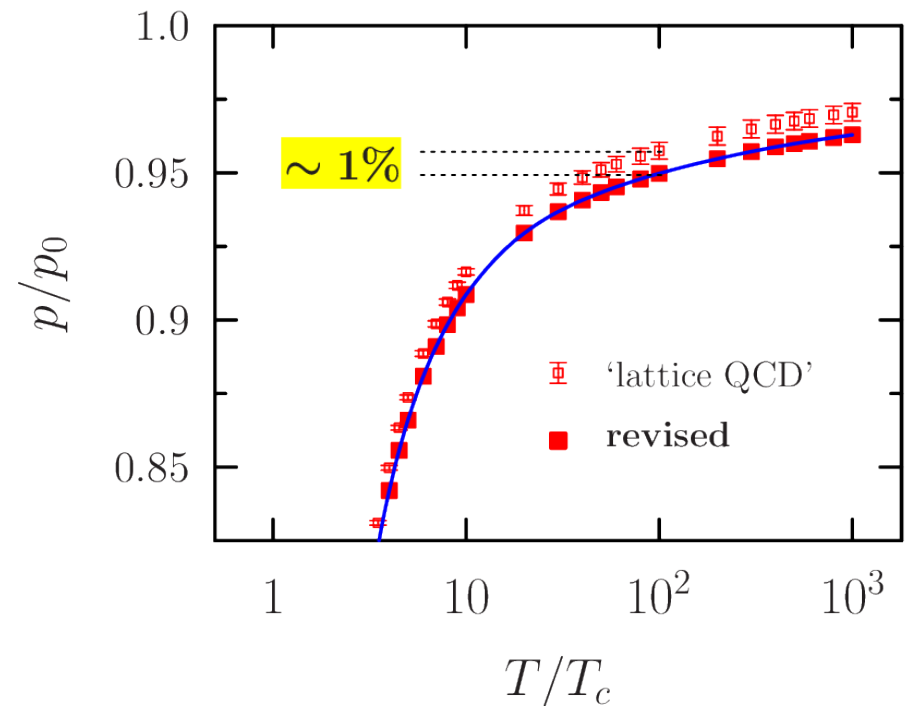
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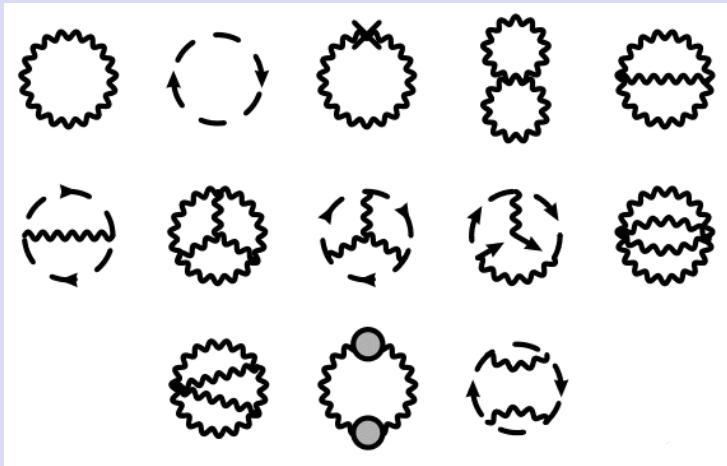
Methods

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu} + \dots \text{ where } F_a^{\mu\nu} = \partial^\mu A_a^\nu - \partial^\nu A_a^\mu - g f_{abc} A_b^\mu A_c^\nu$$

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Perturbation theory

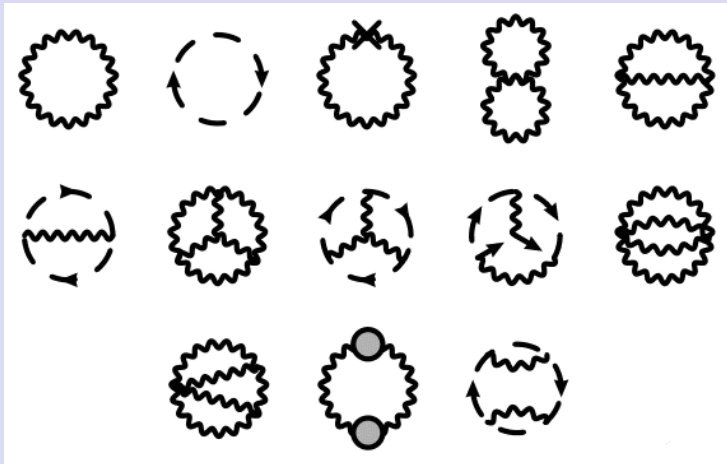


$$p = p_{\text{SB}} \left[1 + c_2 \sqrt{\alpha^2} + \dots \right]$$

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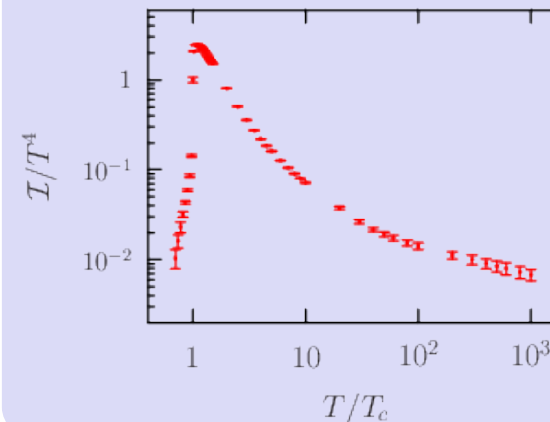
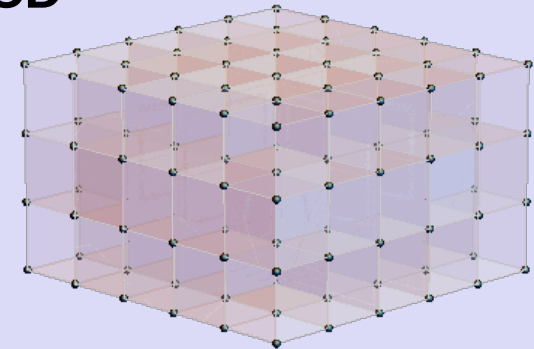
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Lattice QCD



Perturbation theory **CHALLENGES**

$$\frac{p}{p_{\text{SB}}} = 1 + c_2 \alpha^{2/2} + c_3 \alpha^{3/2} + (c_4 + \tilde{c}_4 \ln \alpha) \alpha^{4/2} + c_5 \alpha^{5/2} + (c_6 + \tilde{c}_6 \ln \alpha) \alpha^{6/2} + \dots$$



[Shuryak 1978]

⋮

[Kajantie et al 2003]

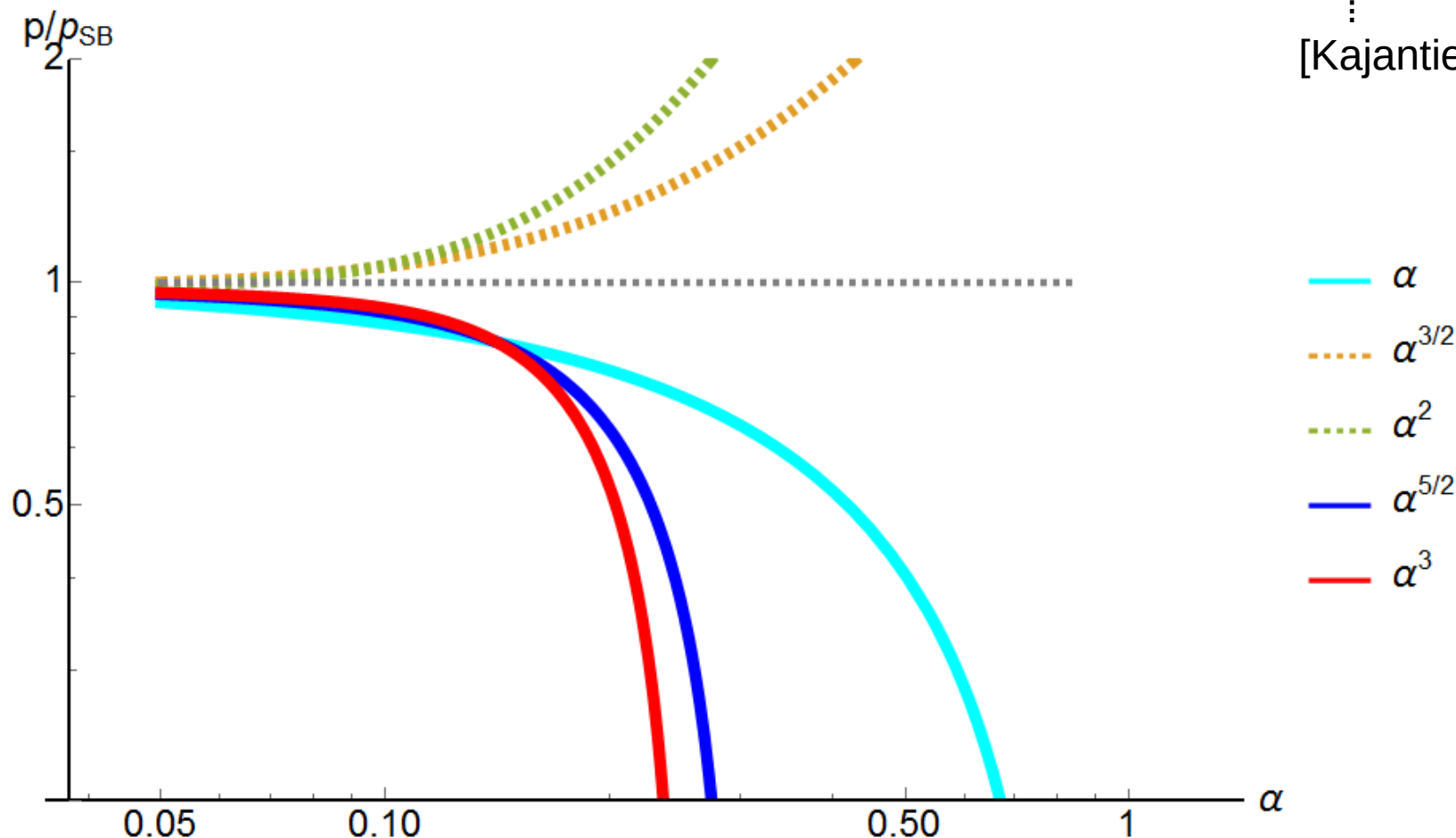
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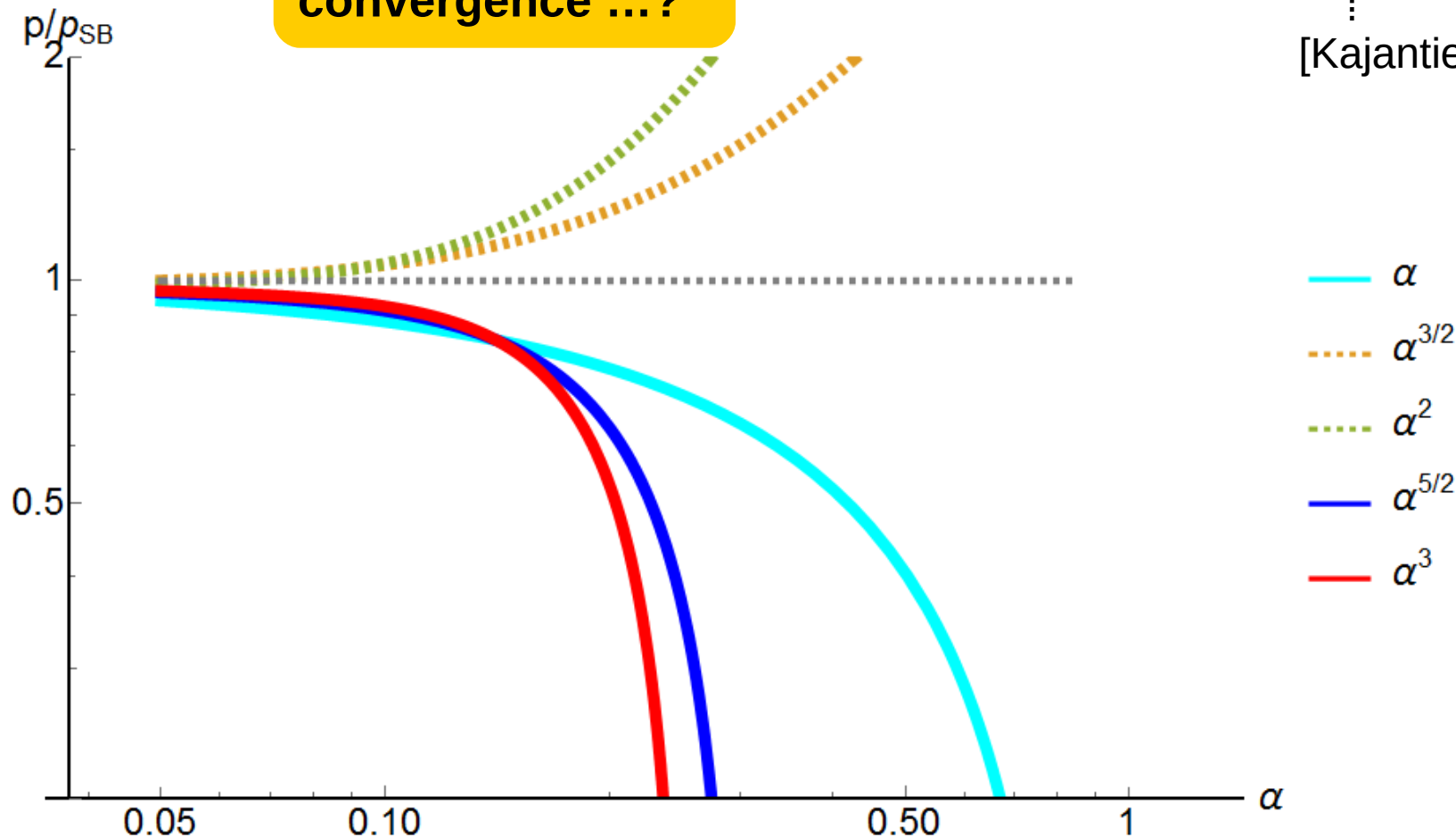
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convergence ...?



[Shuryak 1978]

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Perturbation theory “diverges”

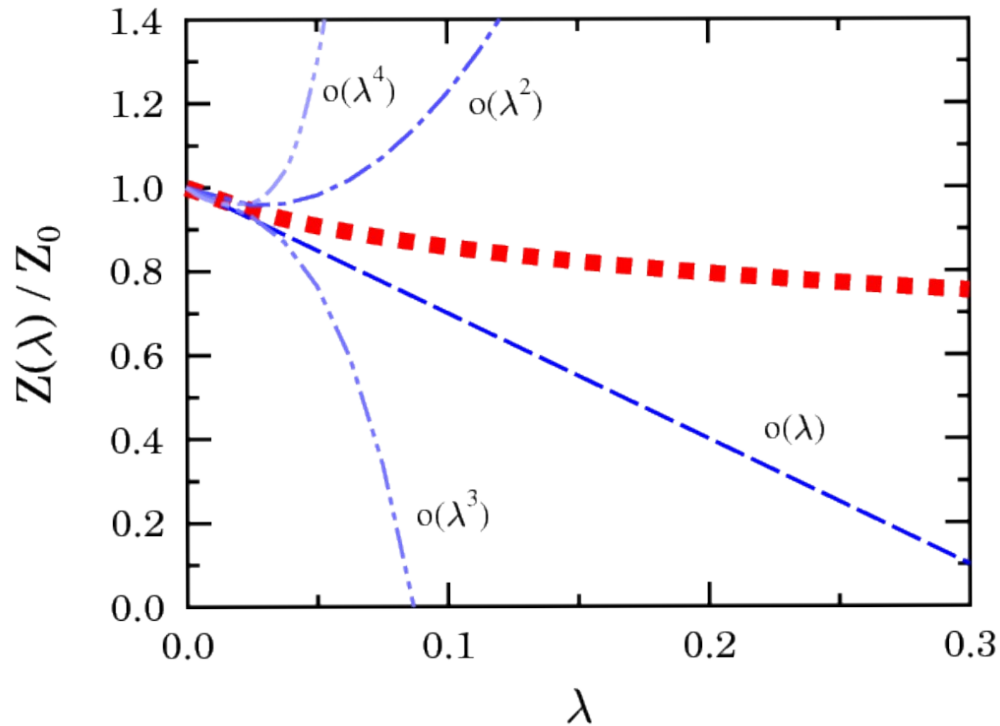
toy model: scalar “QFT” in $d=0$ dimensions

- “Lagrangian” $L = \frac{1}{2}x^2 + \lambda x^4$ → “partition fnc” $Z(\lambda) = \int dx \exp(-L(\lambda))$
- perturbative expansion $Z(\lambda)/Z_0 = 1 - 3\lambda + \frac{1}{2} 105\lambda^2 \mp \dots$

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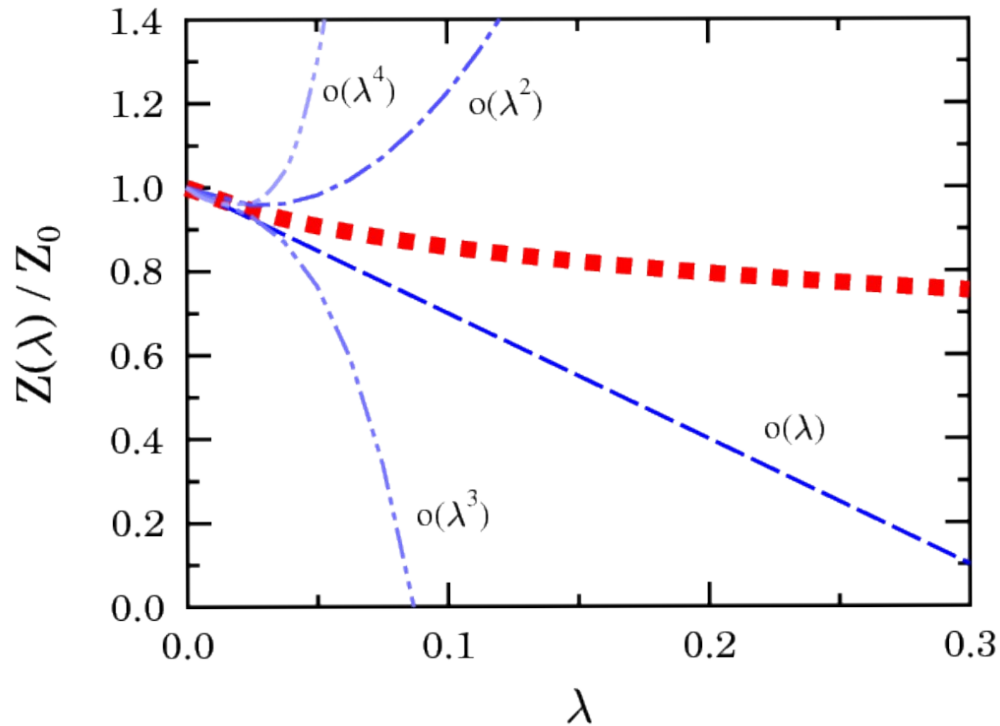


lower order better for larger coupling

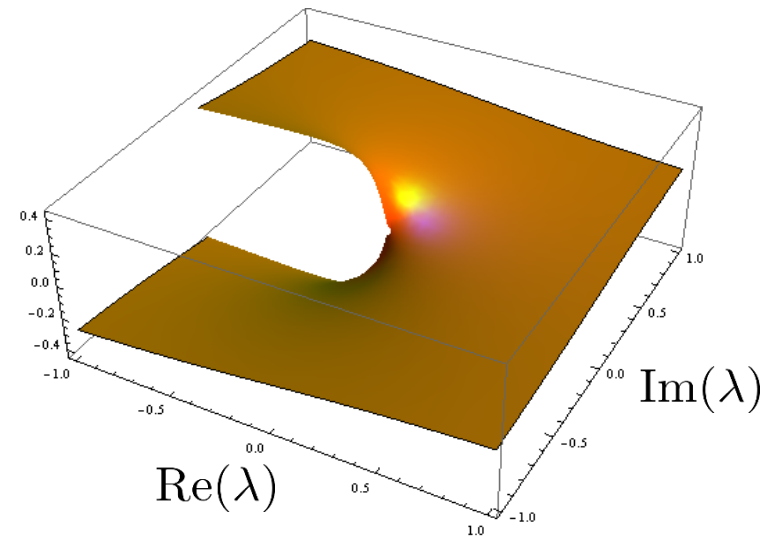
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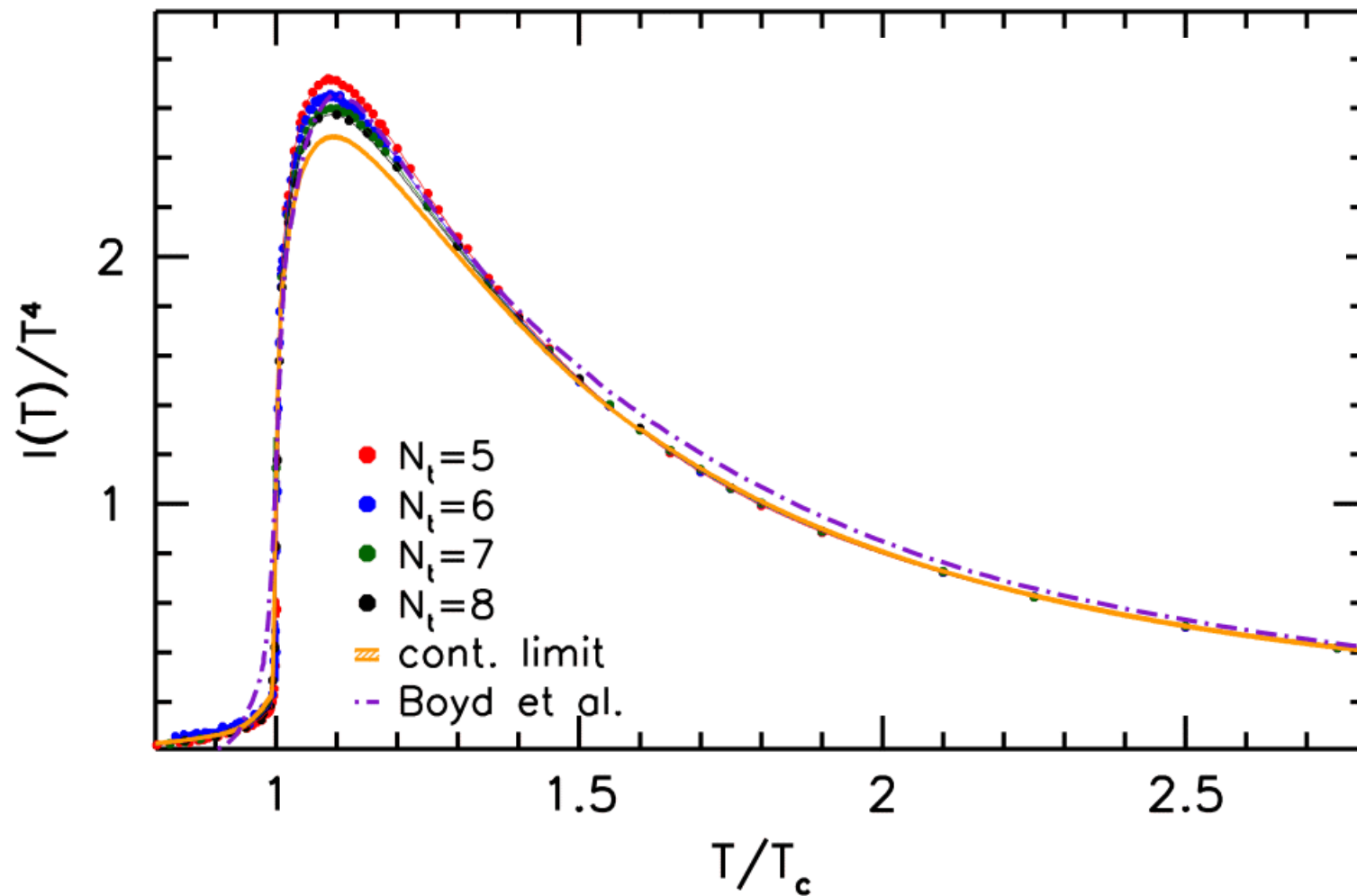


cut λ plane: convergence radius = 0



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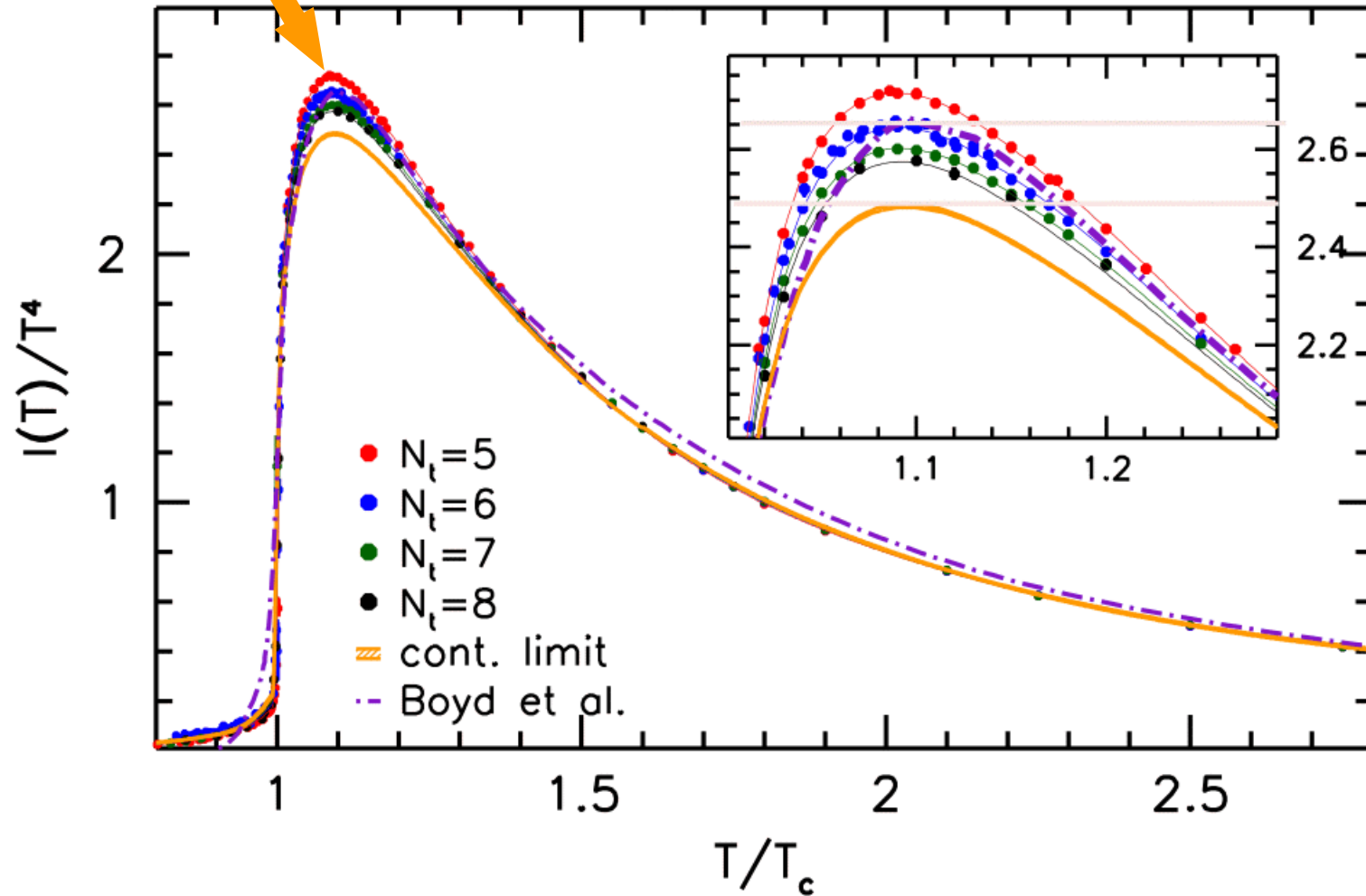
Lattice QCD **CHALLENGES**



[Borsanyi et al, 2012]

Lattice QCD **CHALLENGES**

finite-size artefacts in particular around T_c : **correlations**



[Borsanyi et al, 2012]

Lattice QCD

integral method: **pressure** from **interaction measure** $\mathcal{I} = e - 3p$

$$\frac{p(T)}{T^4} = \sigma + \int_{T_0}^T \frac{dT'}{T'} \frac{\mathcal{I}(T')}{T'^4} \quad \text{where} \quad \sigma = \frac{p(T_0)}{T_0^4}$$

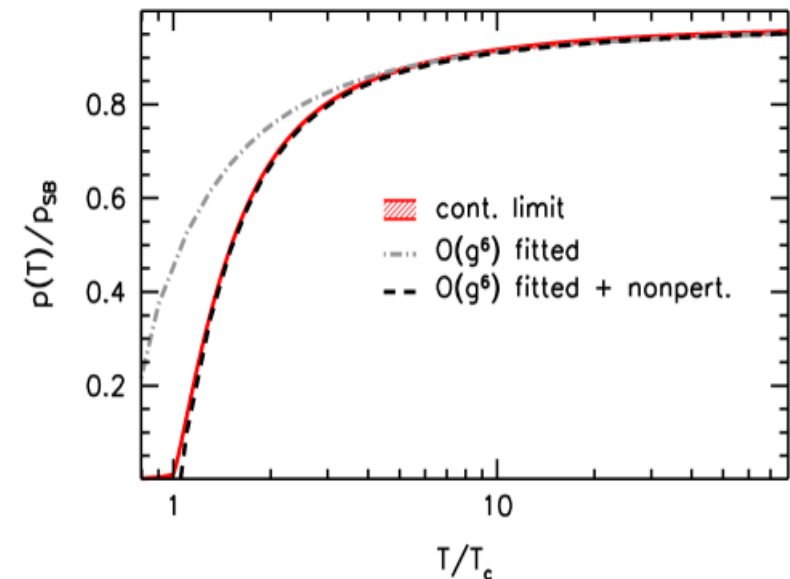
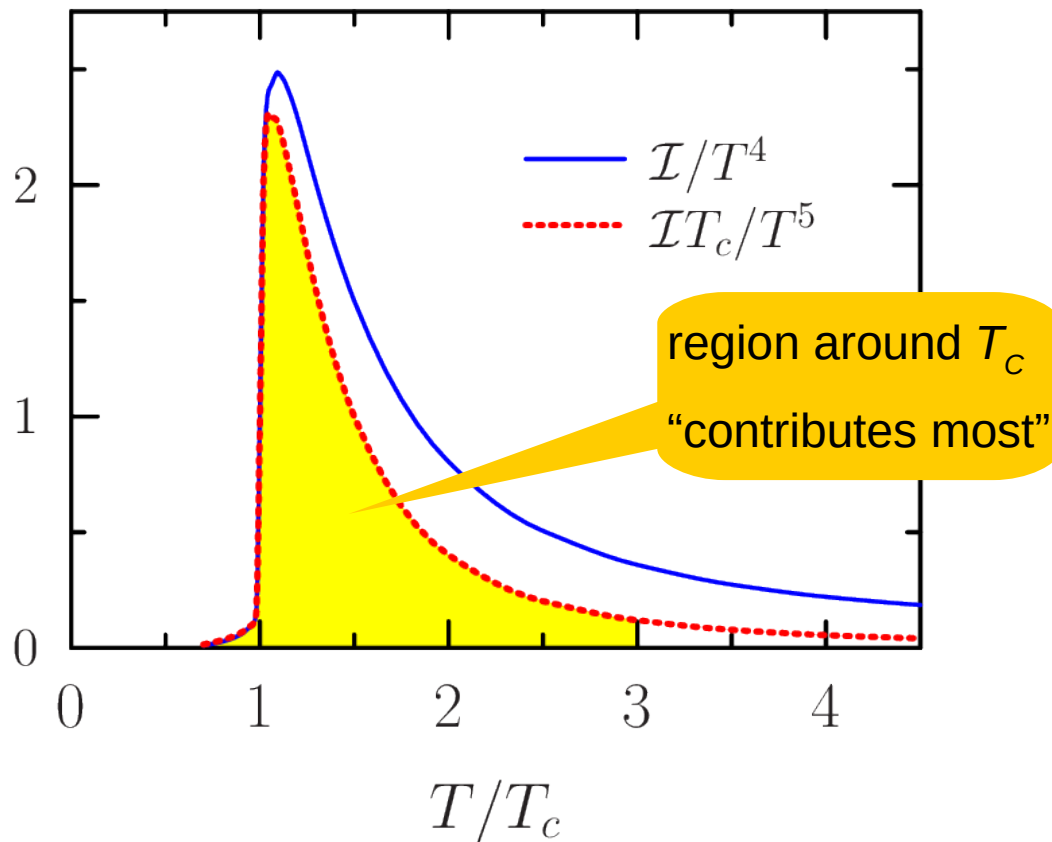
where $e = sT - p$
with $s = \partial p / \partial T$

Lattice QCD

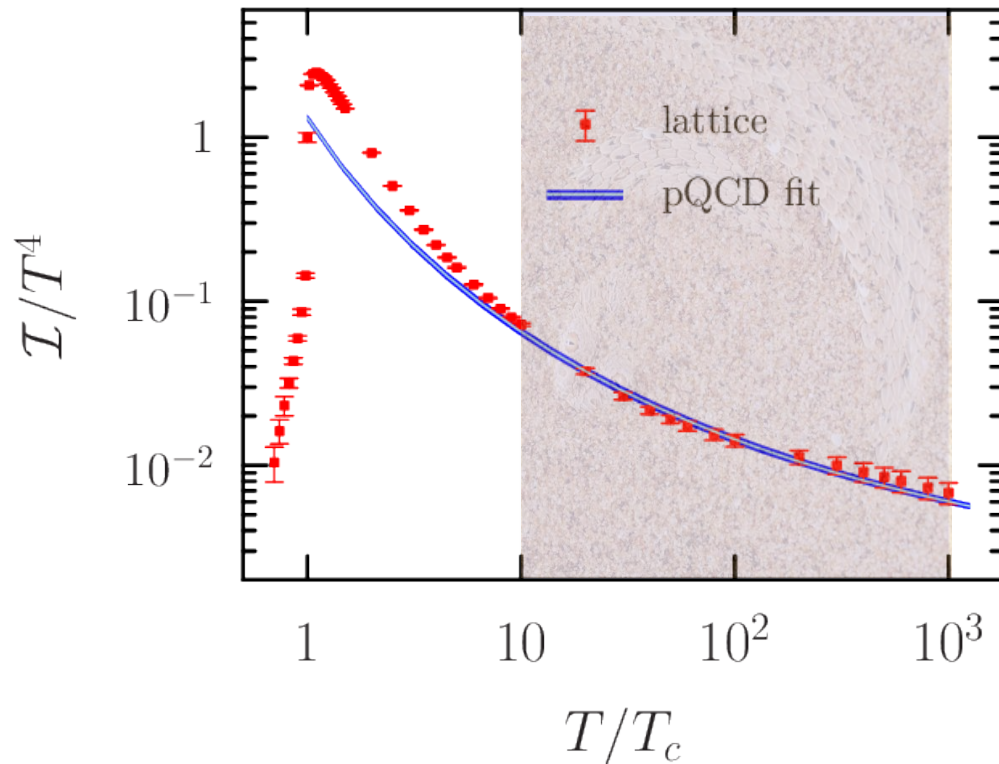
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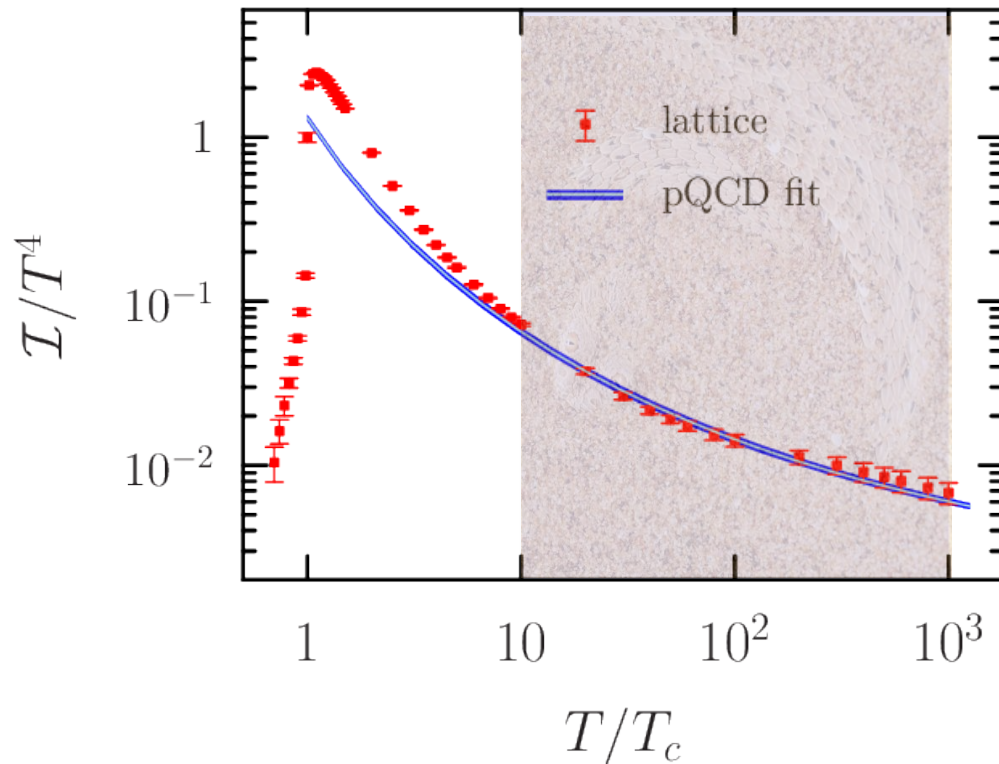
Scrutinize existing results [Borsanyi et al, 2012]



compare to pQCD $\Lambda = 0.73T_c$ known

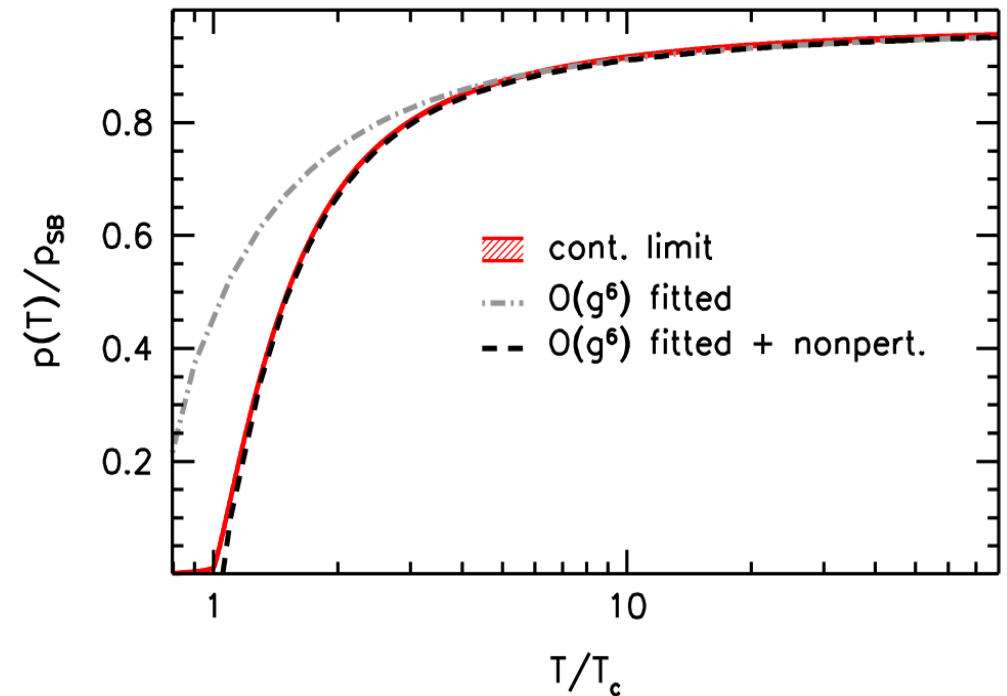
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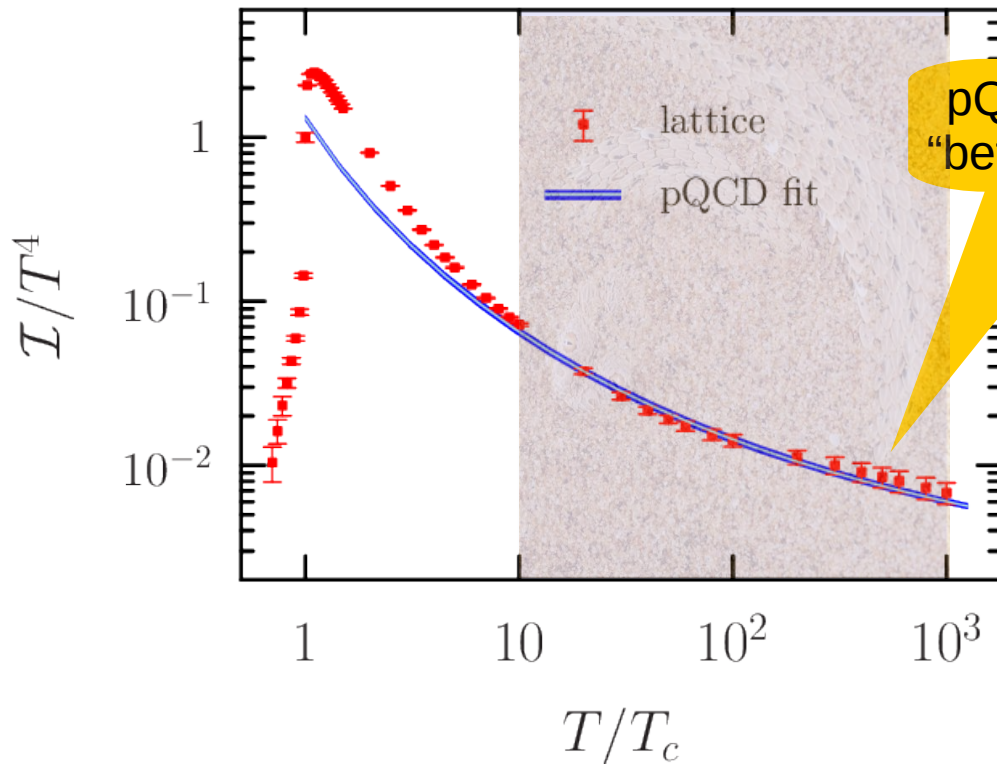


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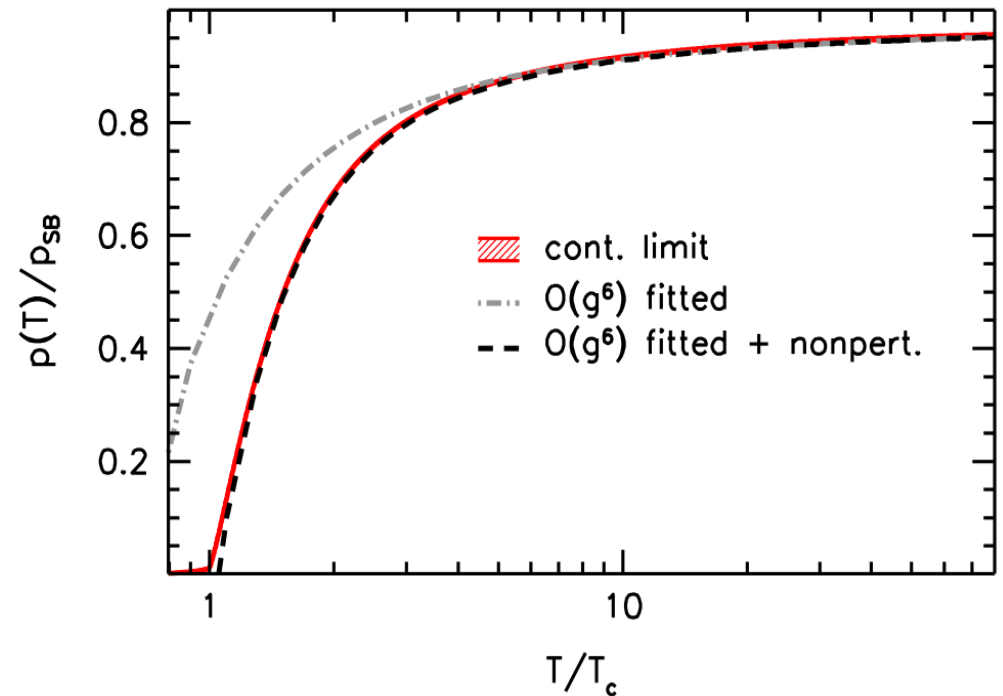


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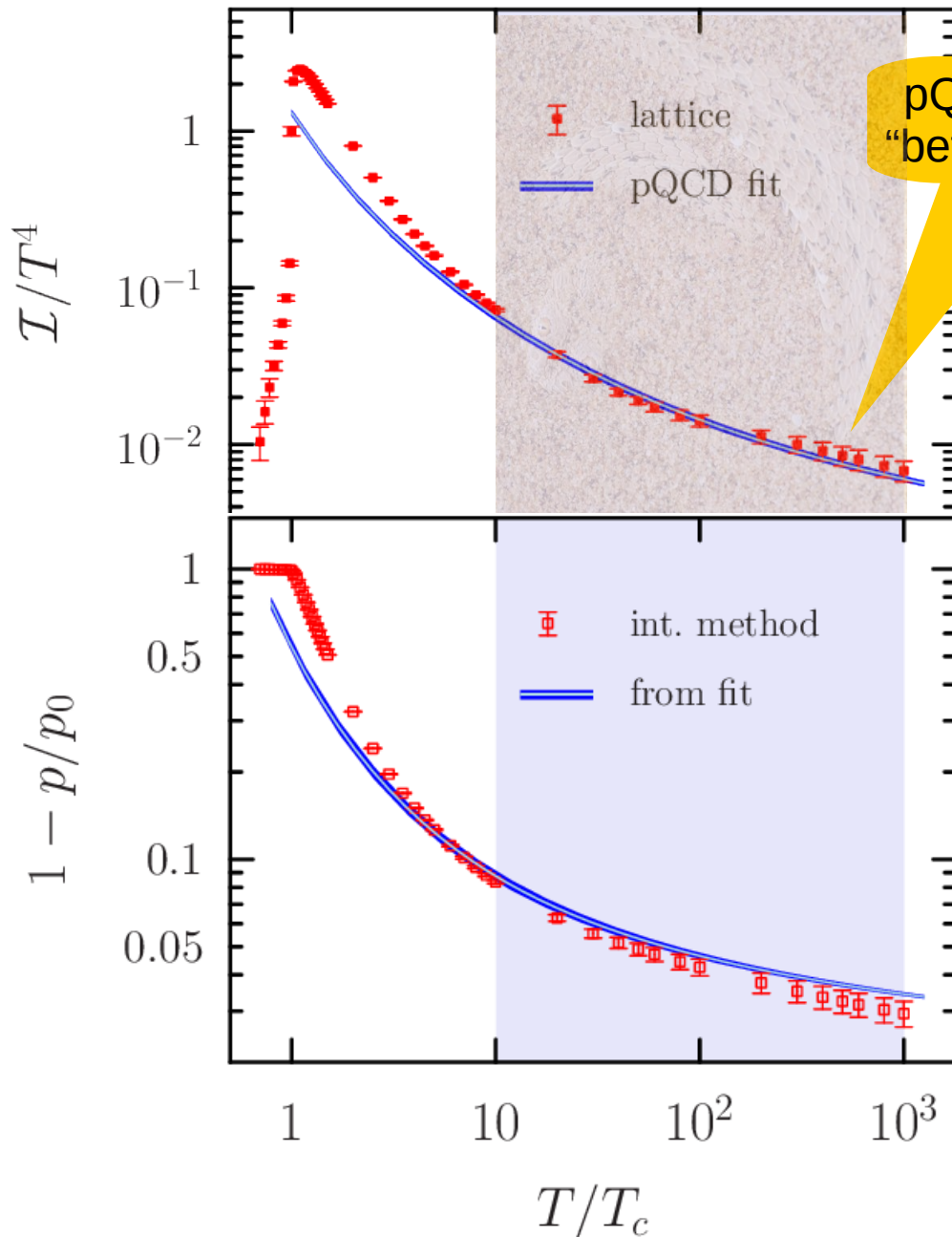


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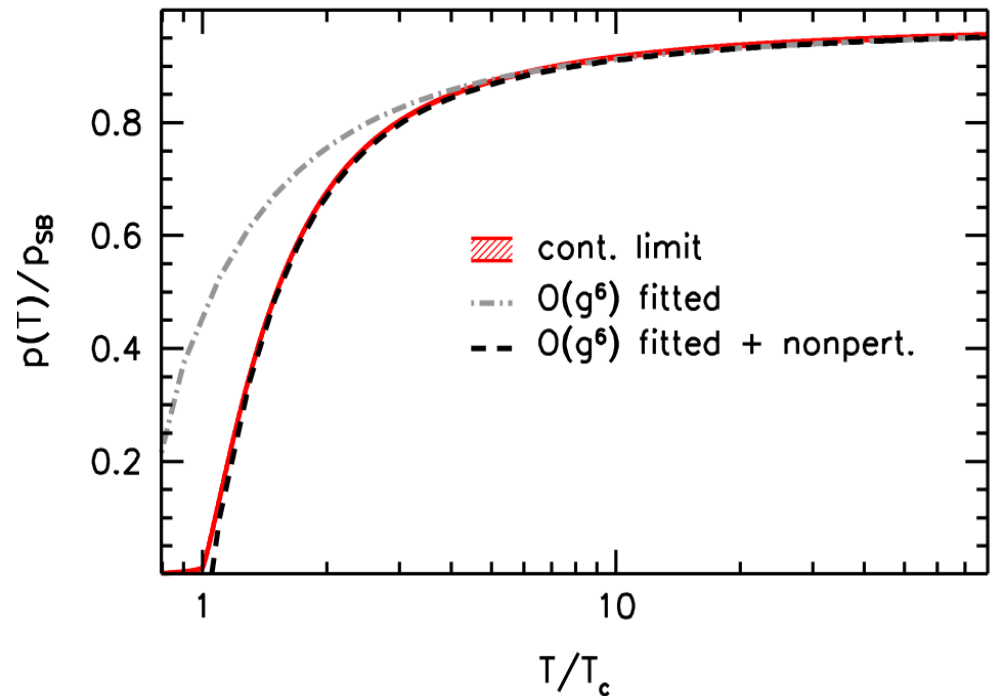


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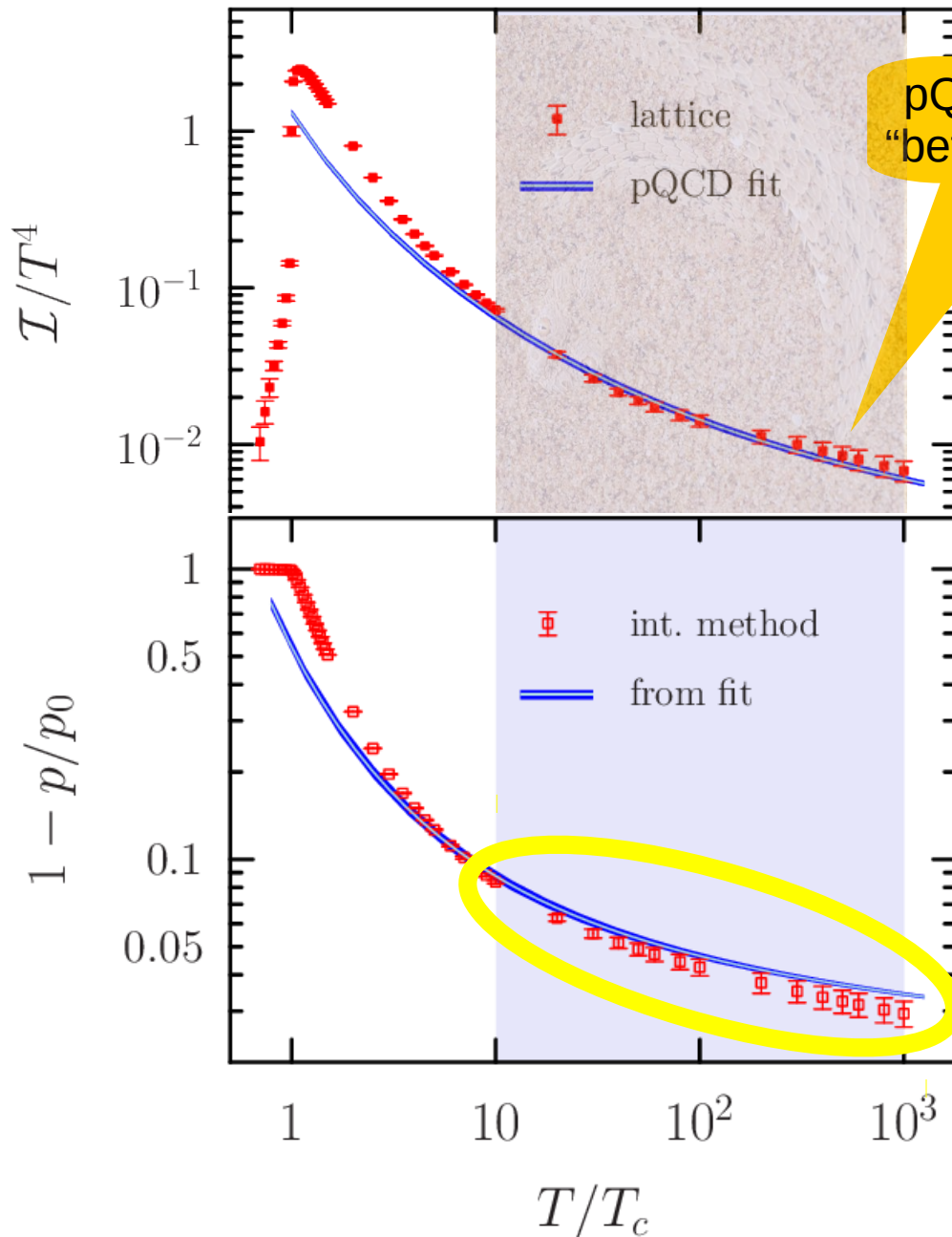


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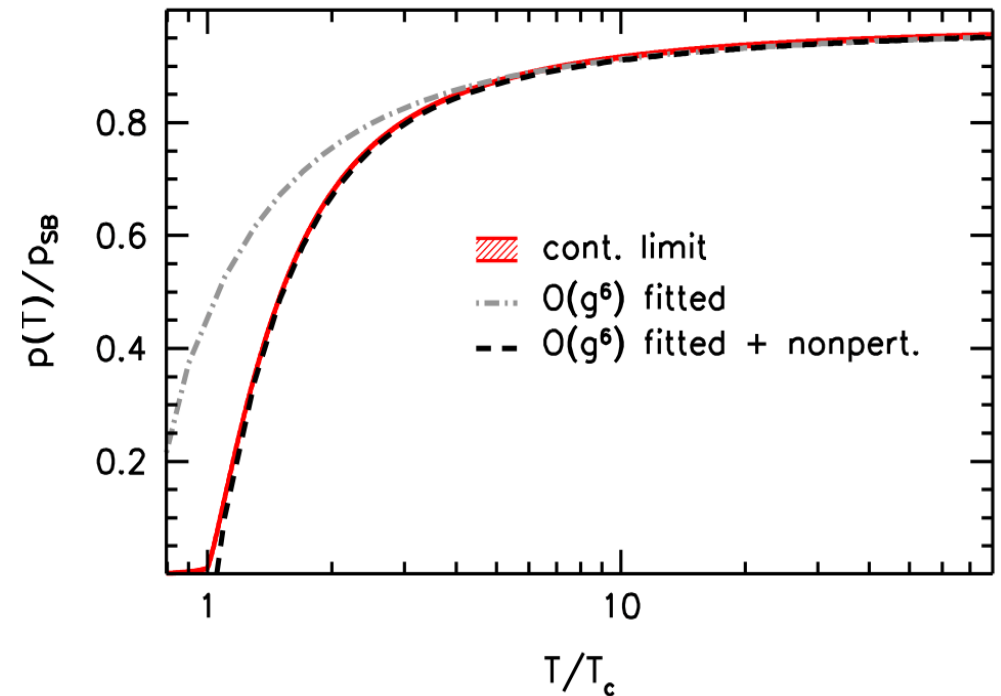


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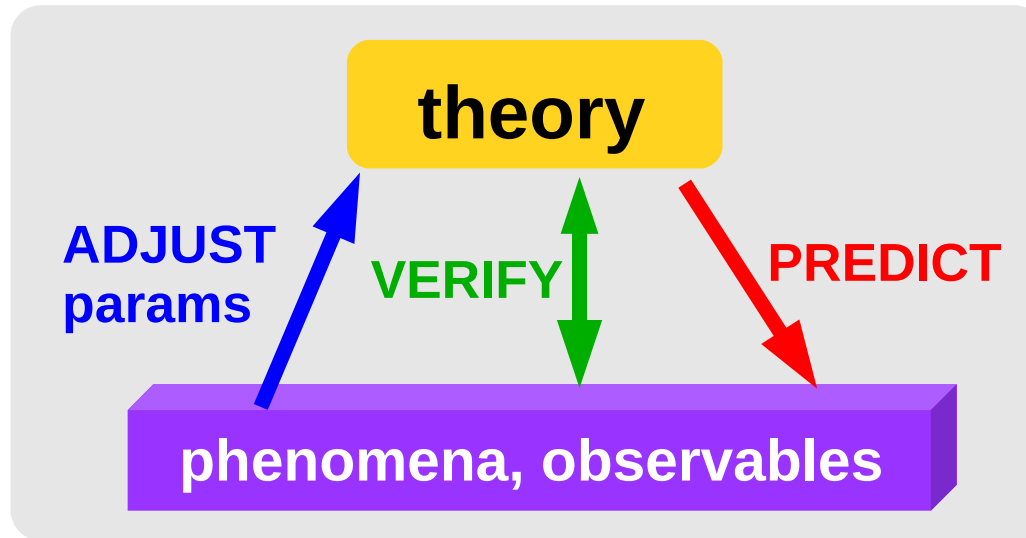


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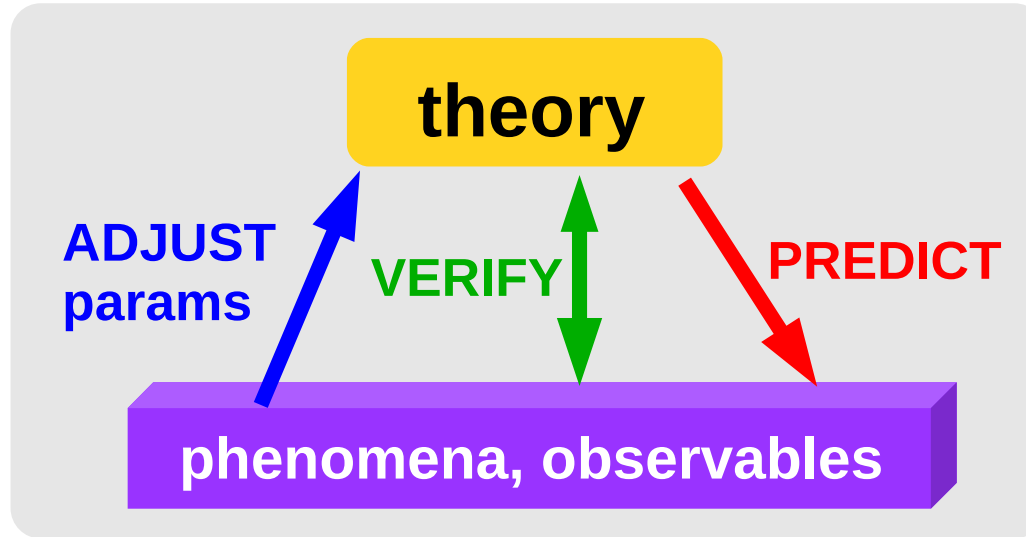
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Theory & models

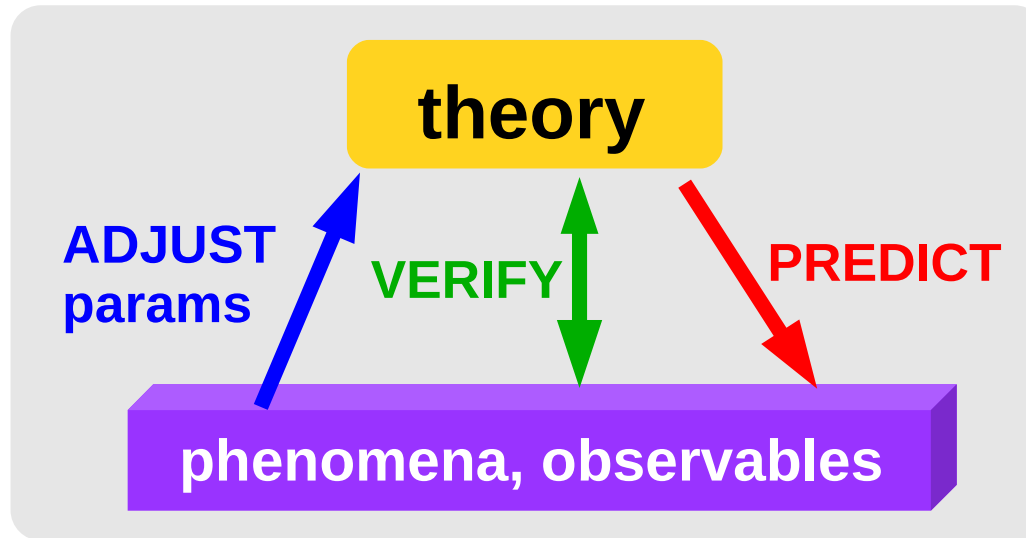


Theory & models



- in practice, we **always use models/approximations**
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Theory & models



- in practice, we **always use models/approximations**
 - ♦ **validity check** is crucial
- in QFT
 - ♦ adjusting parameters = **renormalization**

Our approach

- **thermodynamic renormalization:** match perturbative results to lattice data at sufficiently large “renormalization temperature”
to specify model parameter(s)
use interaction measure (being the actual lattice “observable”)

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- **make predictions for other observables** in applicability range

$(n|l)$ models ($n_f=0$)

pressure (= thermodynamic potential) to order n :

$$p_{(n)} = p_0 \left[1 + \sum_{m=2}^n C_m \alpha^{m/2} \right] \quad \text{where } p_0 = 8 \times 2 \frac{\pi^2}{90} T^4$$

$$C_2 = -1.2$$

$$C_3 = +5.4$$

$$C_4 = 6.8 \ln \alpha + 16.2$$

$$C_5 = -45.7$$

$$C_6 = -36.6 \ln \alpha + c_6 \quad (\text{for } \mu = 2\pi T)$$

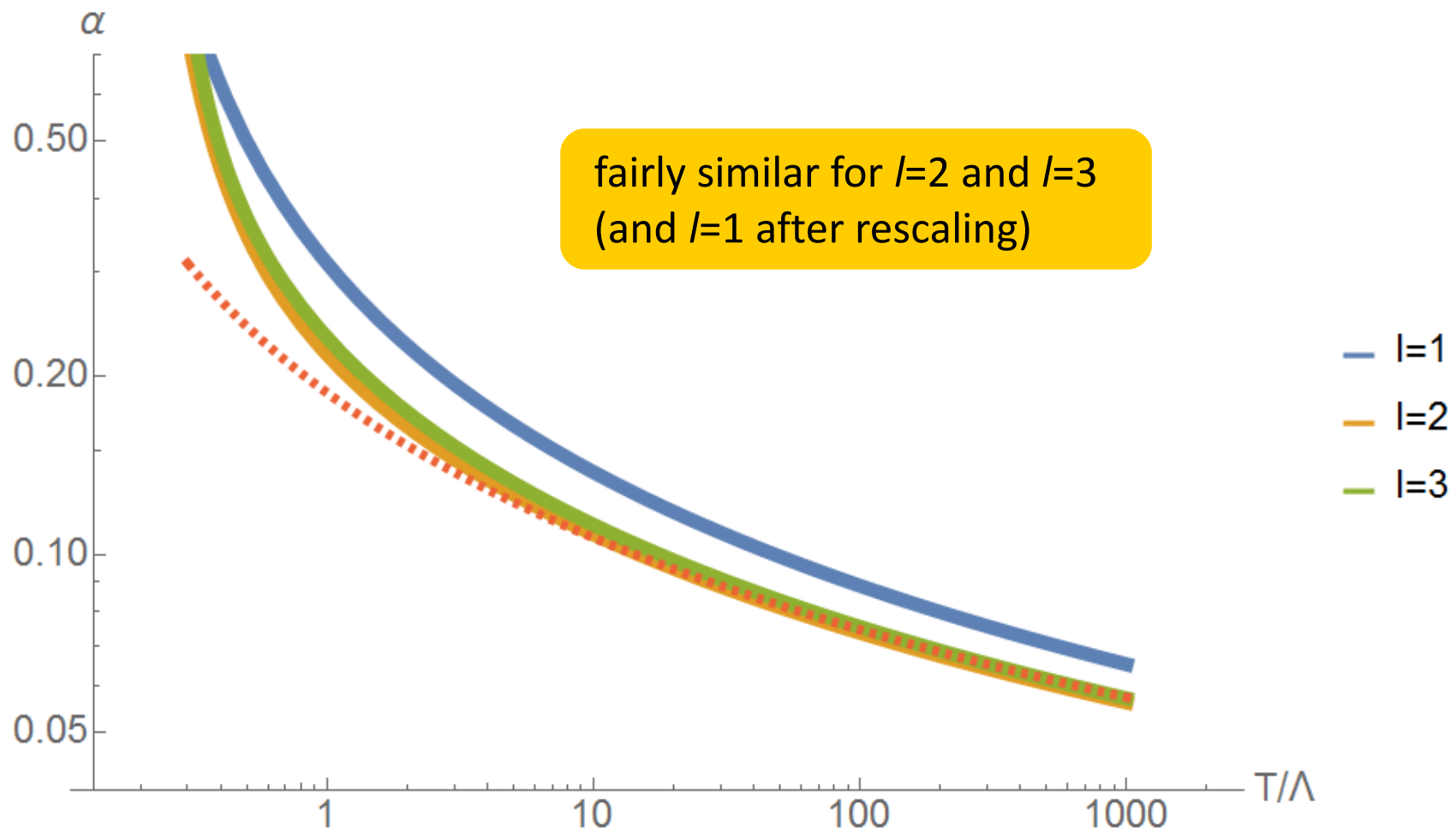
running coupling to order l :

$$\alpha_{(\ell)} = \sum_{k=1}^{\ell} a_k(L) L^{-k}$$

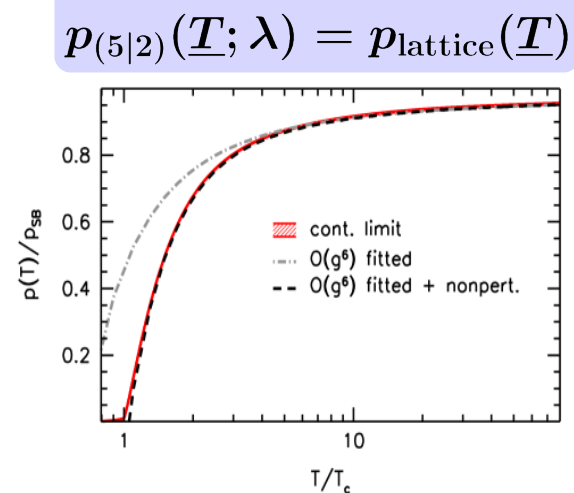
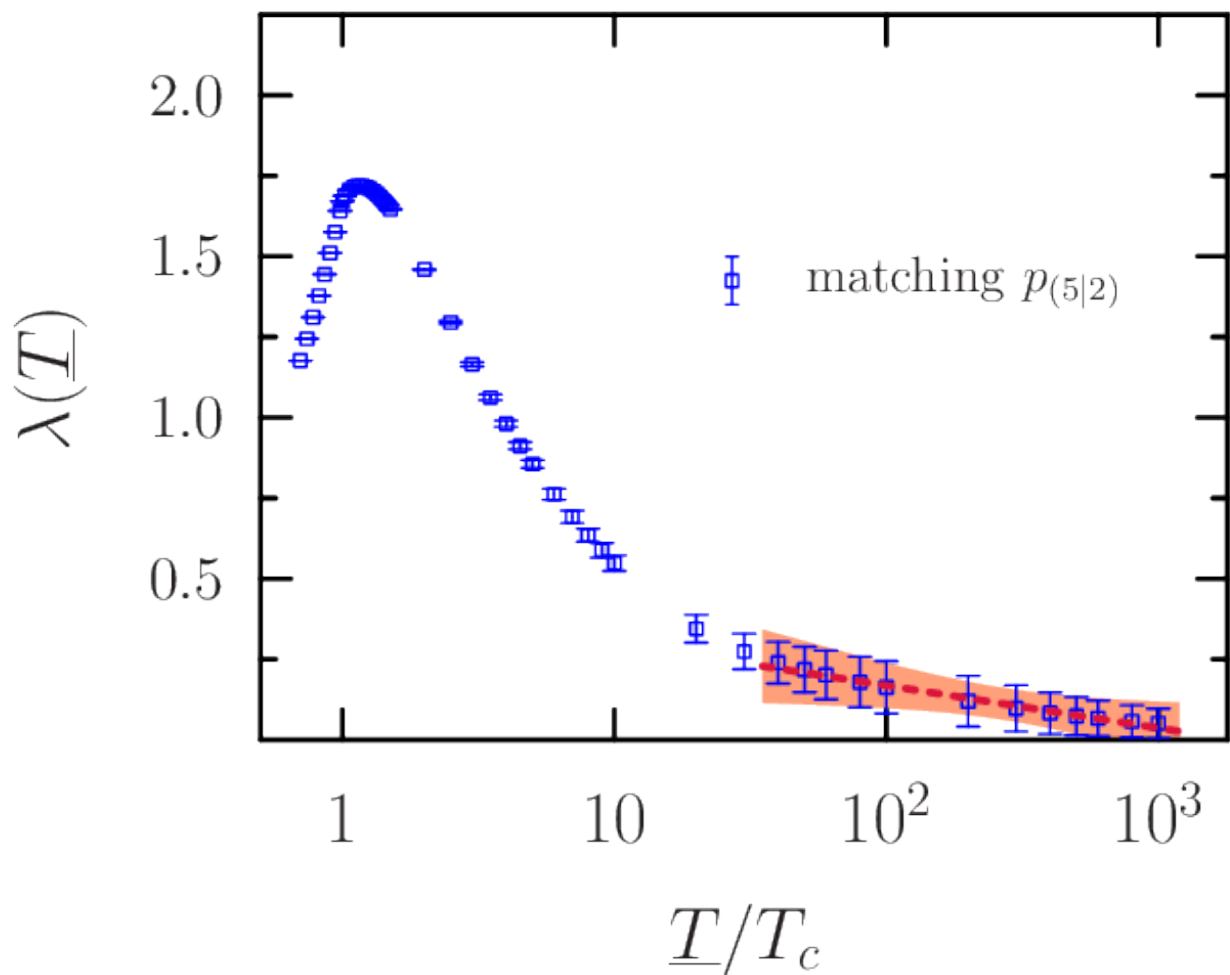
$$a_1 = 1.14, a_2 = -0.96 \ln L, a_3 = 0.41 + 0.81(\ln L - 1) \ln L$$

$$L(T) = \ln \left(\frac{2\pi T}{\lambda T_c} \right)^2 \quad \text{where } \lambda = \Lambda/T_c$$

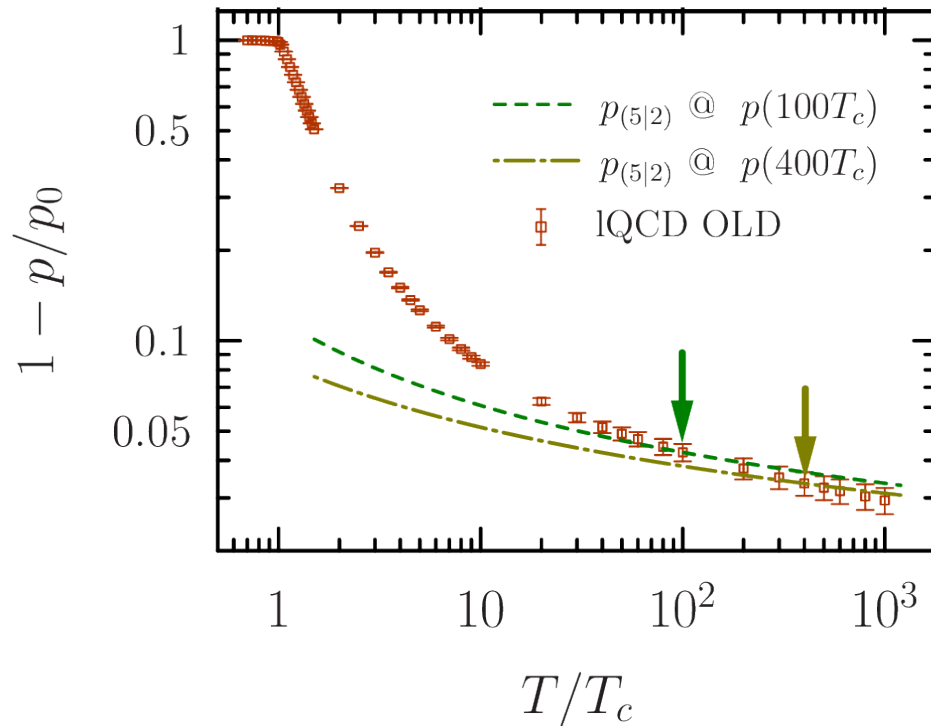
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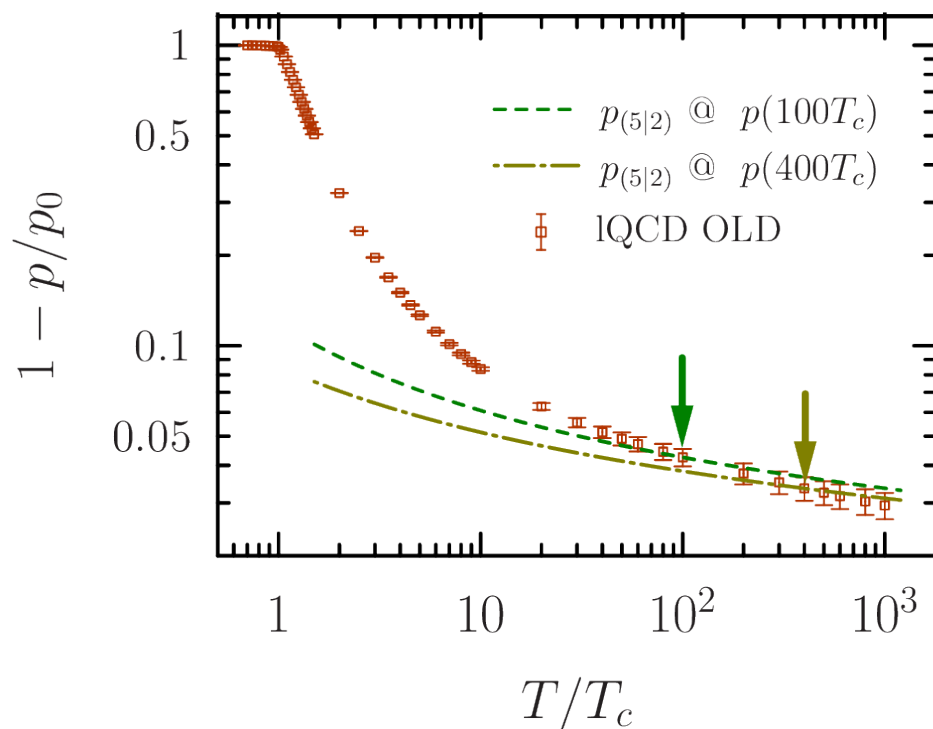
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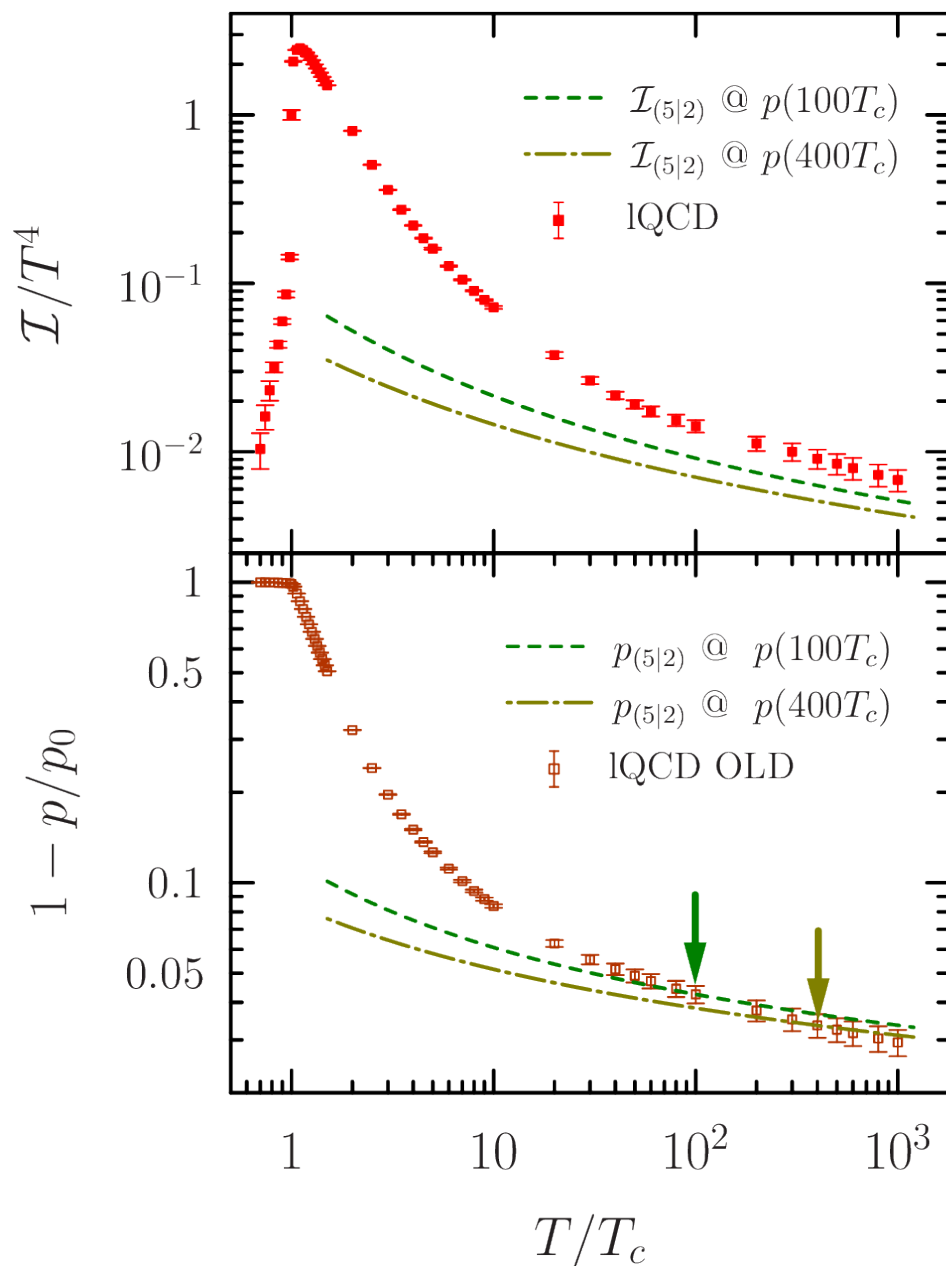


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small applicability range,
not to $T \rightarrow \infty$

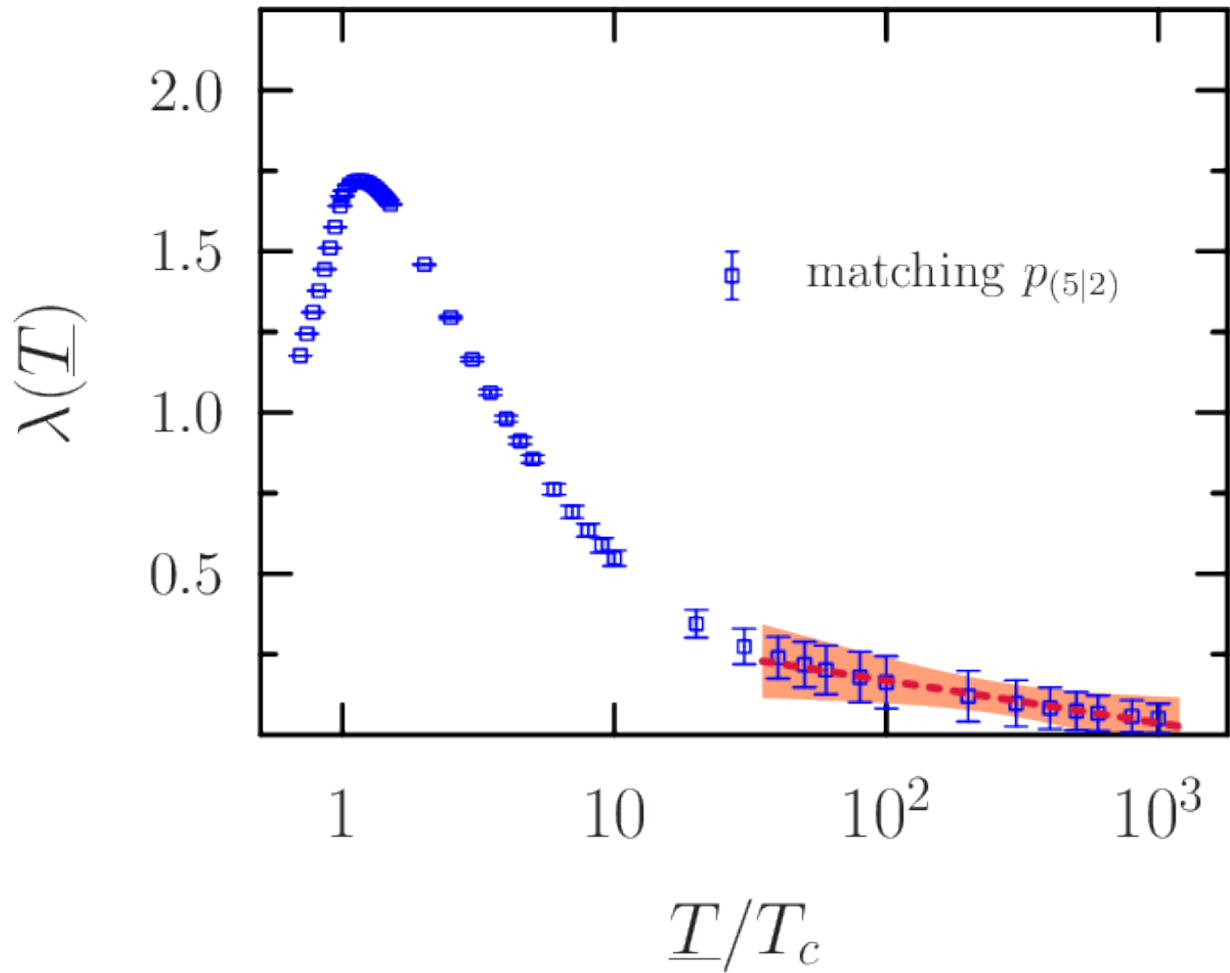
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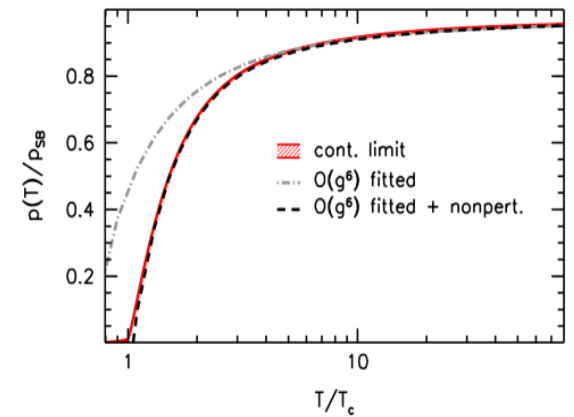
systematic discrepancy for interaction measure as actual lattice “observable”

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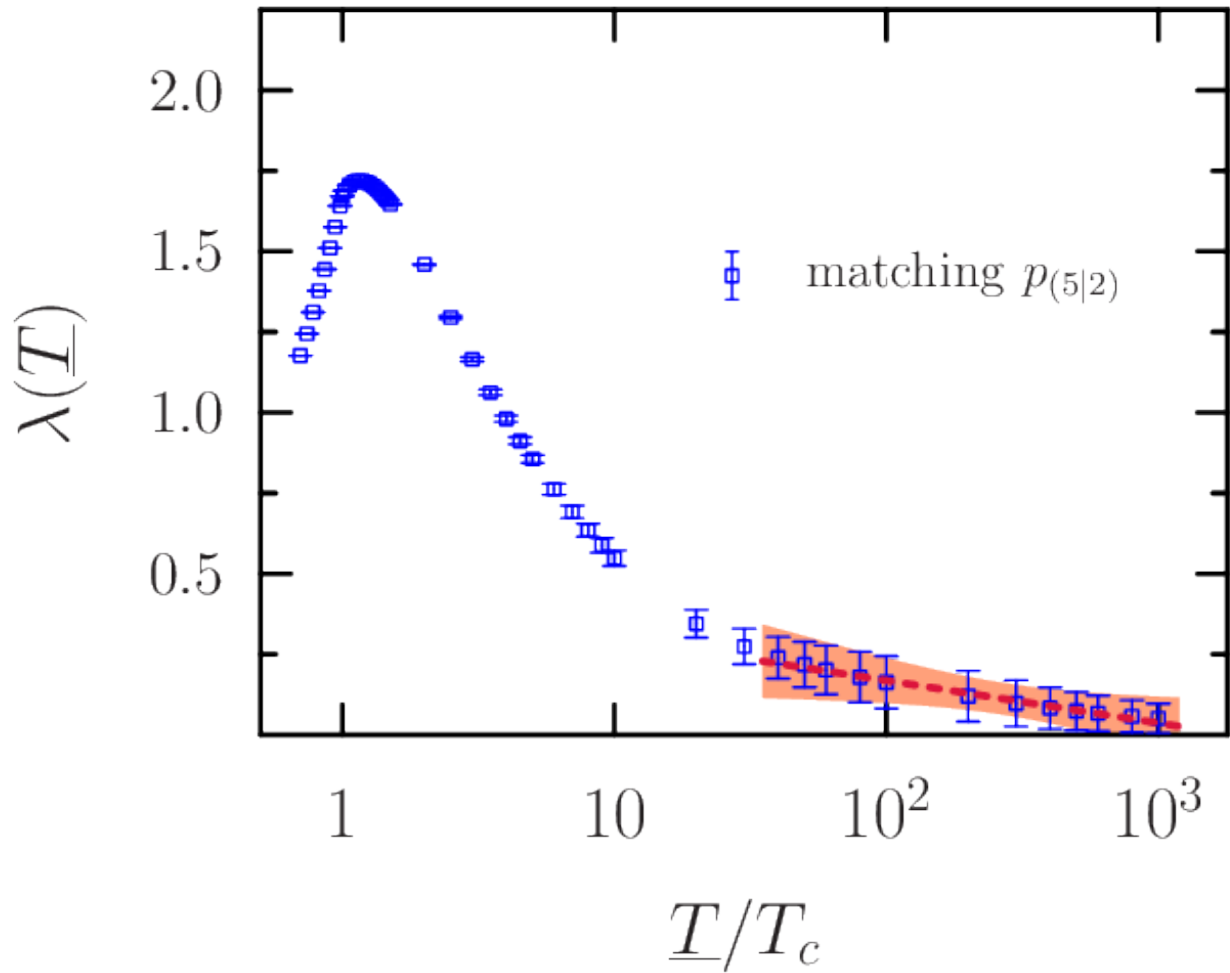
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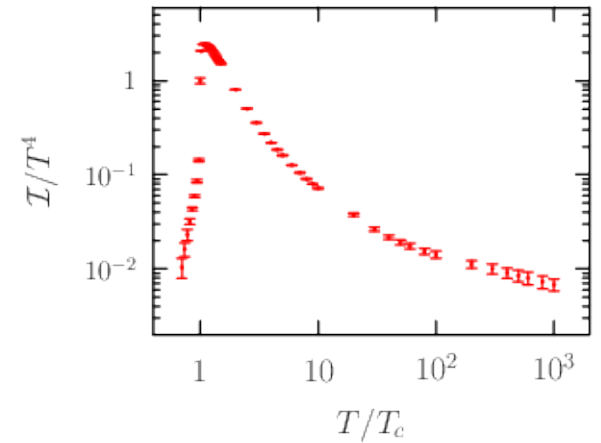
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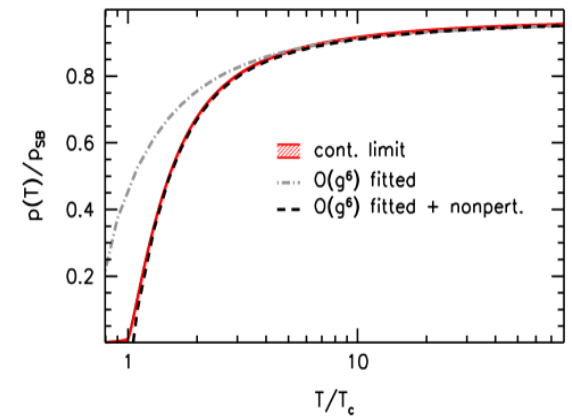
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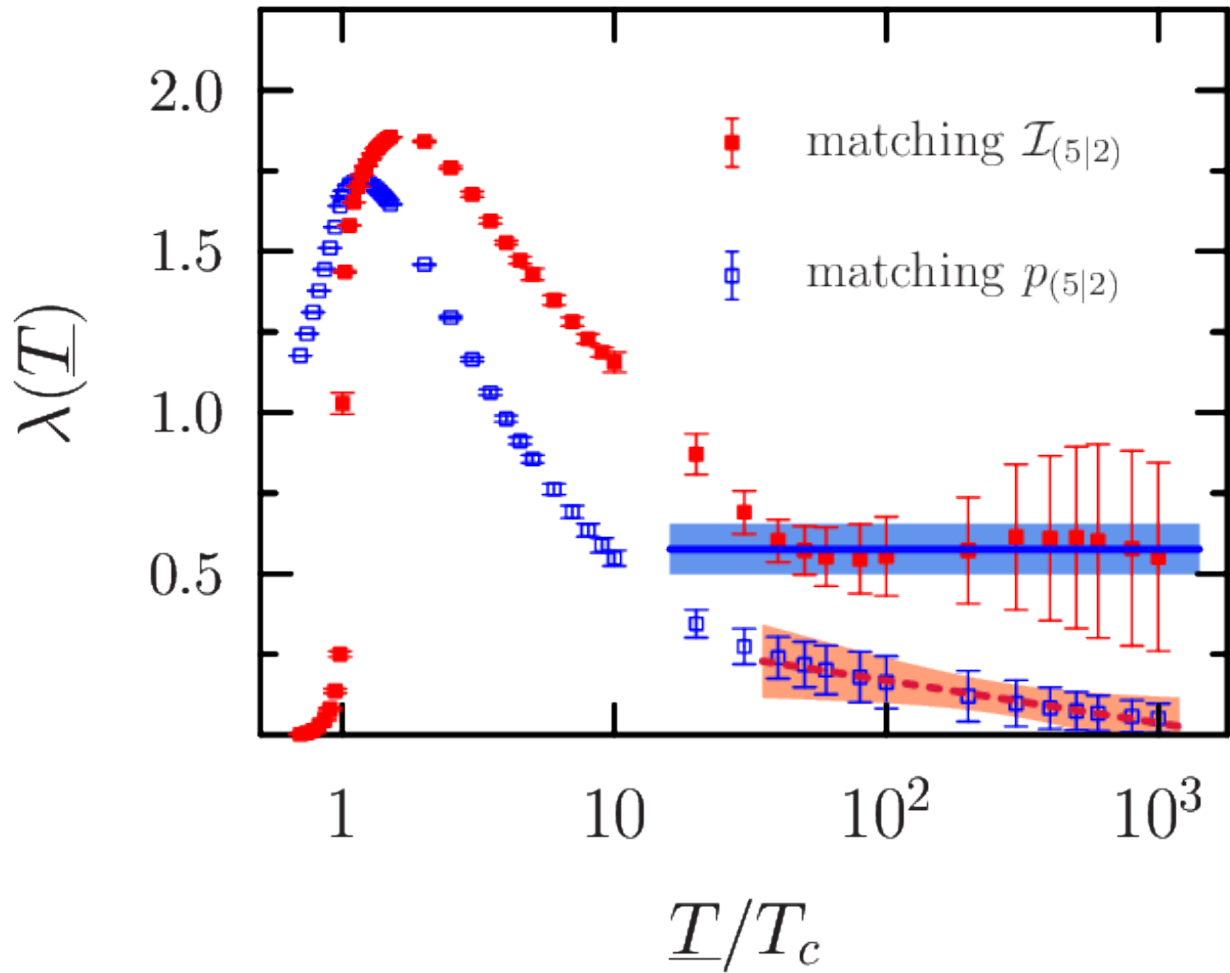
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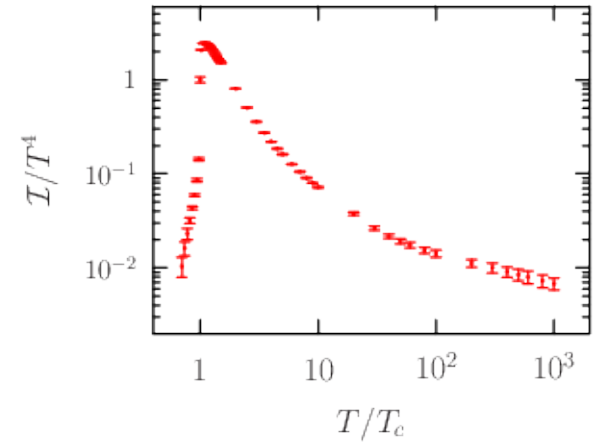
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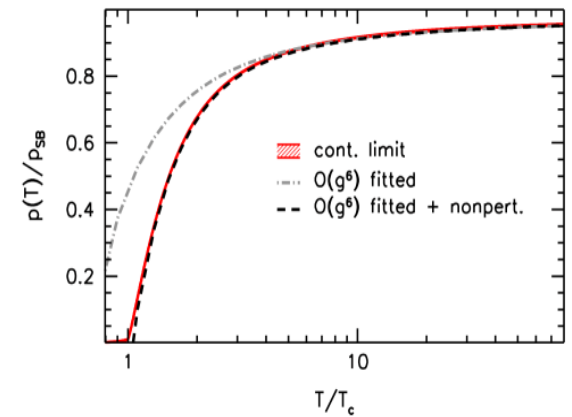
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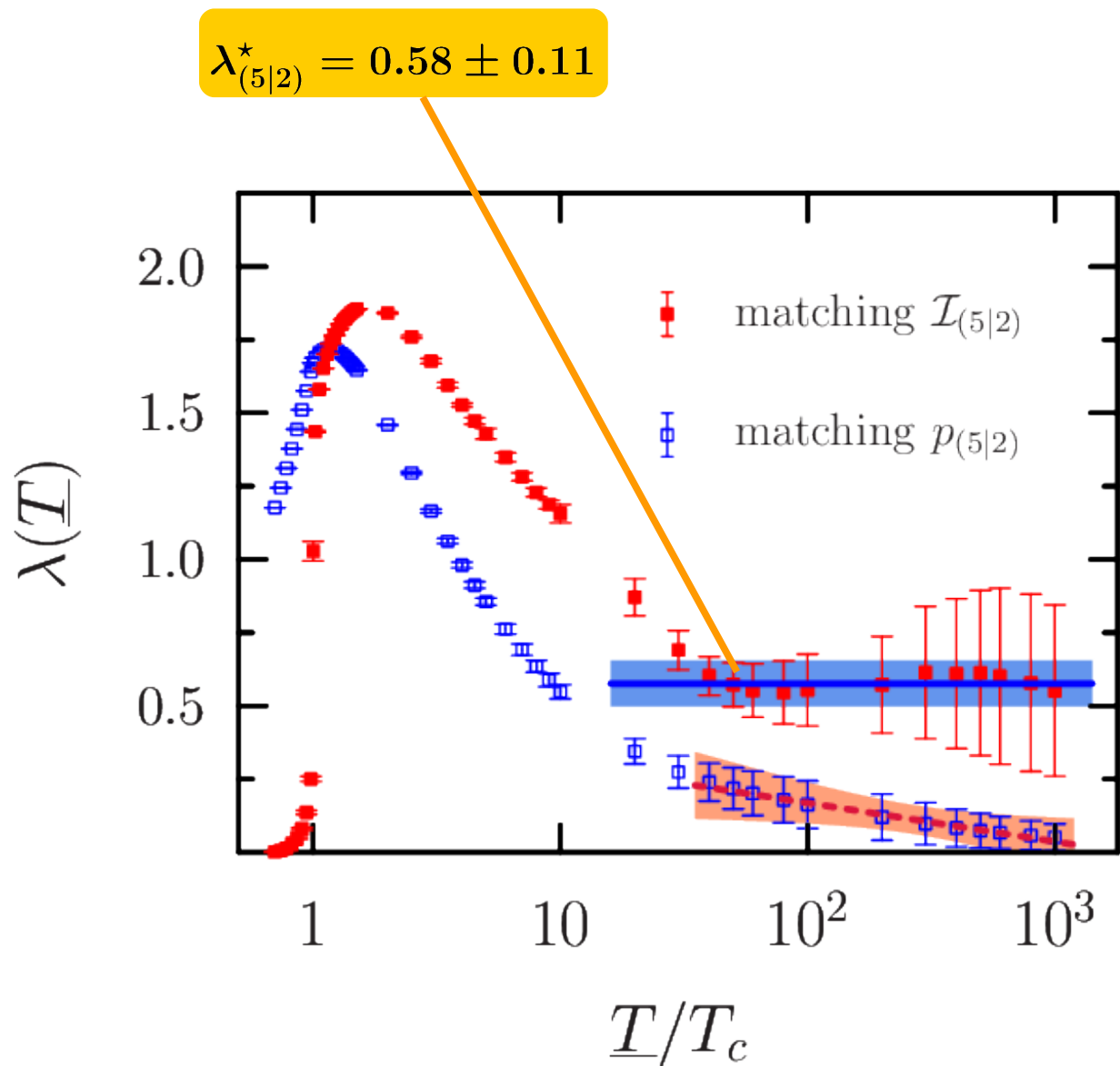
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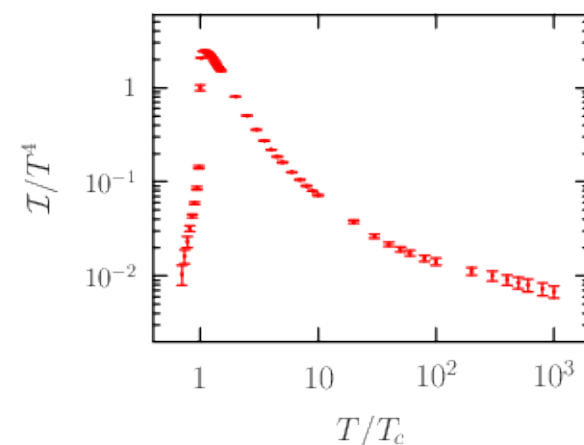
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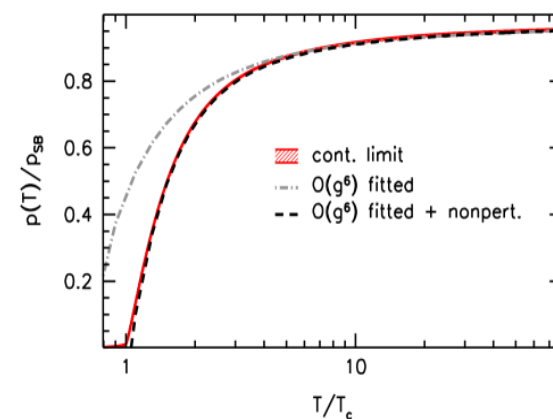
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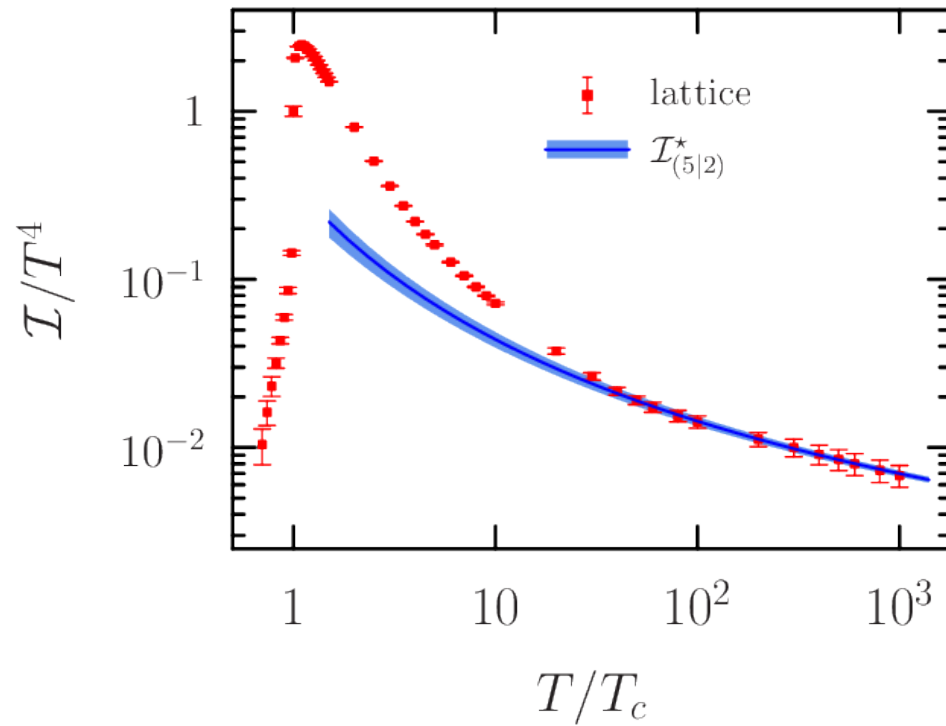
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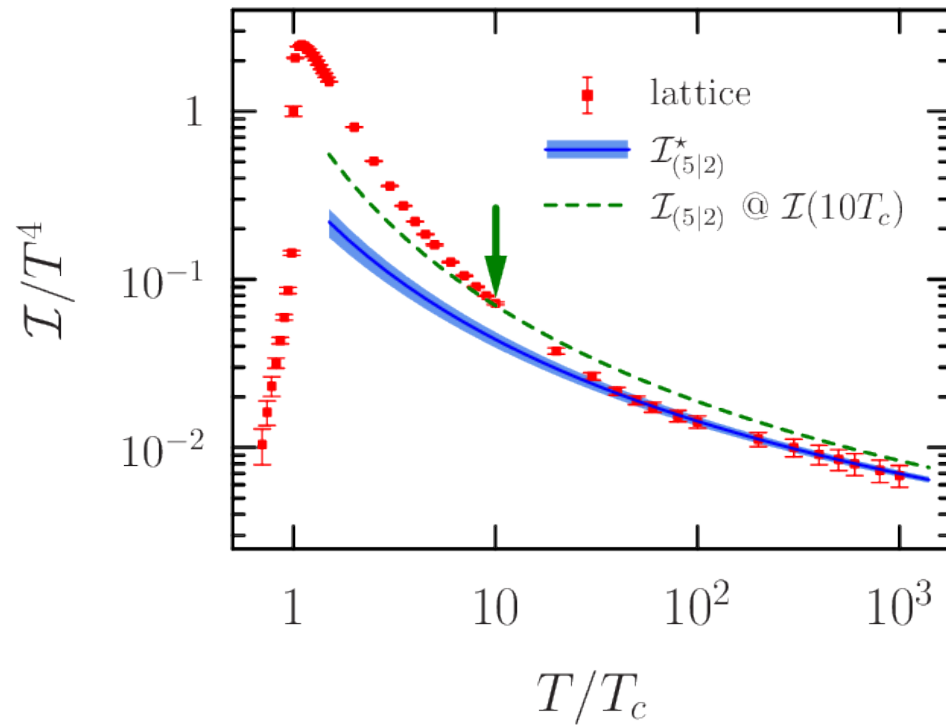


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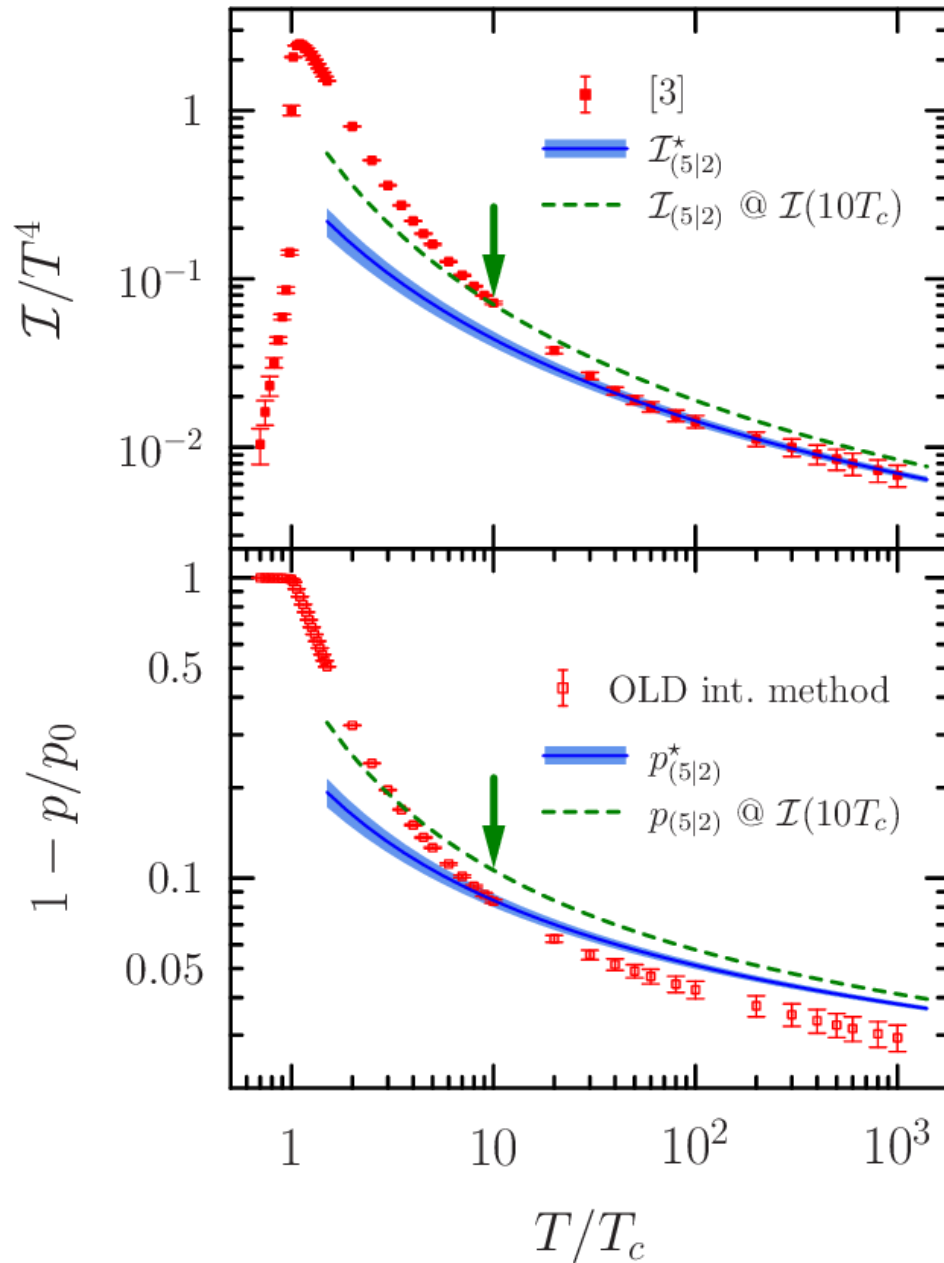
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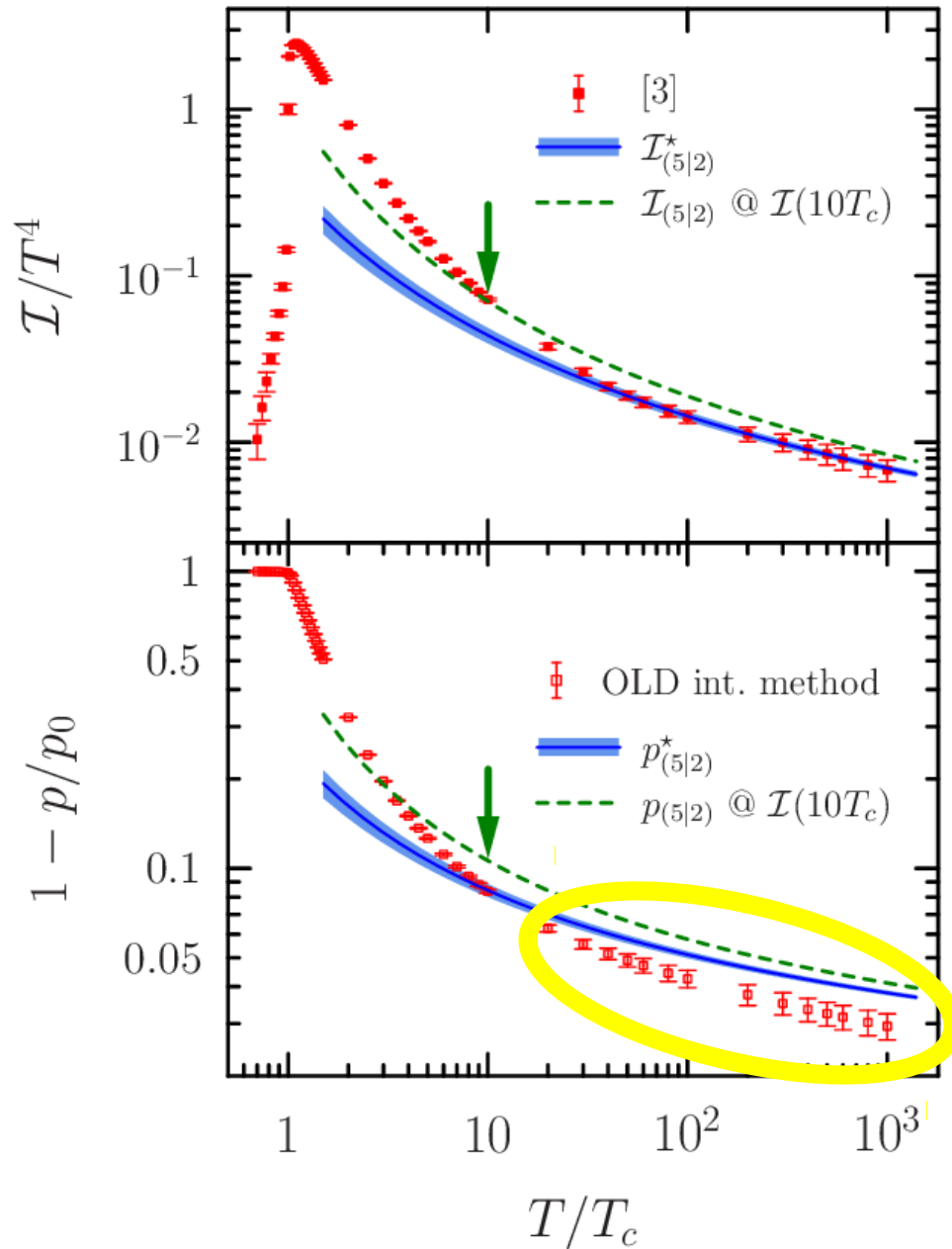
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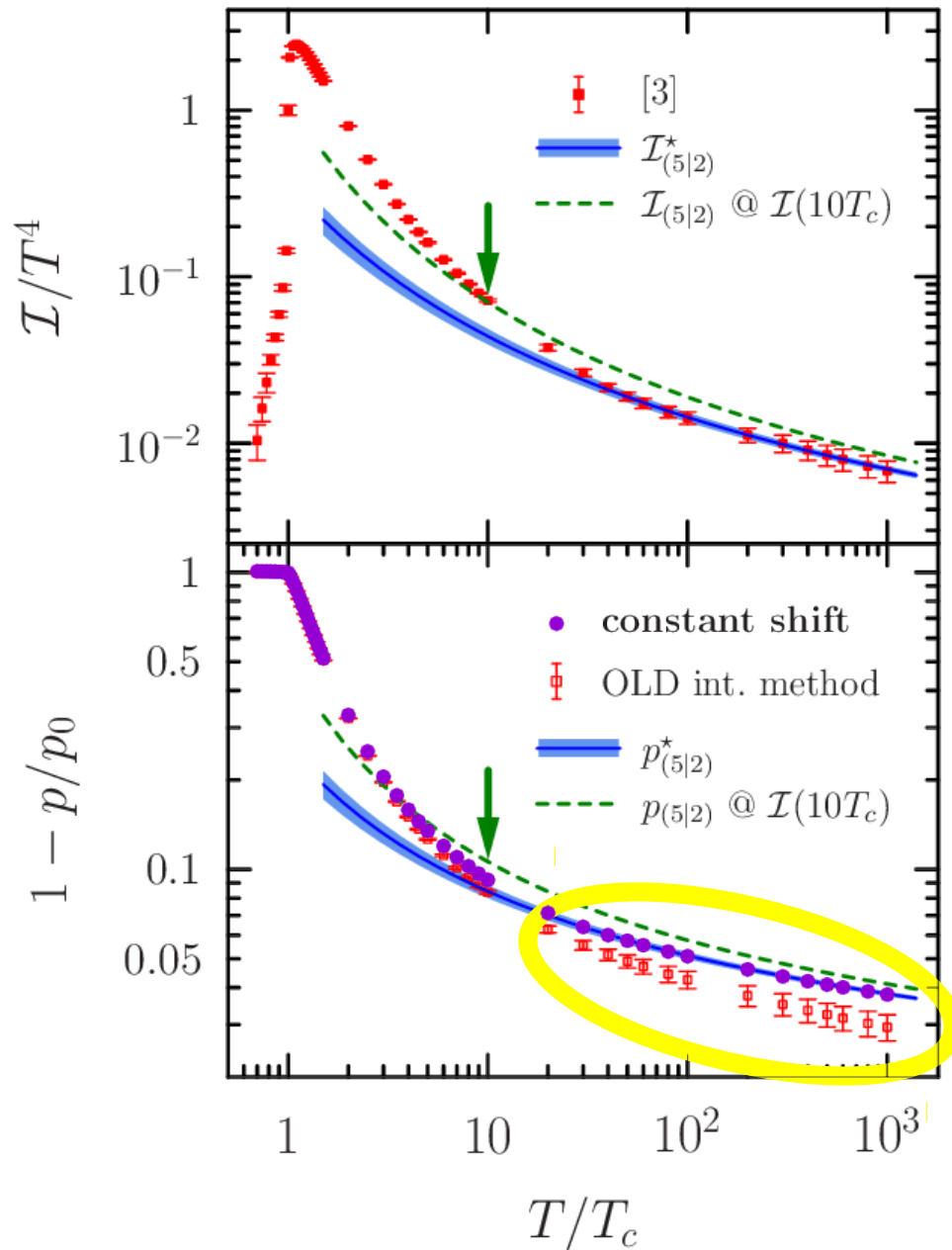
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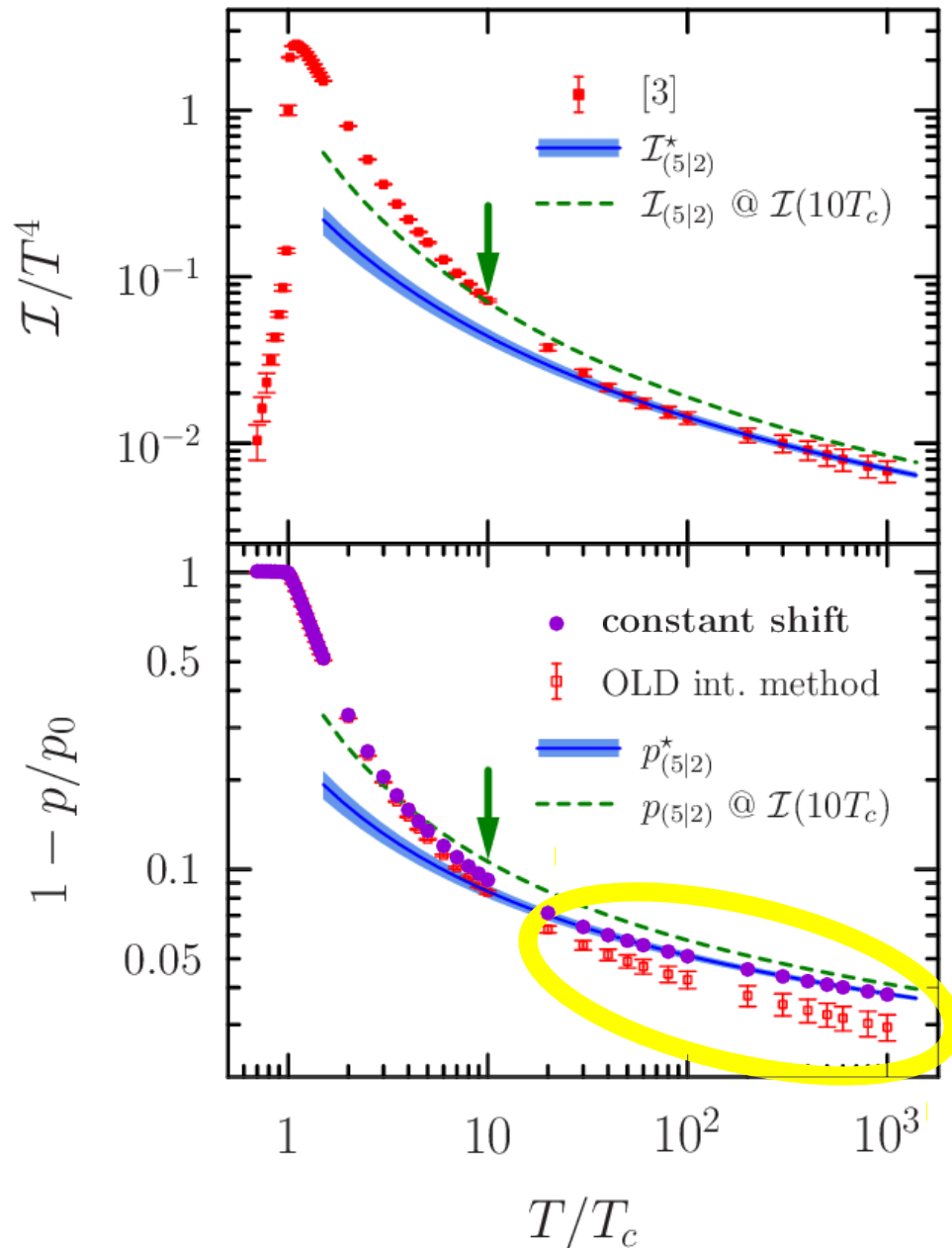
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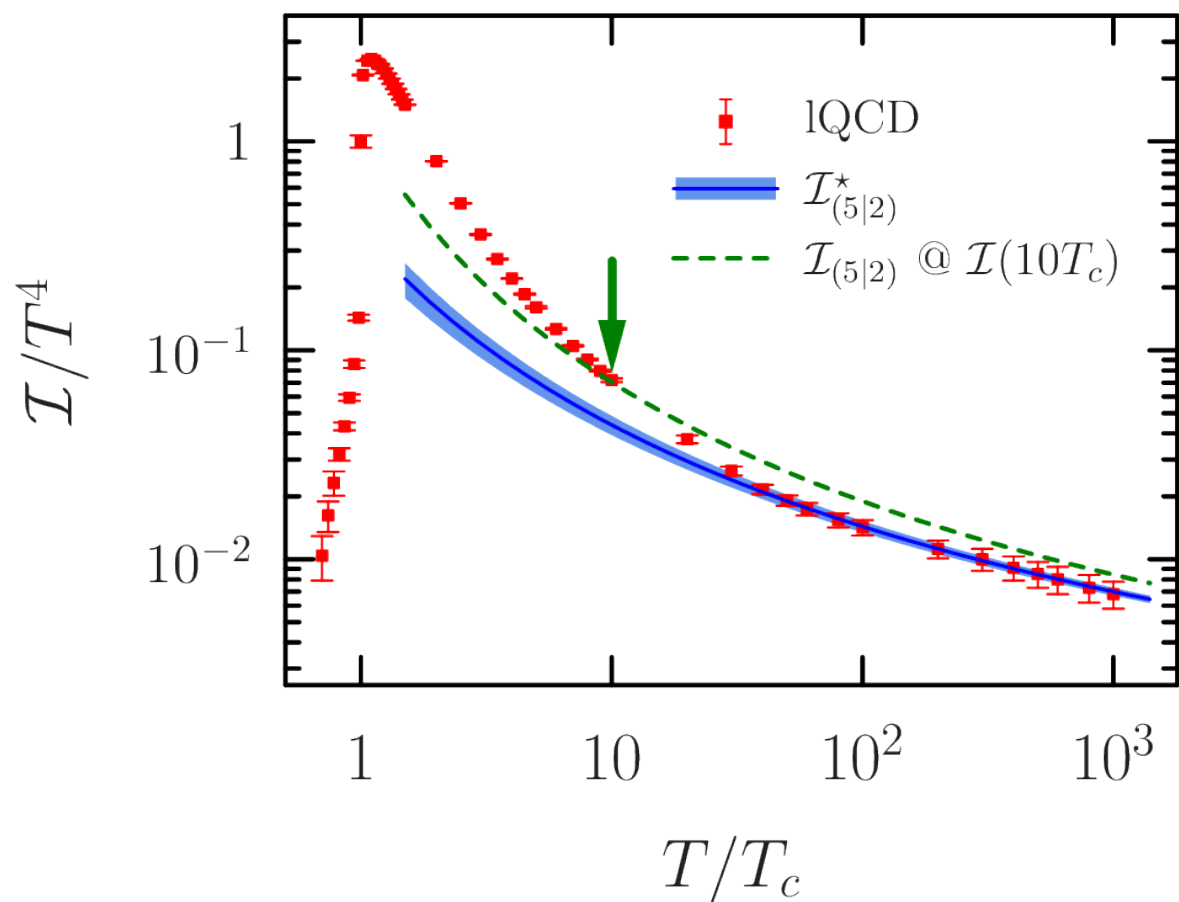
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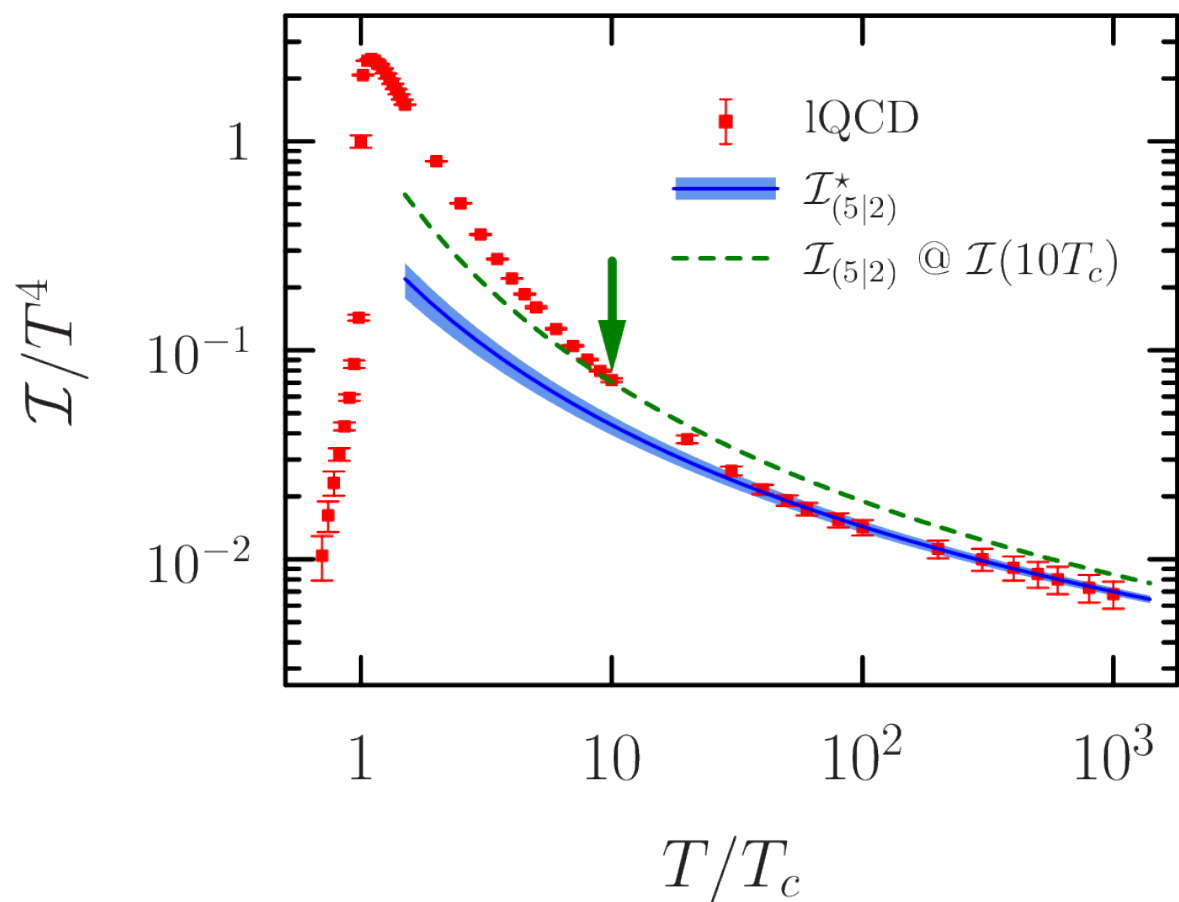
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breakdown at $T^* \sim 40T_c$ because “coupling too large” ...?



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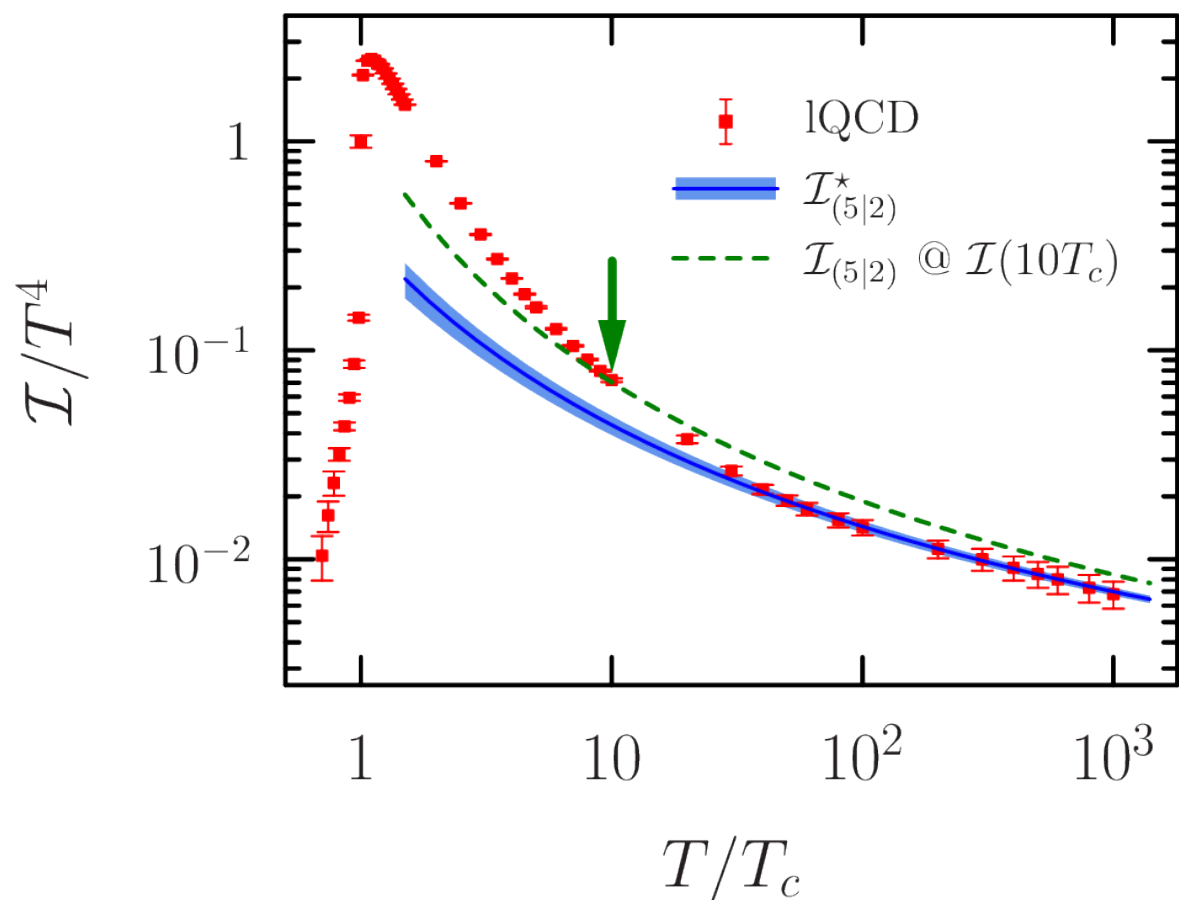
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$$\begin{aligned}\alpha(10T_c) &= 0.10 \\ \alpha(40T_c) &= 0.08 \\ \alpha(400T_c) &= 0.06\end{aligned}$$

(5|2) model

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$$\alpha(10T_c) = 0.10$$

$$\alpha(40T_c) = 0.08$$

$$\alpha(400T_c) = 0.06$$

similar properties as
asymptotic series

$$p_{(5|2)}^*(40T_c) = p_0[1 - 0.09 + 0.12 - 0.01 - 0.08] \approx p_0[1 - \frac{1}{2}0.09]$$

(6|3) model

more difficult

- 2 parameters

$$\lambda_{(6|3)}, c_6$$

- expect smaller applicability range

(6|3) model

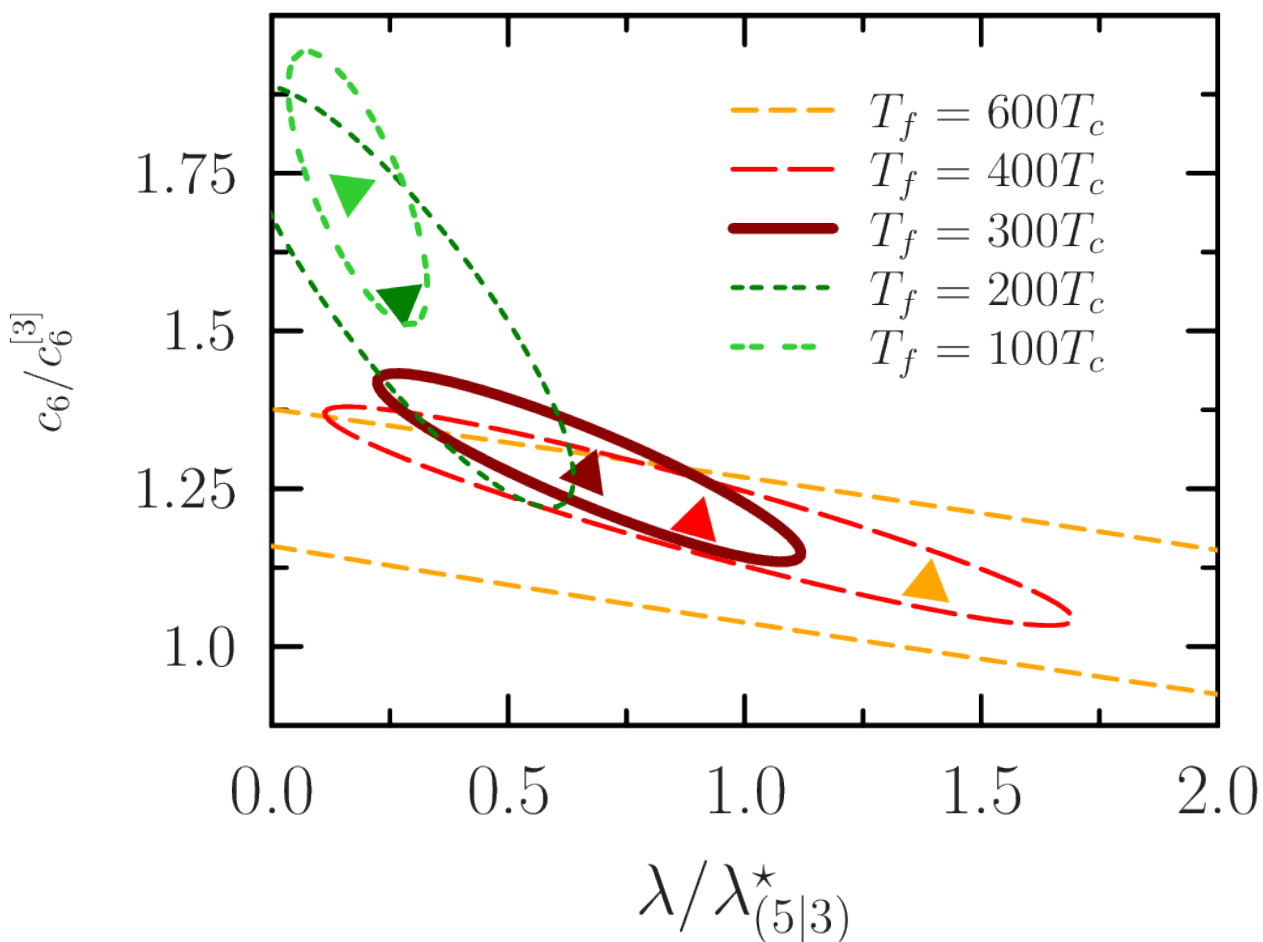
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fit over interval $[T_f, T_{\max}]$



(6|3) model

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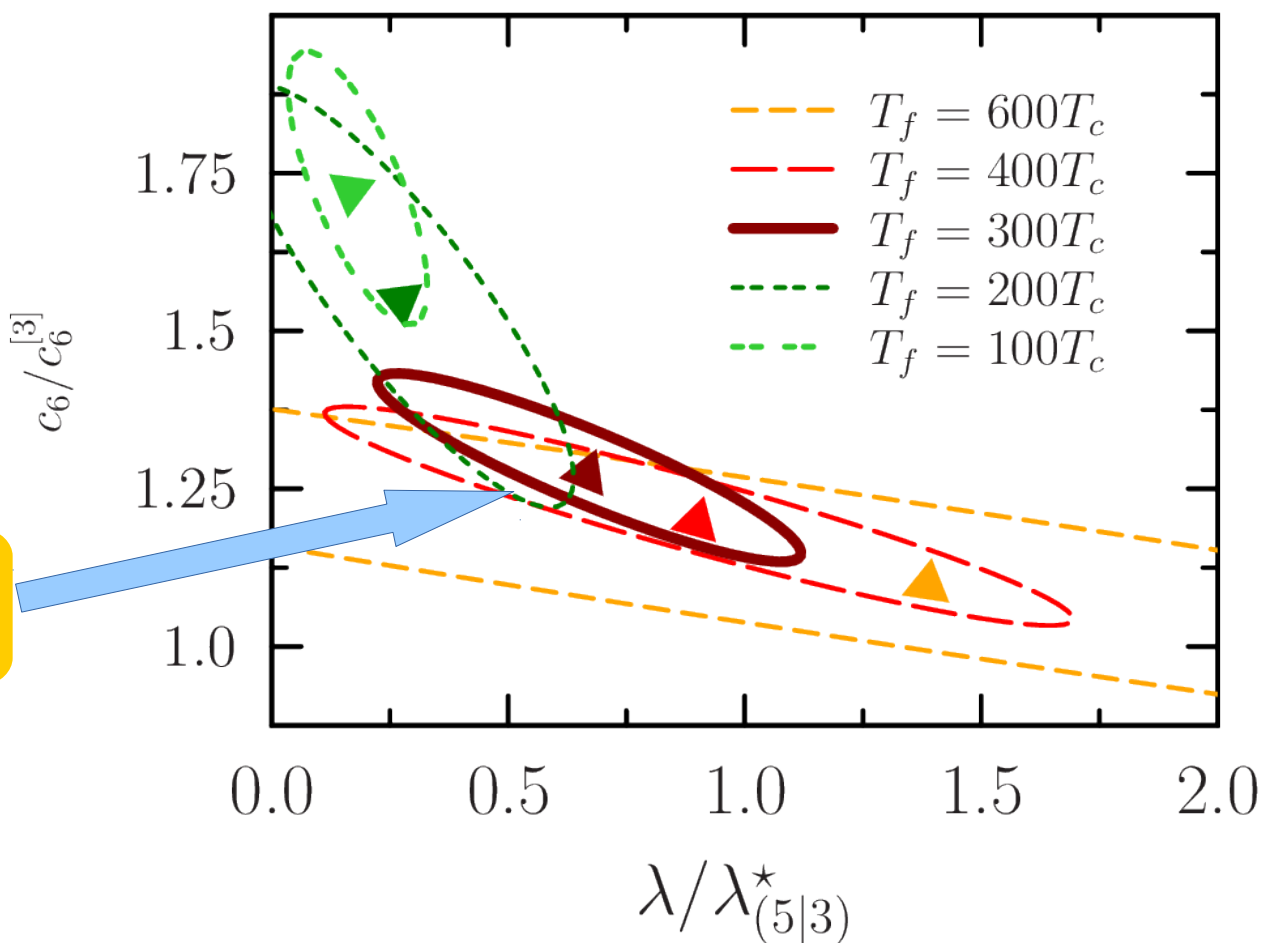
- 2 parameters

$$\lambda_{(6|3)}, c_6$$

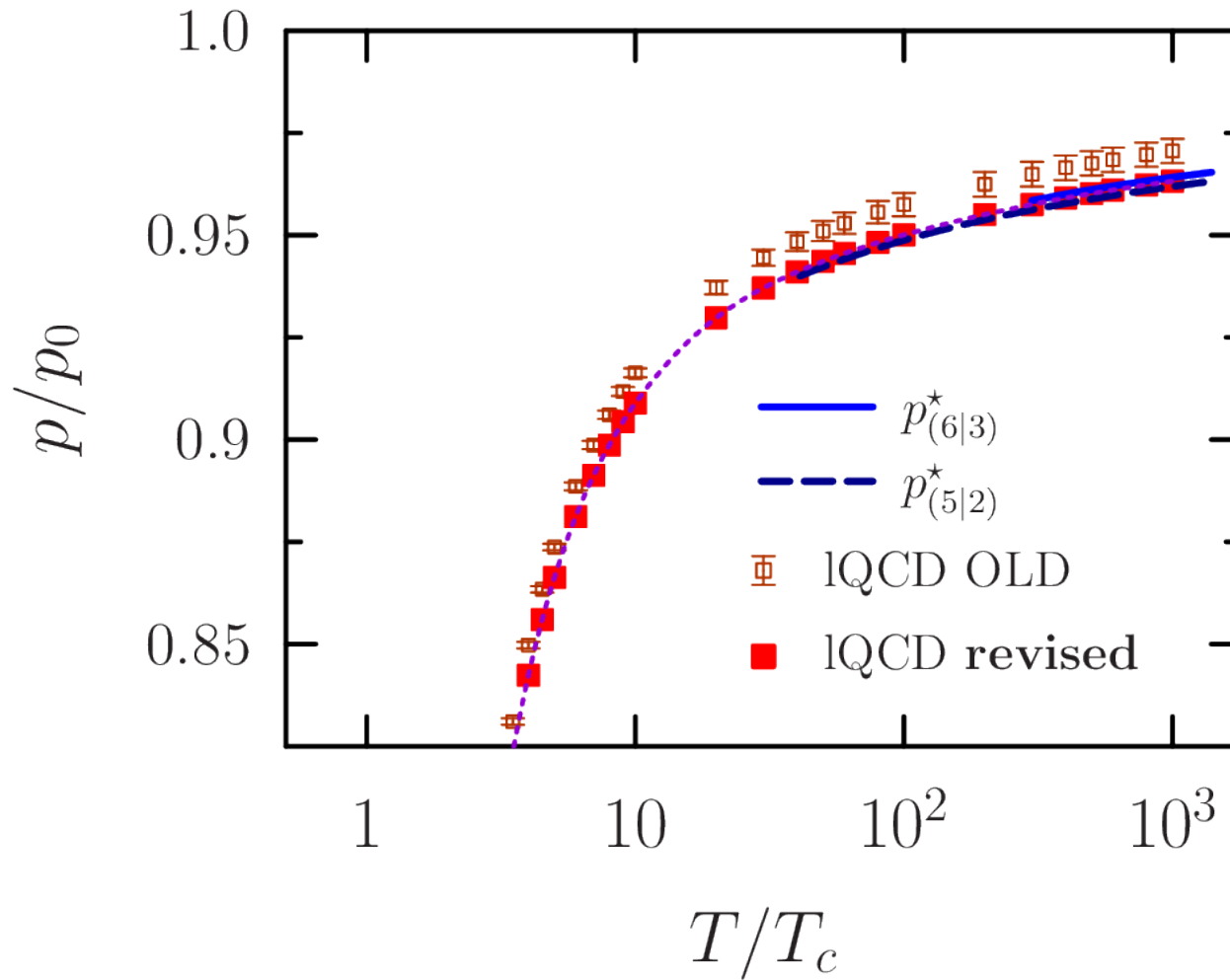
- expect smaller applicability range

fit over interval $[T_f, T_{\max}]$

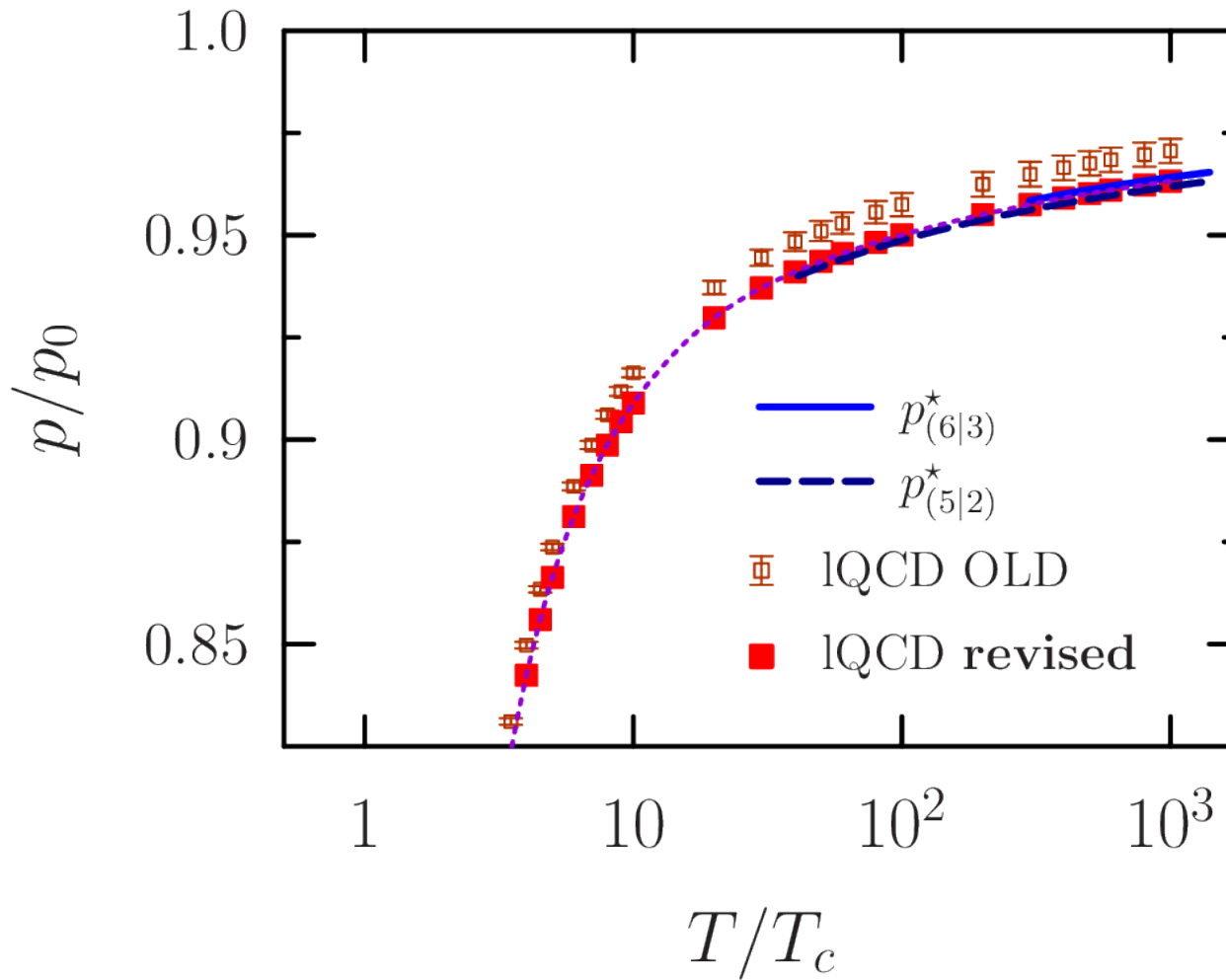
applicability range: $T > 300T_c$



Results



Results



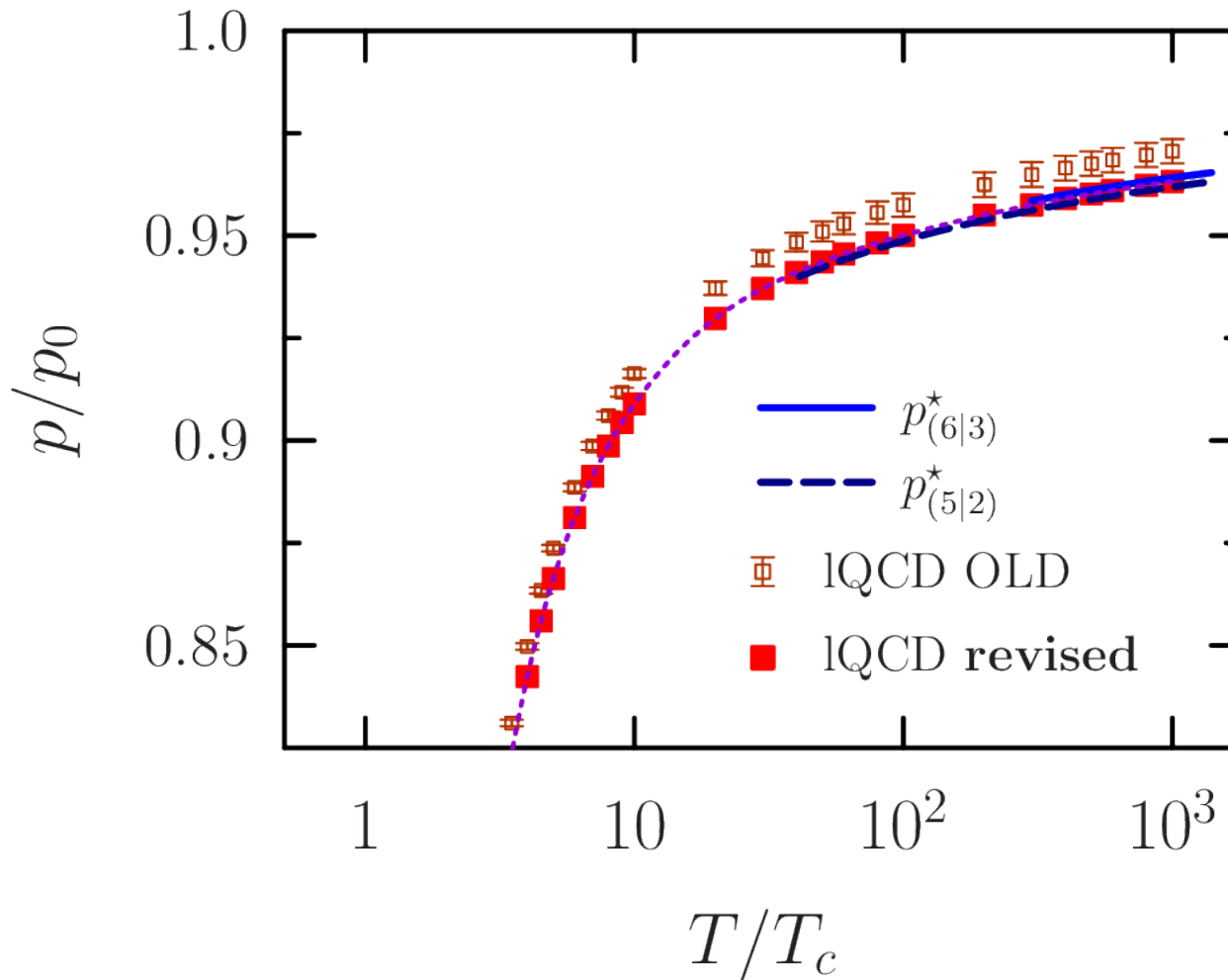
matching
interaction measure at large T



“re-calibration” of
pressure at $T > 4T_c$

slower approach to free limit

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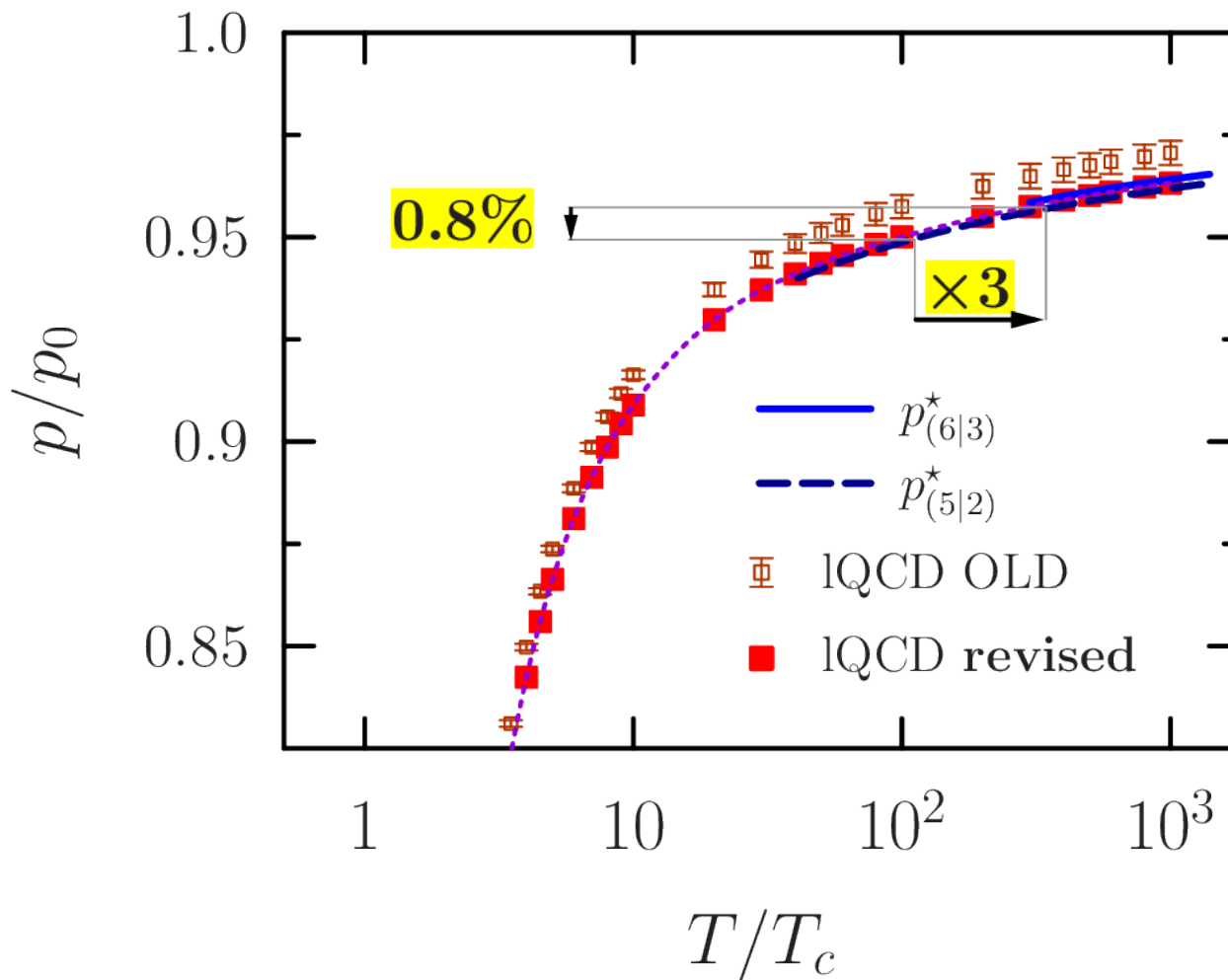
non-perturbative coefficient

$$c_6 = \mathcal{O}(-40)$$

$$c_6 = -72 \pm 3$$

$$c_6 = -95 \pm 6$$

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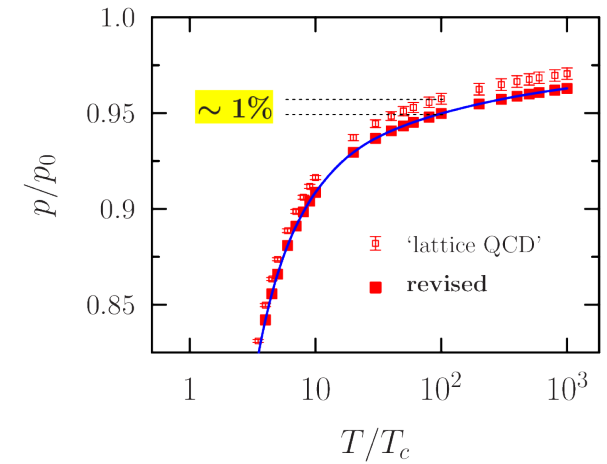
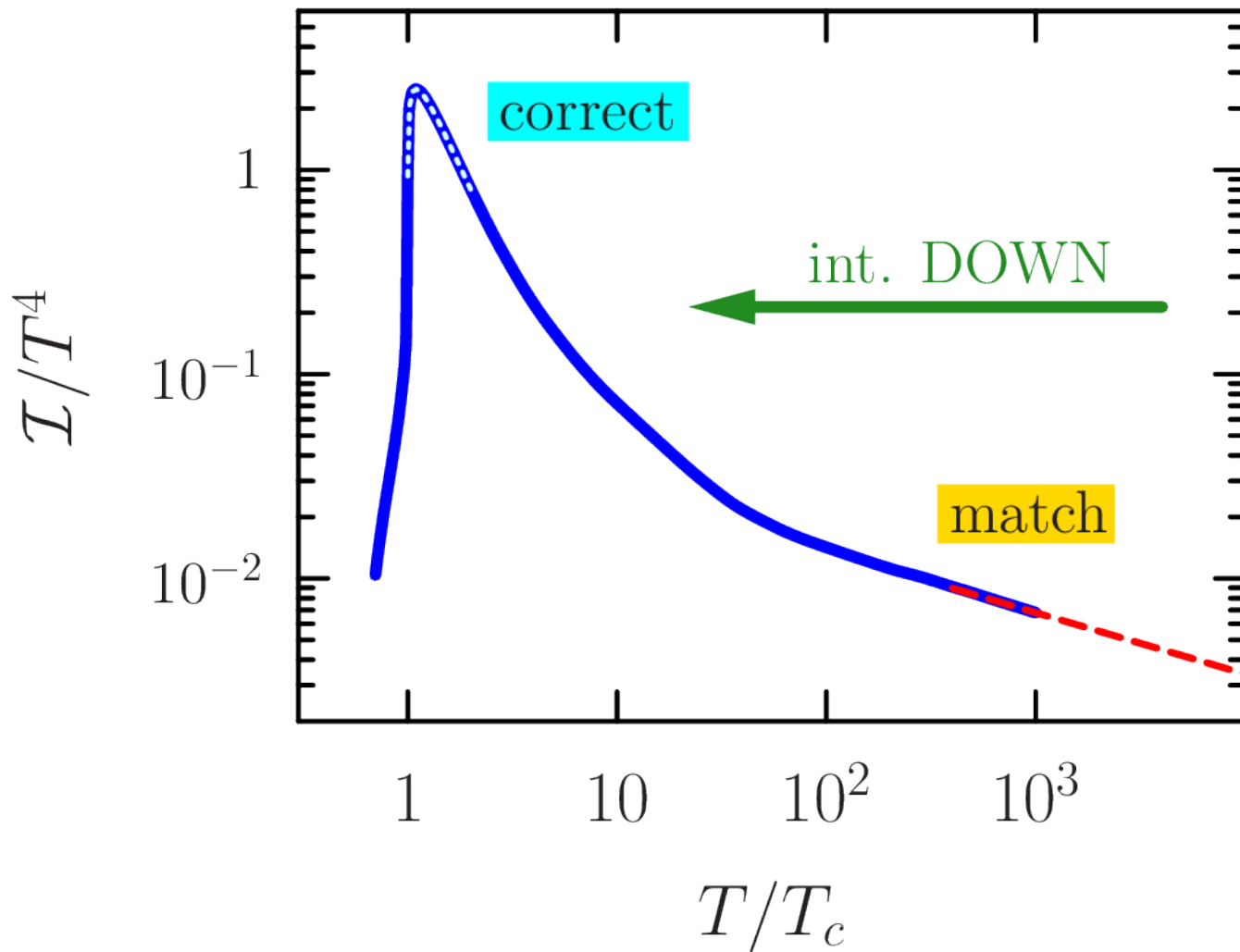
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- **outlook: phenomenological implications for physical case**