Cosmological Constant Problem: Deflation During Inflation B. Koch with F. Canales, C. Laporte, & A. Rincon, based on: JCAP no 1, 21, 2020, ArXiv:1812.10526



Frankfurt, Nuclear Physics Colloquium

16.01.2020

acknowledge: VRI, Fondecyt



Collaboration

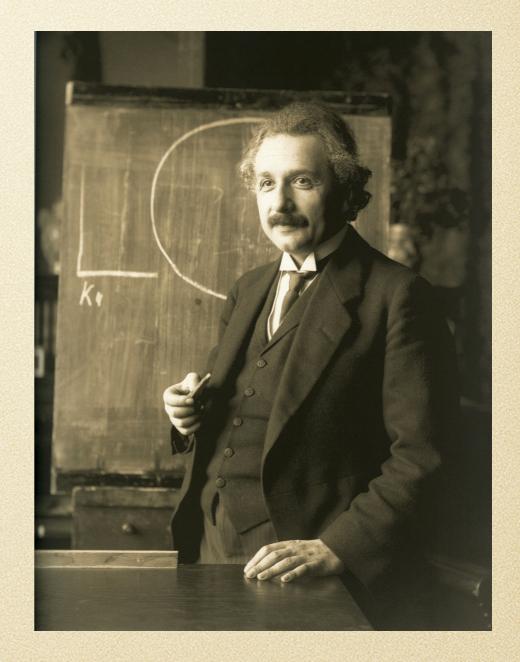
• Angel Rincon, Cristobal Laporte, & Felipe Canales



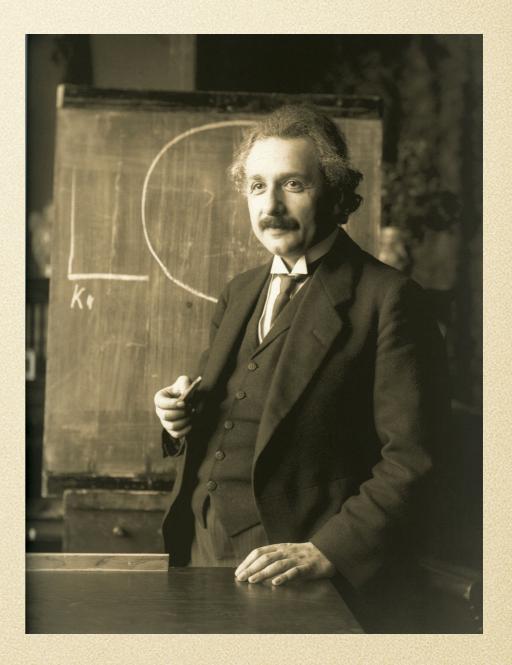
René Araneda

Content

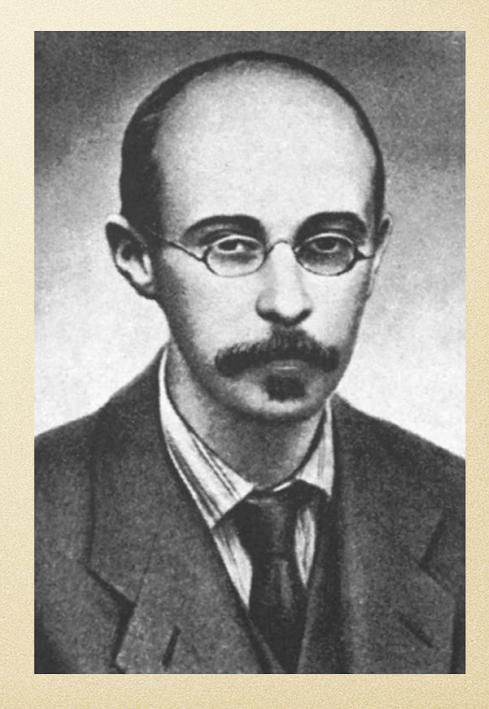
- Cosmological constant problem, status
- Conceptual problem, in evolving Universe
- Scale dependent framework & evolving Universe
- Possible solution: Deflation during inflation
- Link to Asymptotic Safety
- Conclusion



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi G T_{\mu\nu}$$

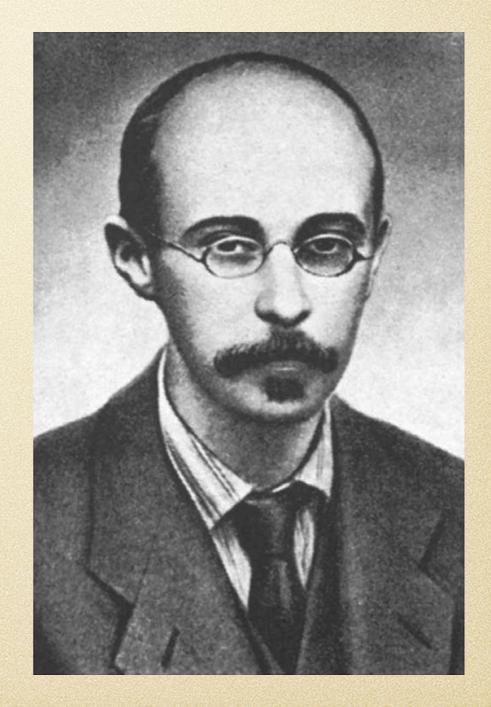


Alexander Friedmann

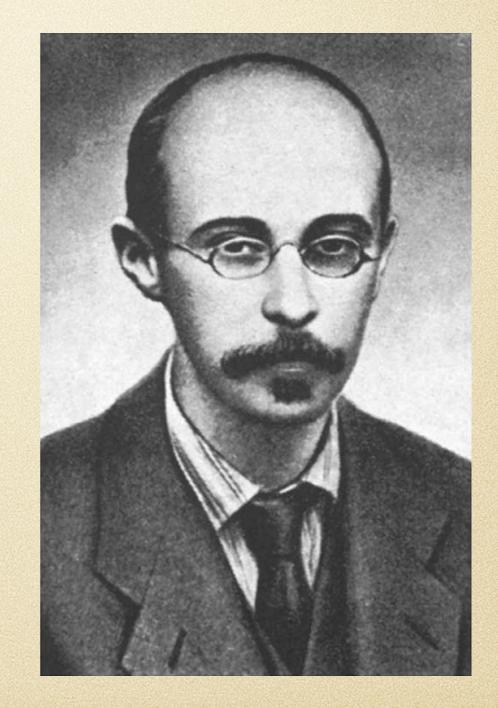


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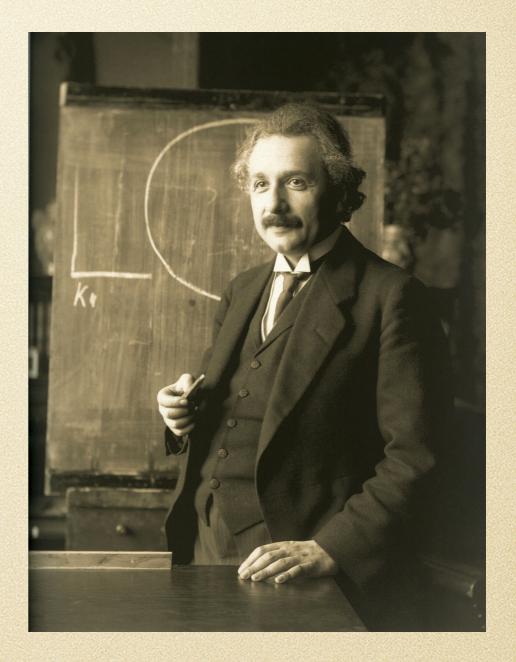
 $ds^2 = a(t)ds_3^2 - dt^2$



Alexander Friedmann $ds^2 = a(t)ds_3^2 - dt^2$ $\frac{\dot{a}^2 + k}{a^3} = \frac{1}{3} 8\pi G\rho$ $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$

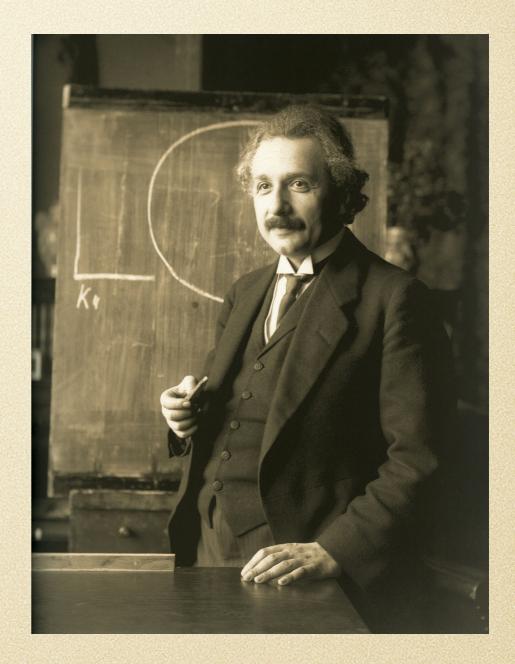


 $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi G T_{\mu\nu}$



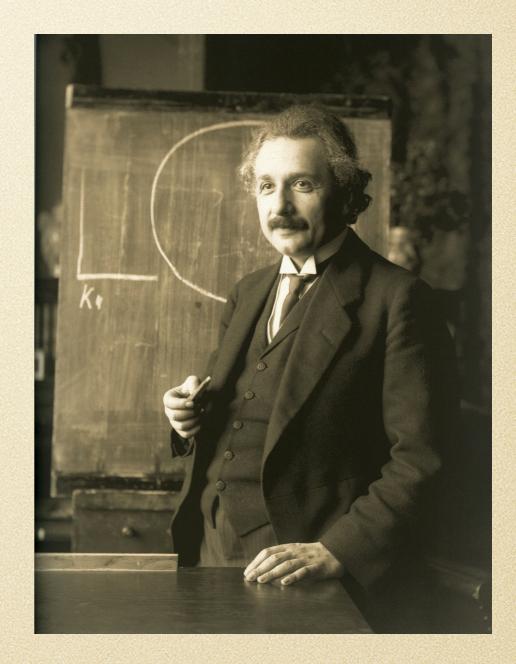
$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) \neq 0$$



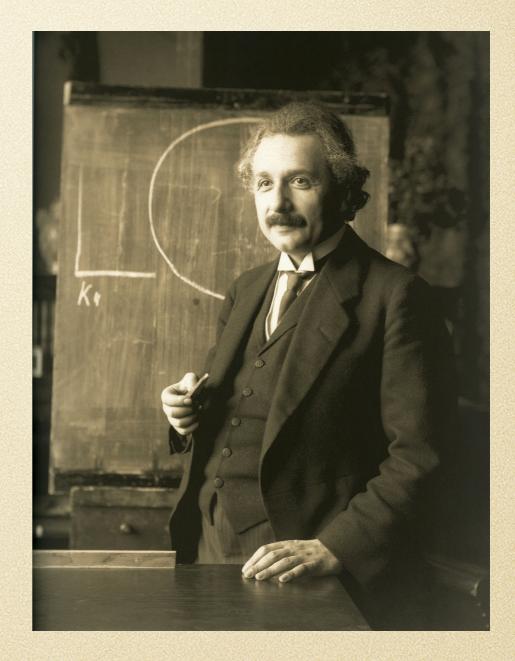
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$$\Rightarrow \dot{a} \neq 0$$



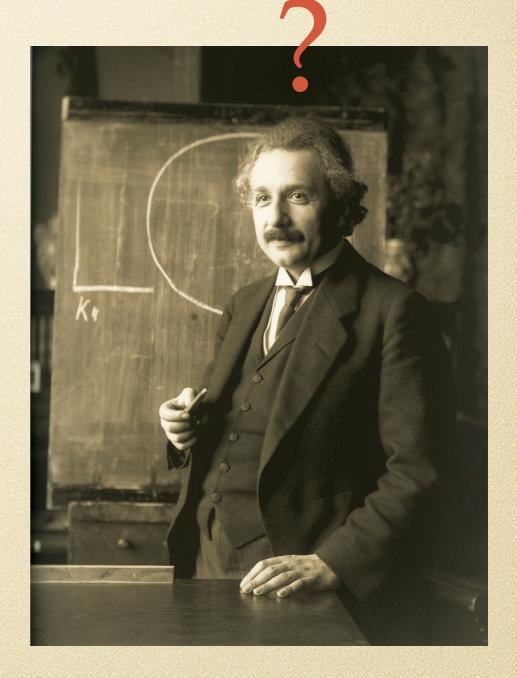
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$$\Rightarrow \dot{a} \neq 0 \text{ not static}$$



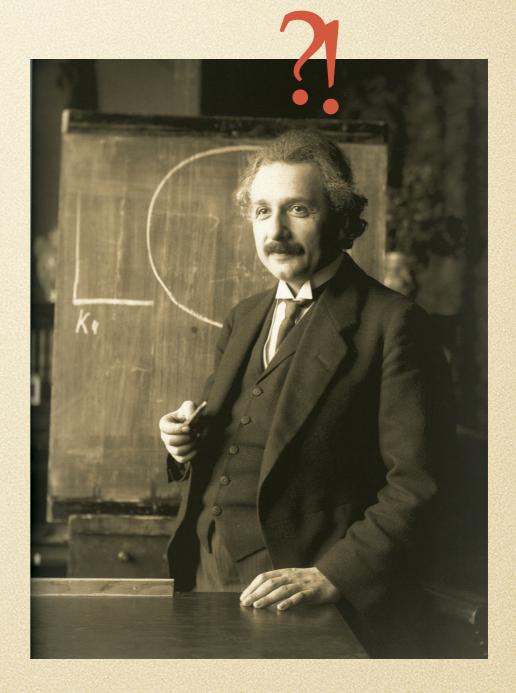
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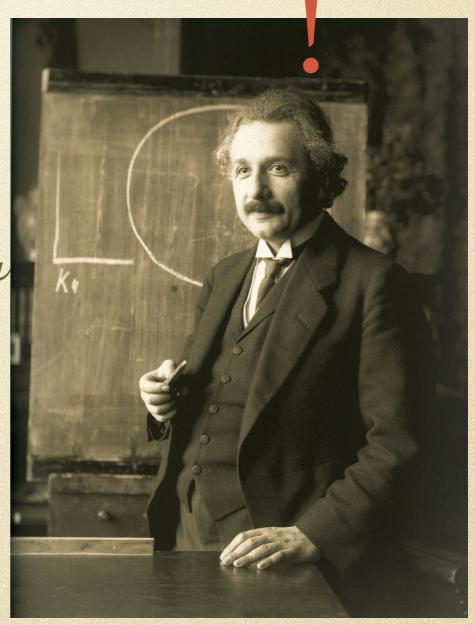


$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

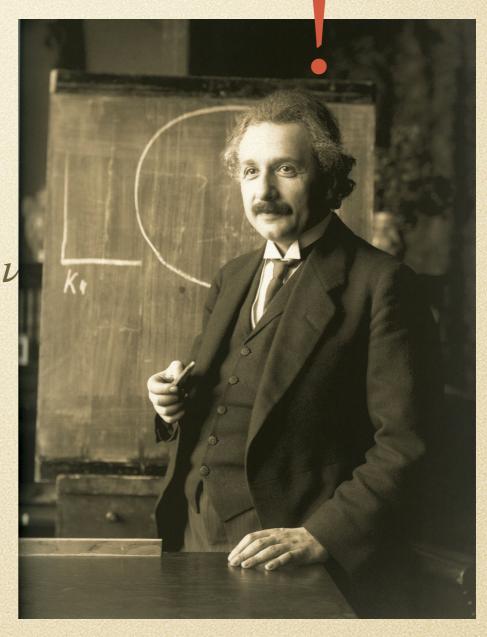
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 $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu}$$
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0$$
static possible $\dot{a} \equiv 0$



Edwin Hubble, Georges Lemaítre

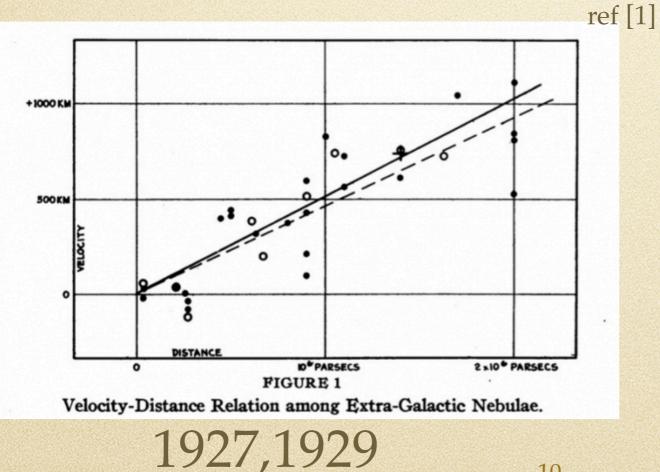


Edwin Hubble, Georges Lemaítre measurement:



Edwin Hubble, Georges Lemaítre

measurement:





Edwin Hubble

measurement:

 $\dot{a} > 0$



Edwin Hubble

measurement:

 $\dot{a} > 0$

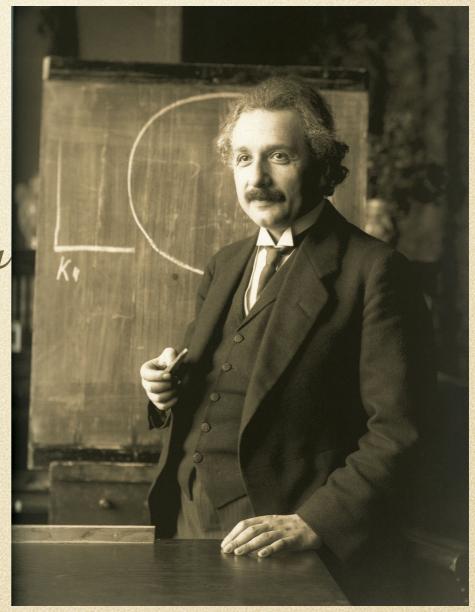
later:

 $\dot{a} = 67.66 \pm 0.42 \frac{km/s}{Mpc}$ (Planck collaboration 2018)

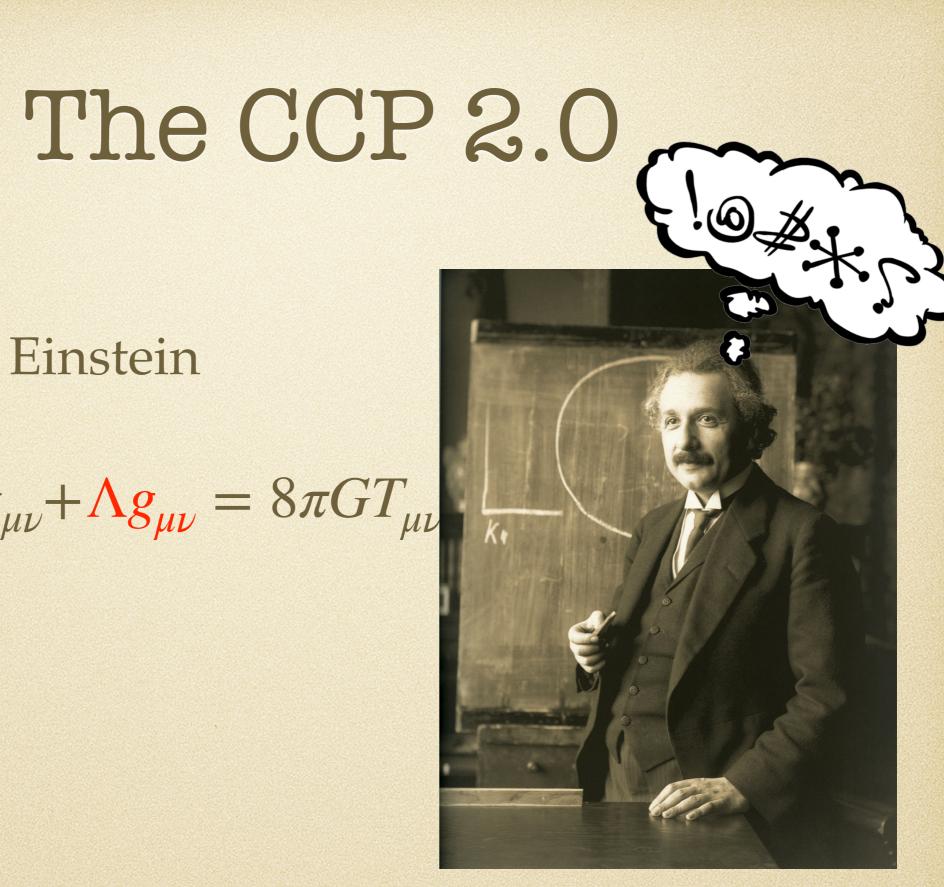


Edwin Hubble measurement: not static $\dot{a} > 0$ later: $\dot{a} = 67.66 \pm 0.42 \frac{km/s}{m}$ *Mpc* (Planck collaboration 2018)

 $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$

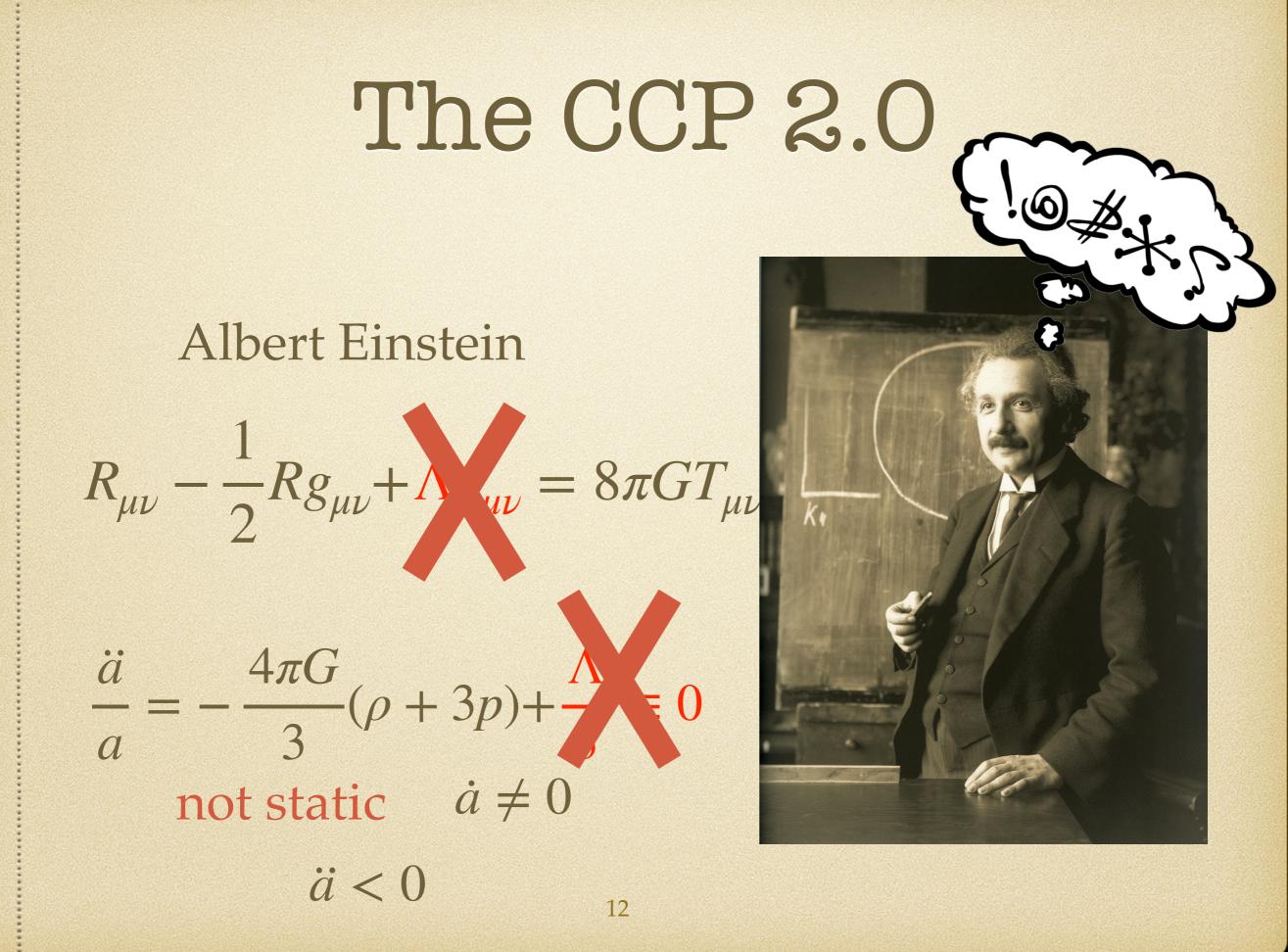


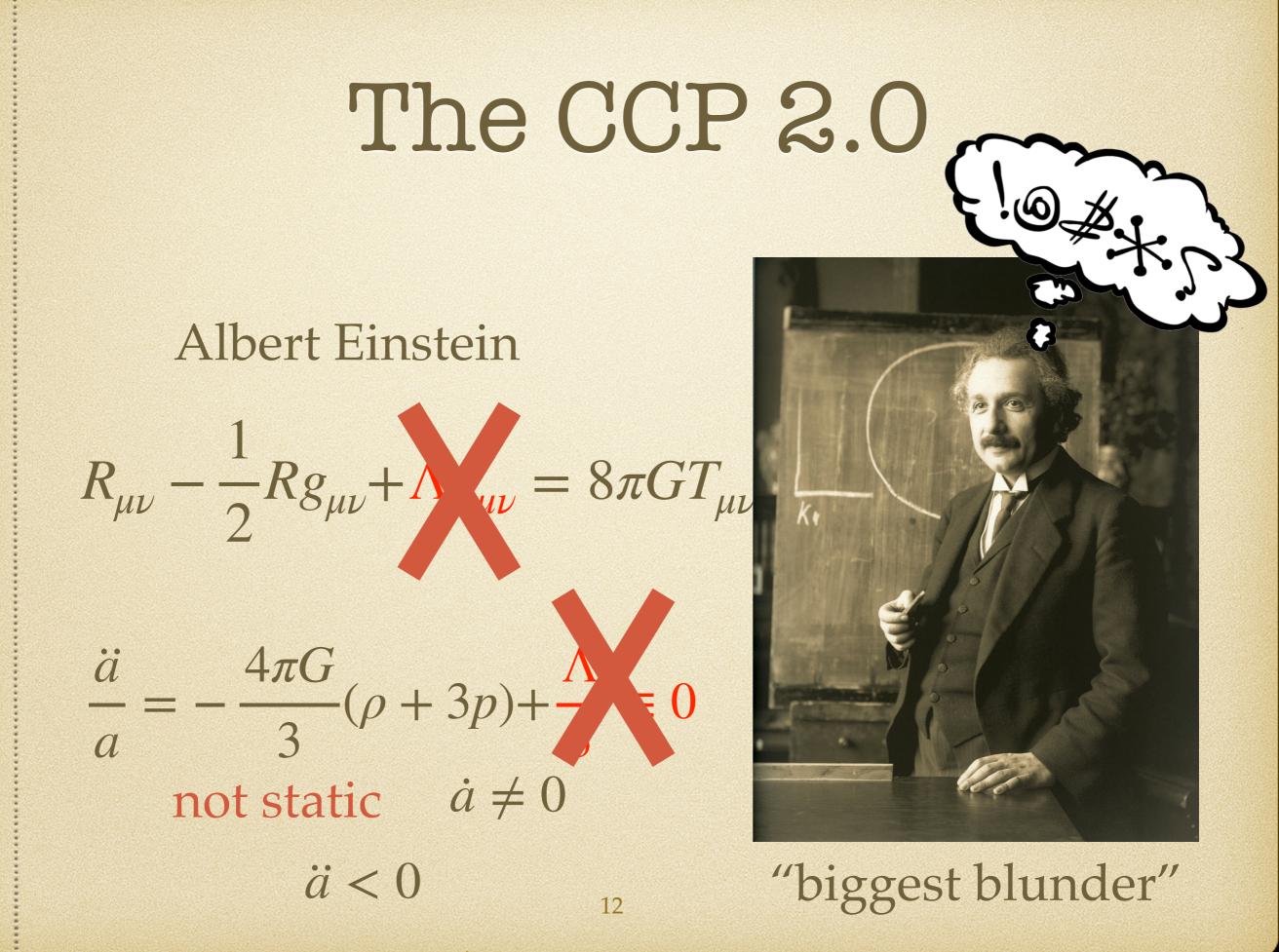
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12

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda_{\mu\nu} = 8\pi G T_{\mu\nu}$$





S. Perlmutter, A. Riess, B. Schmidt, & others

ref [2]



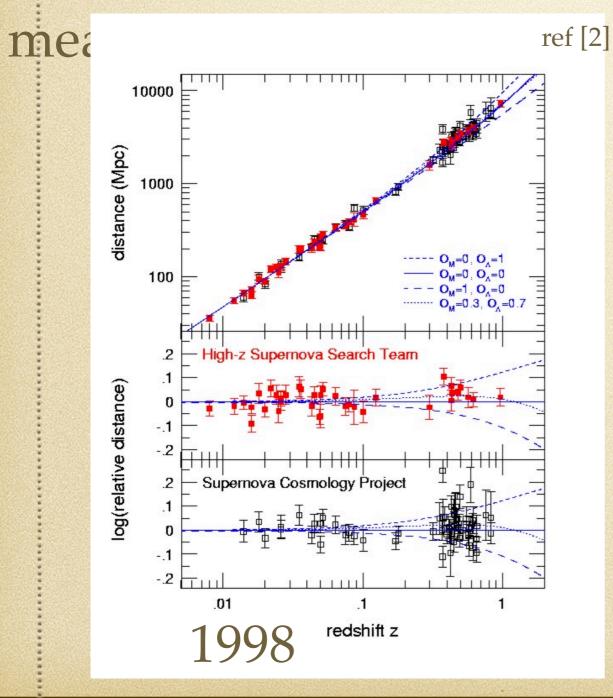
S. Perlmutter, A. Riess, B. Schmidt, & others

measurements: ref [2]



14

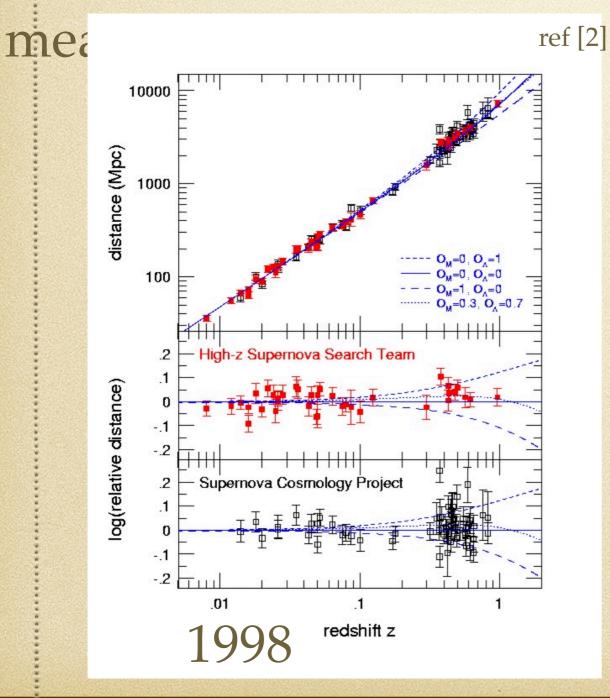
S. Perlmutter, A. Riess, B. Schmidt, & others





14

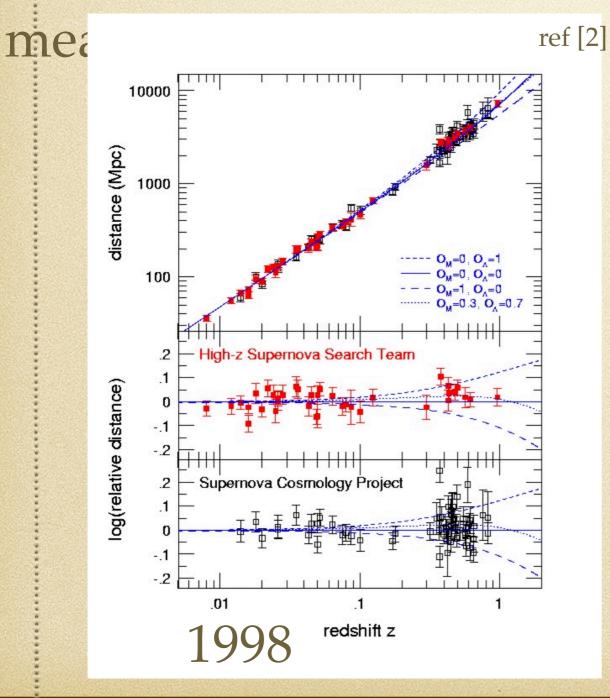
S. Perlmutter, A. Riess, B. Schmidt, & others





 $\dot{a} \neq 0 \qquad \ddot{a} < 0$

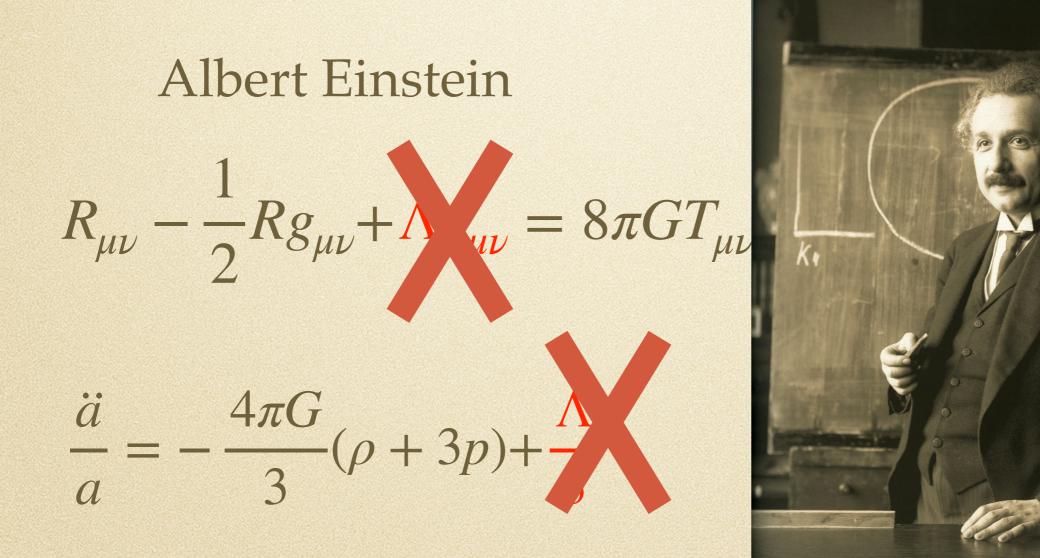
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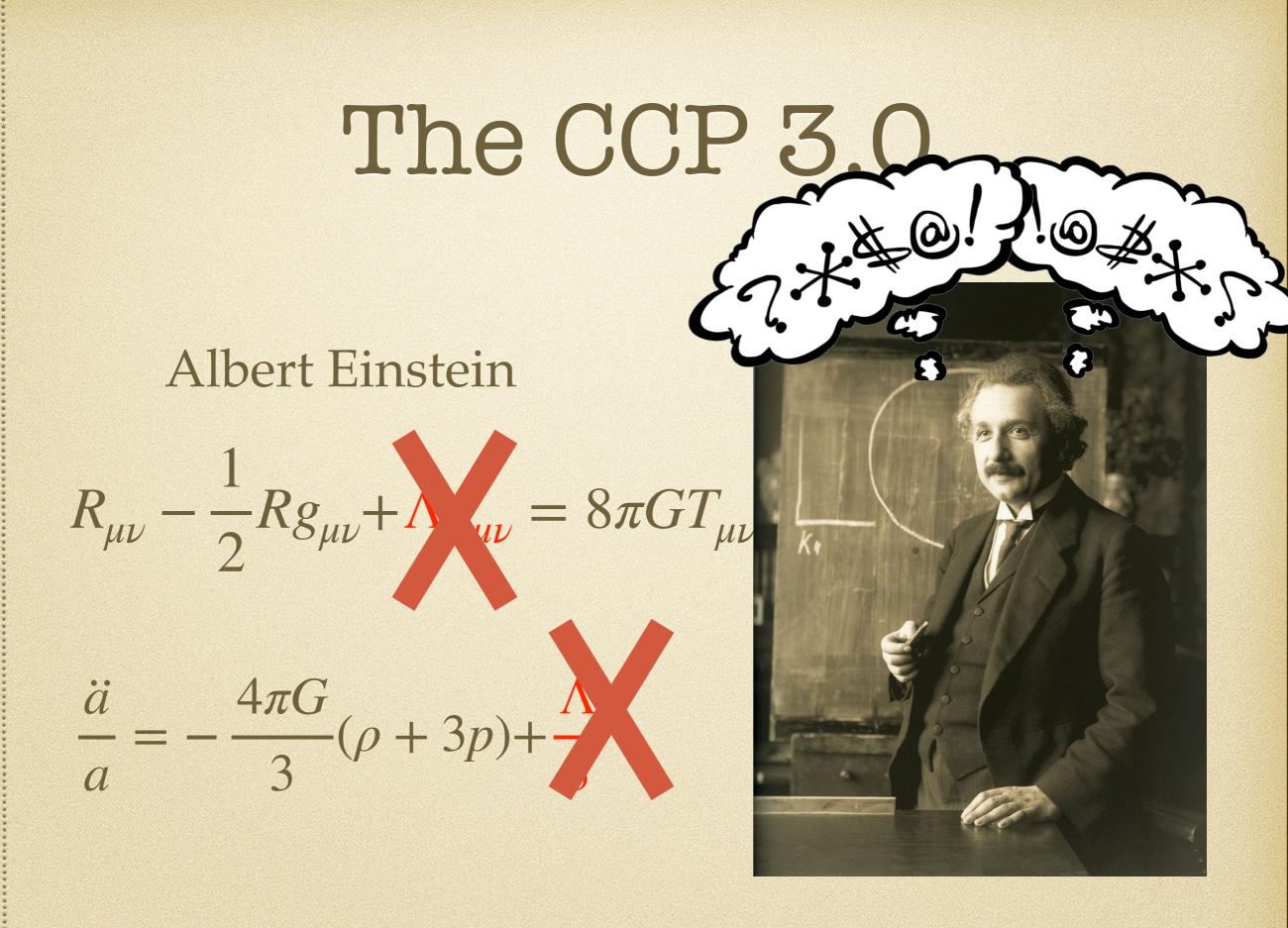




 $\dot{a} \neq 0$ $\ddot{a} > 0$

14

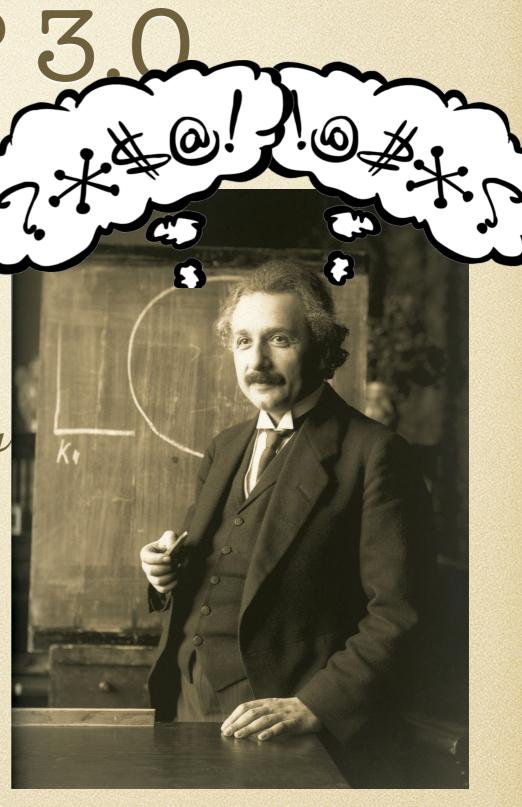




Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

 $-\frac{4\pi G}{3}(\rho+3p)+\frac{\Lambda}{3}$ ä a



<u>k</u>@!<u>7</u>!@_

E

Albert Einstein

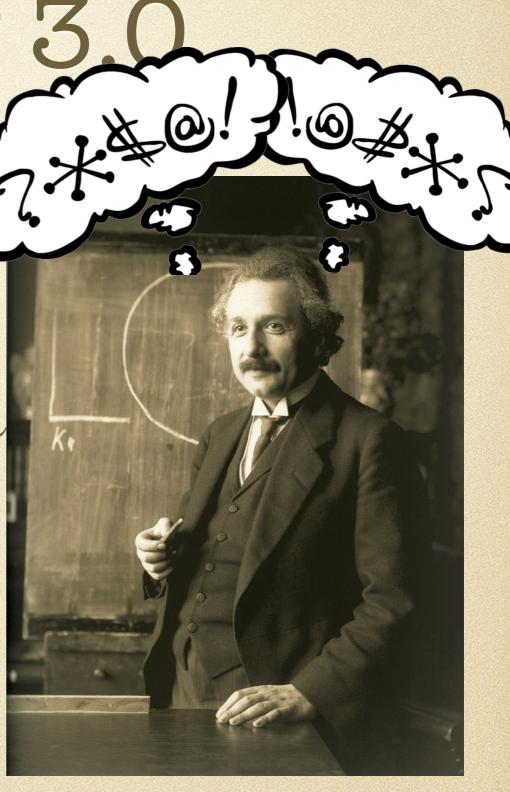
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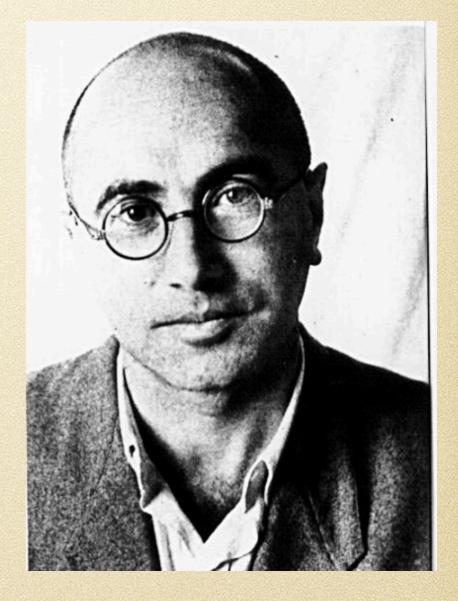
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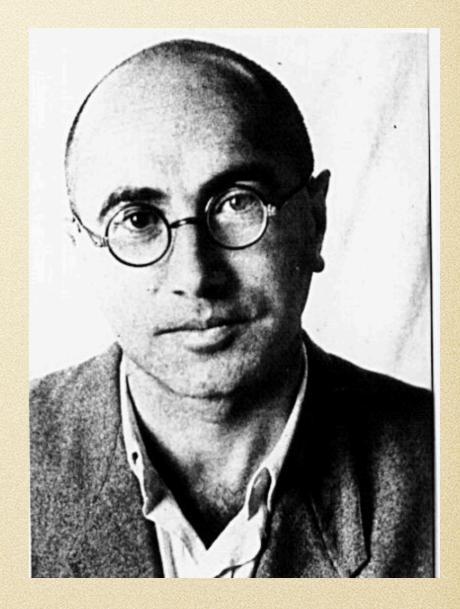
Yakov Zeldovich



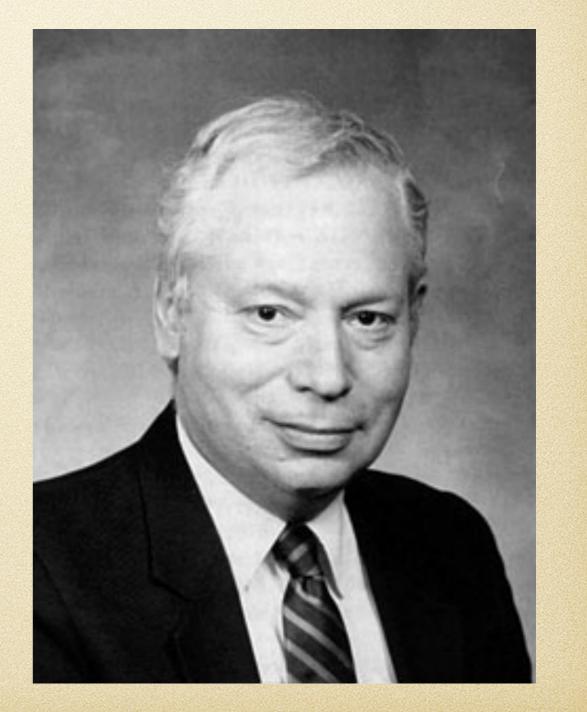
Yakov Zeldovich

Quantum fluctuations predict value of Λ

1967

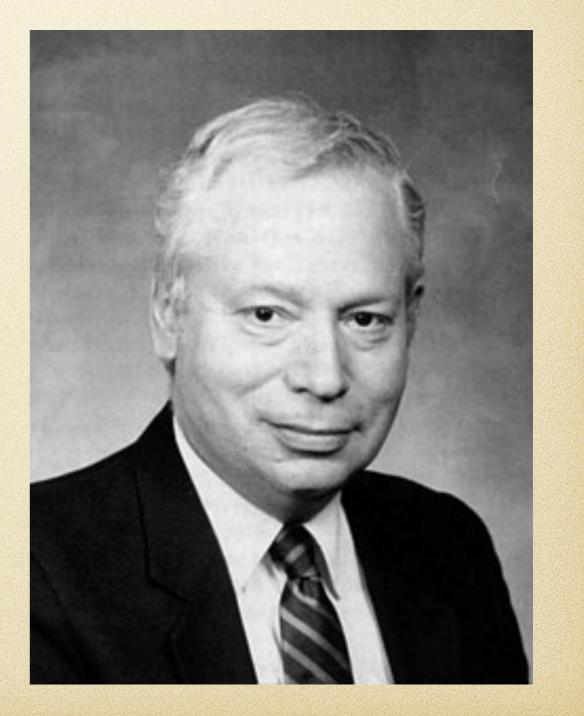


Steven Weinberg



Steven Weinberg

Quantum fluctuations predict value of Λ ref [3] Problem since 1998



Quantum fluctuations predict value of Λ

Quantum fluctuations predict value of Λ

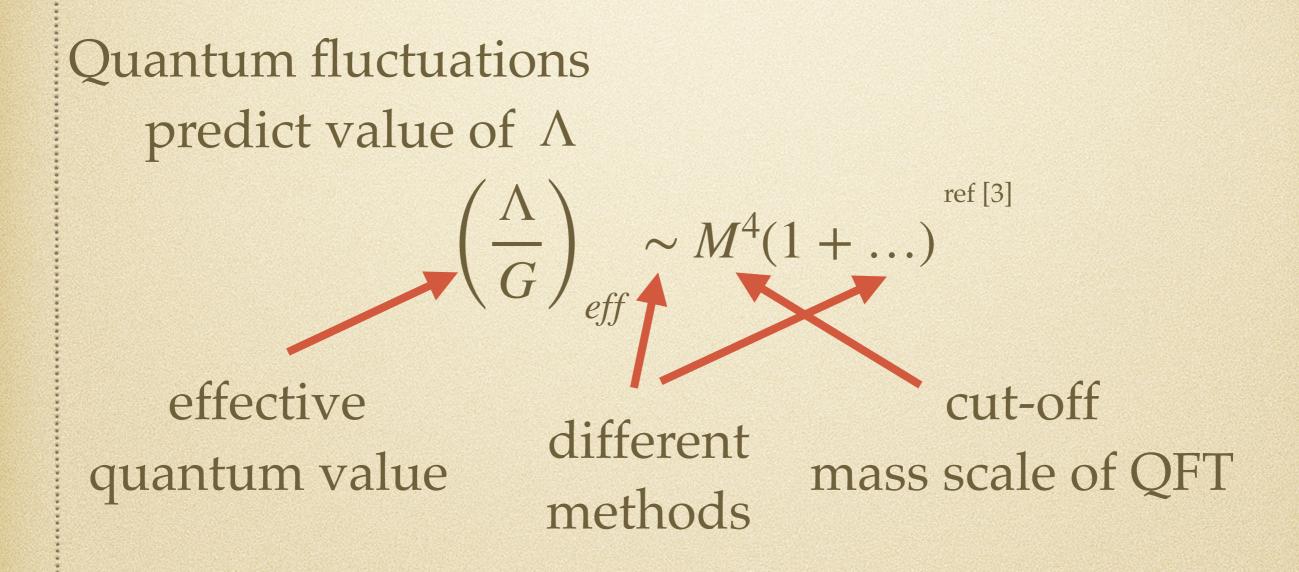
 $\left(\frac{\Lambda}{G}\right)_{eff} \sim M^4(1+\ldots)^{\text{ref [3]}}$

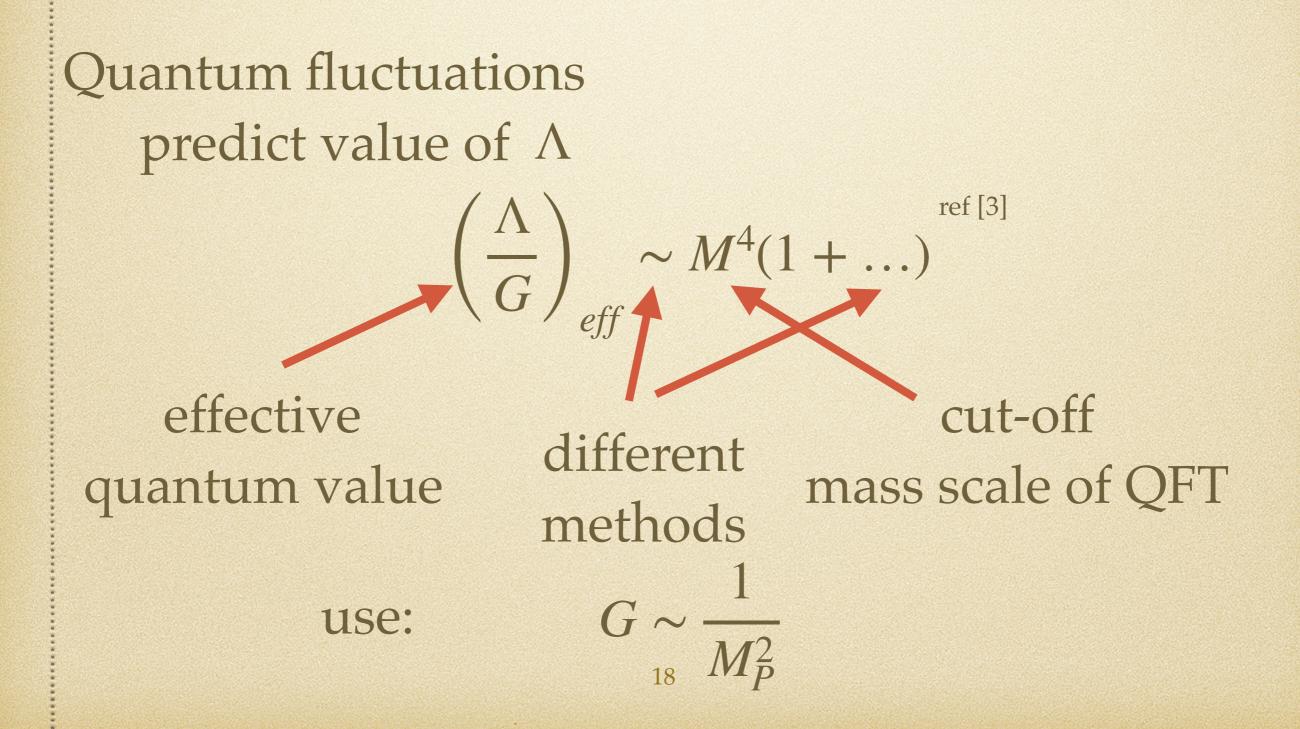
Quantum fluctuations predict value of Λ $\left(\frac{\Lambda}{G}\right)_{eff} \sim M^4(1+...)^{\text{ref [3]}}$

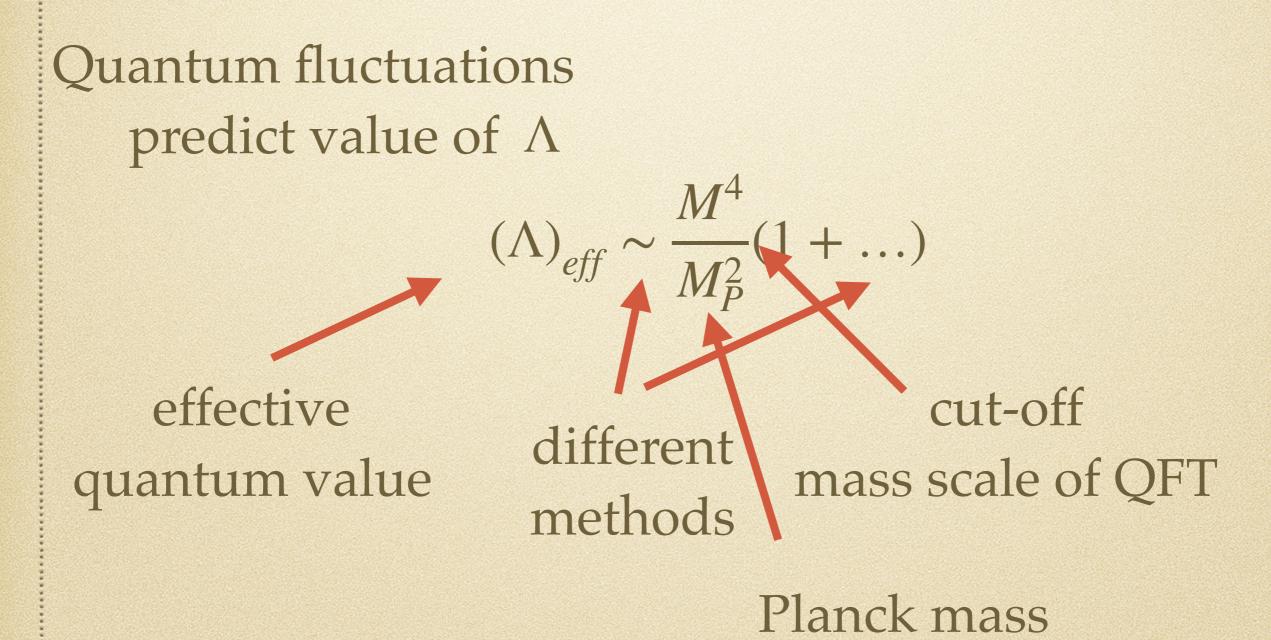
effective quantum value

Quantum fluctuations predict value of Λ $\left(\frac{\Lambda}{G}\right)_{eff} \sim M^4(1+...)^{ref [3]}$

effective quantum value cut-off mass scale of QFT







Quantum fluctuations predict value of Λ

$$(\Lambda)_{eff} \sim \frac{M^4}{M_P^2} (1 + \dots)$$

Quantum fluctuations predict value of Λ

Highest physical mass scale

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Quantum fluctuations predict value of Λ

Highest physical mass scale

Observed value

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Highest physical mass scale

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Observed value

$$\Lambda_o = \frac{\rho_c}{M_P^2} \approx \frac{10^{-47} GeV^4}{M_P^2}$$

Quantum fluctuations predict value of Λ

Highest physical mass scale

 $(\Lambda)_{eff} \sim \frac{M^4}{M_P^2} (1 + \ldots)$

Observed value

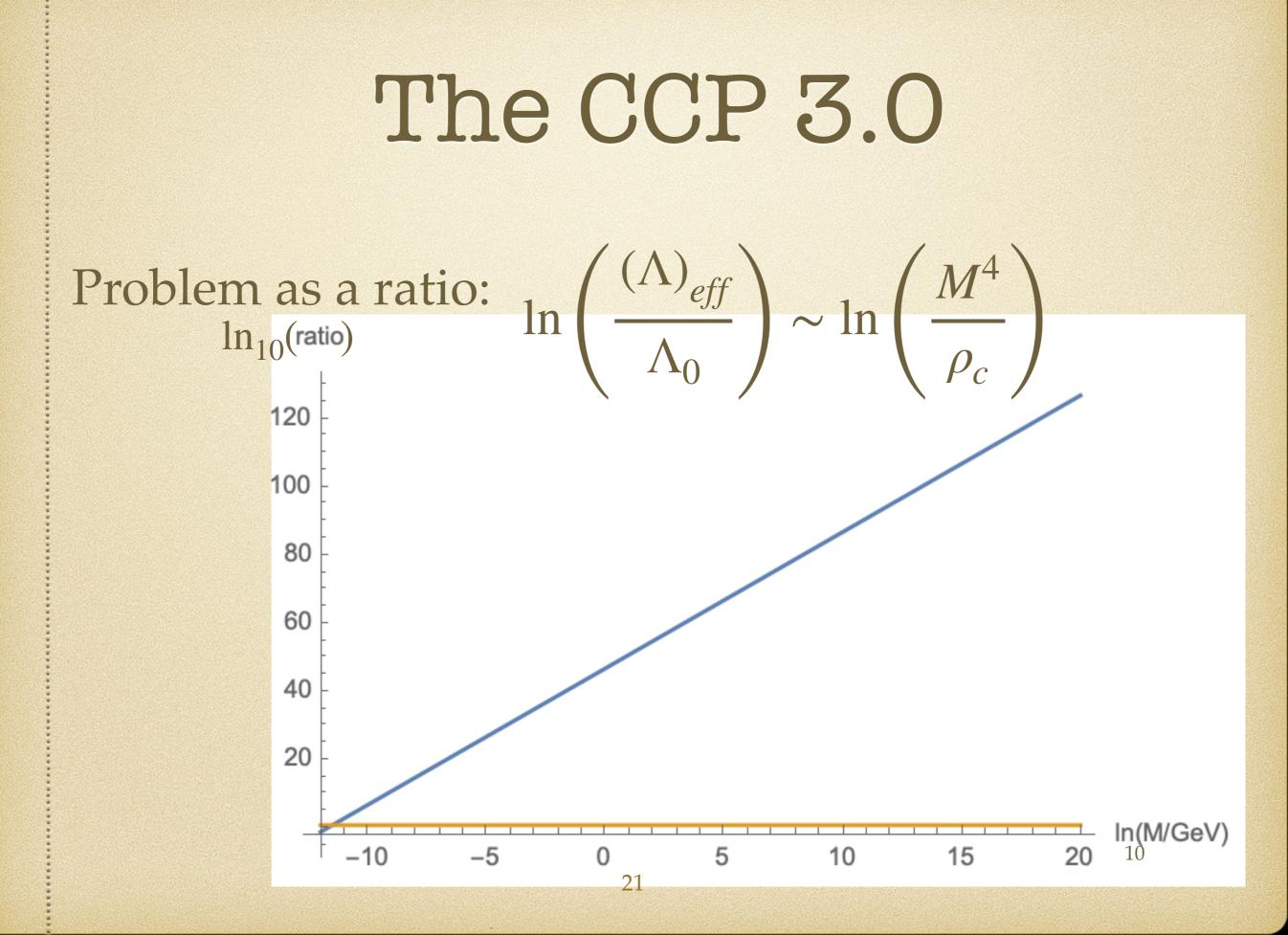
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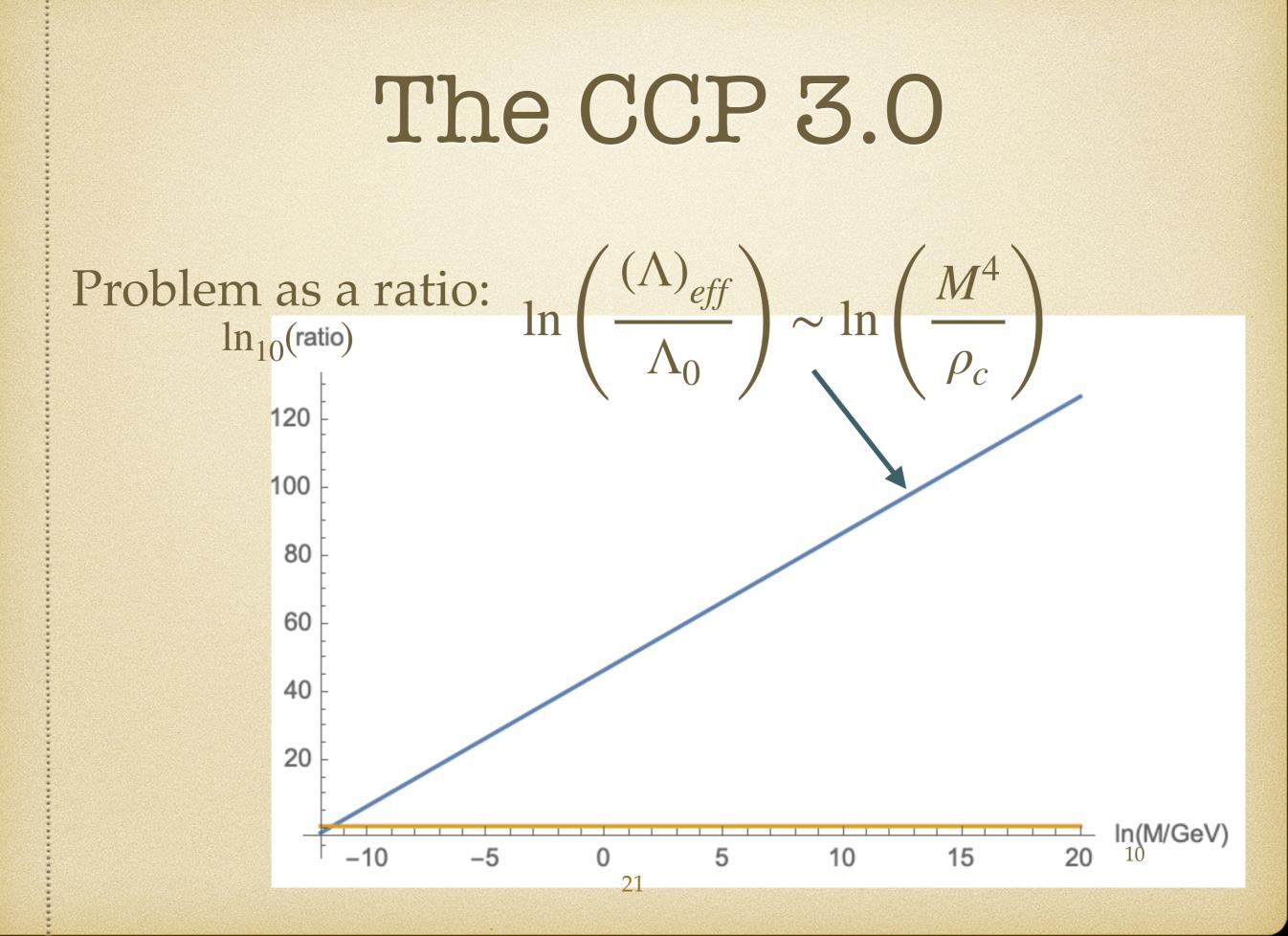
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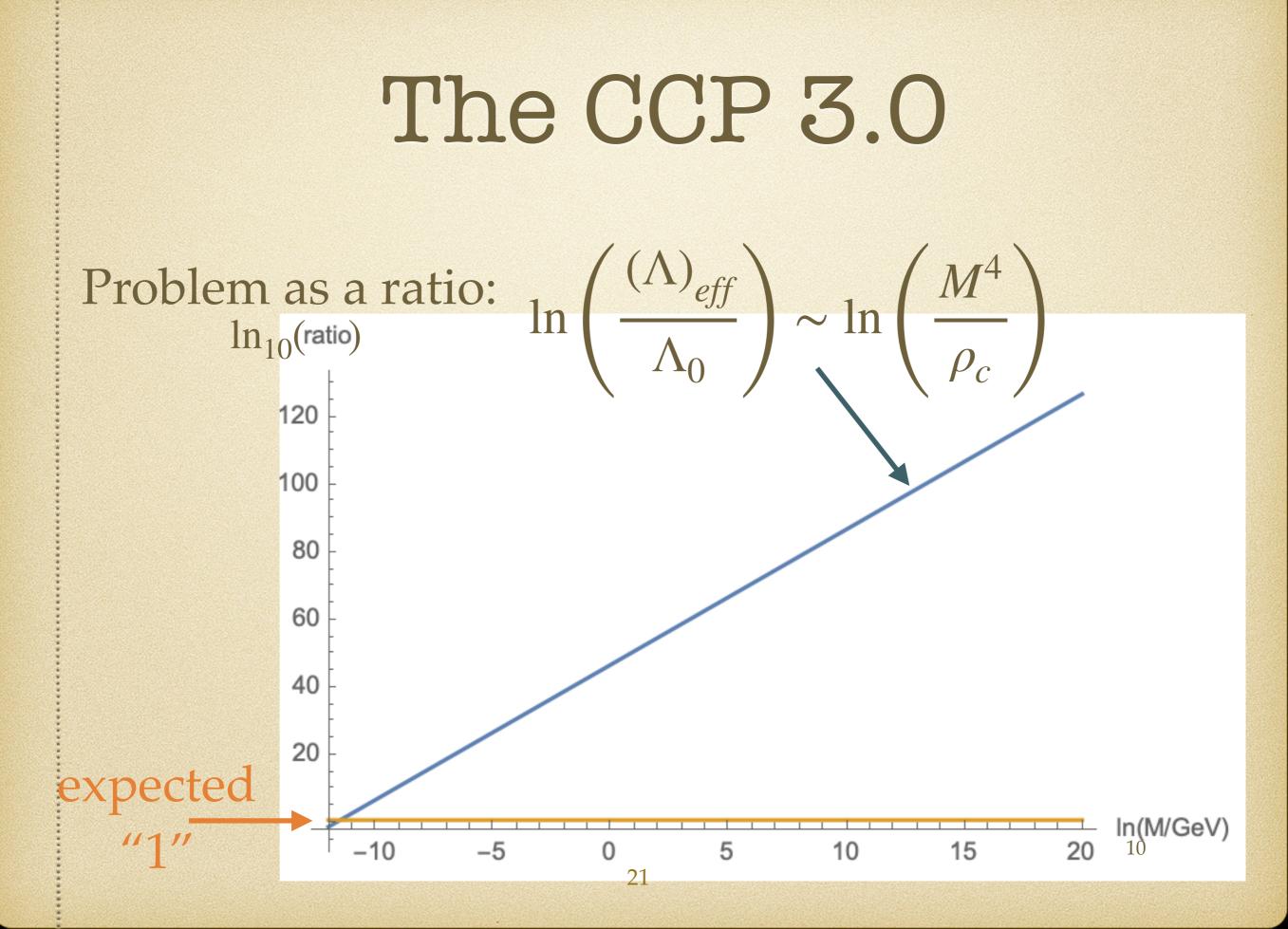
observed critical energy density

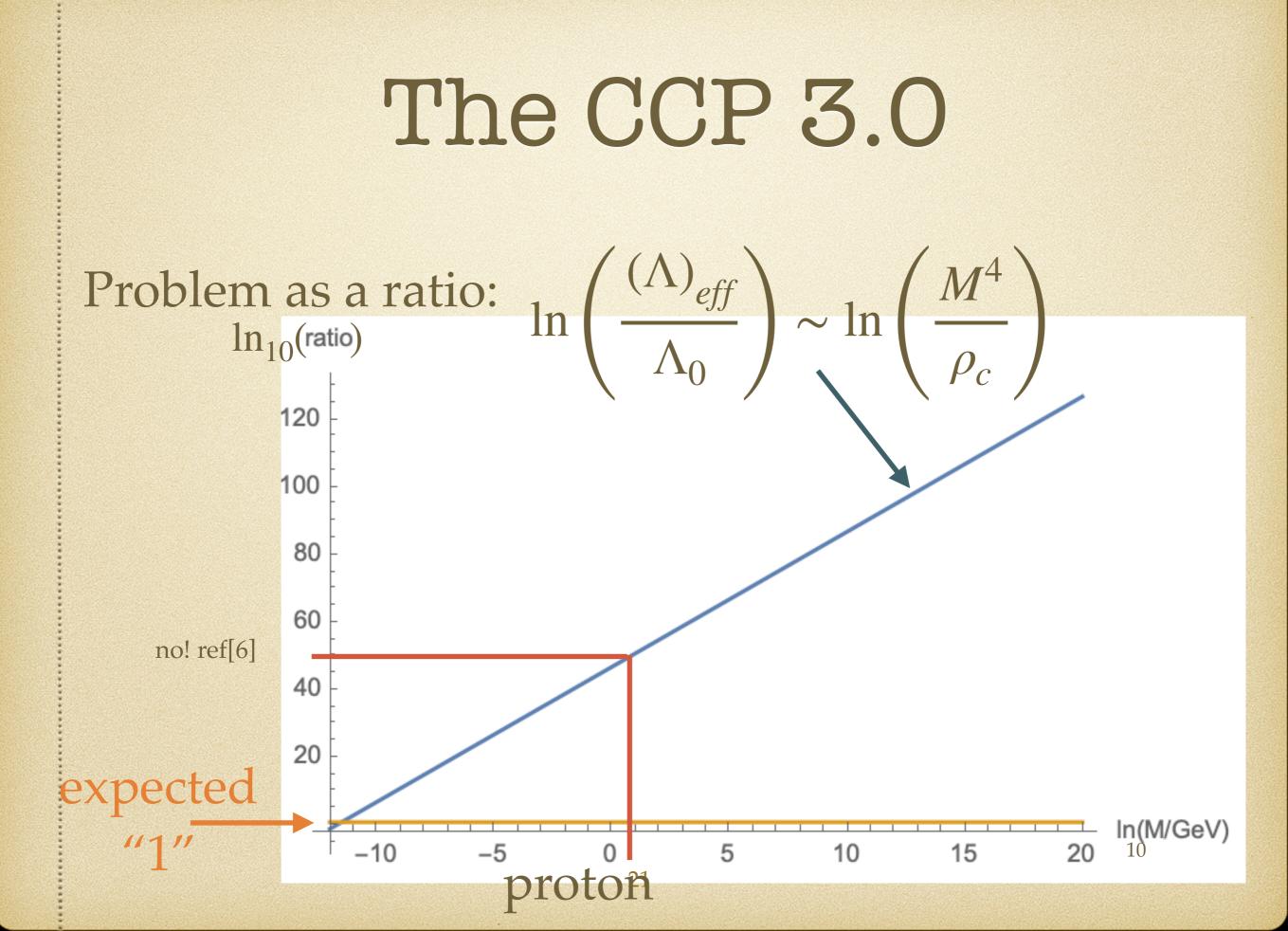
Problem as a ratio:

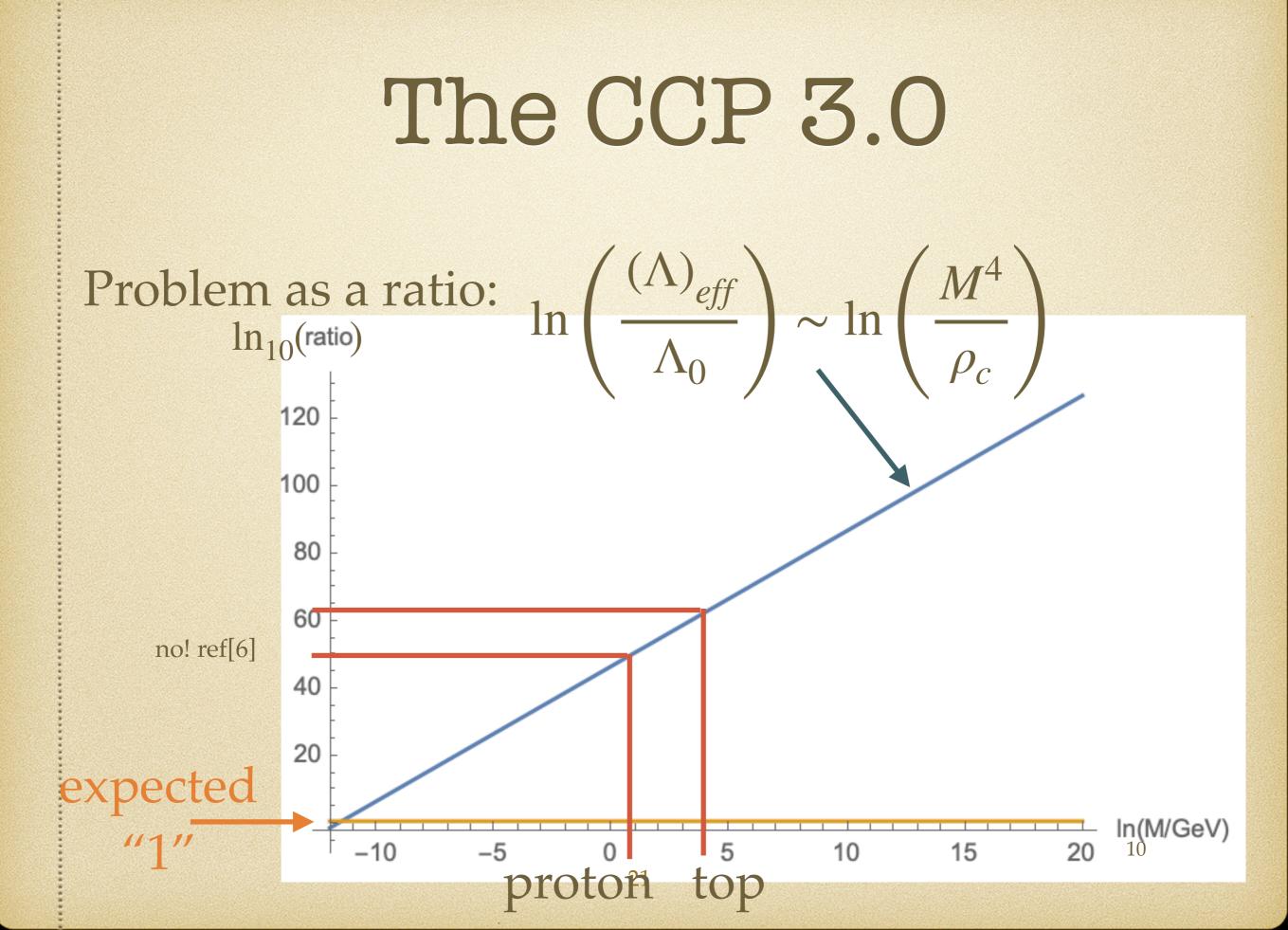
Problem as a ratio: $\ln\left(\frac{(\Lambda)_{eff}}{\Lambda_0}\right) \sim \ln\left(\frac{M^4}{\rho_c}\right)$

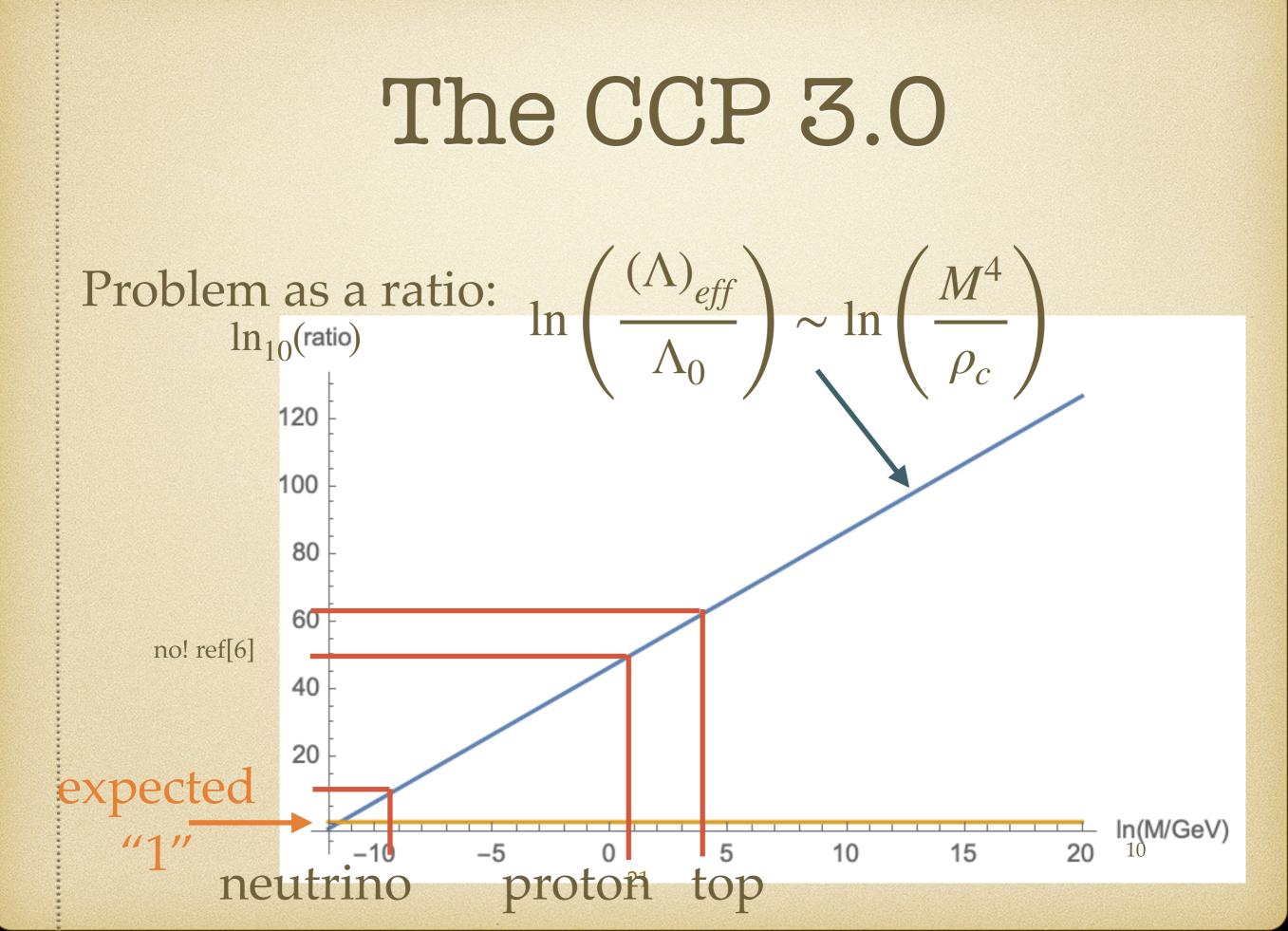


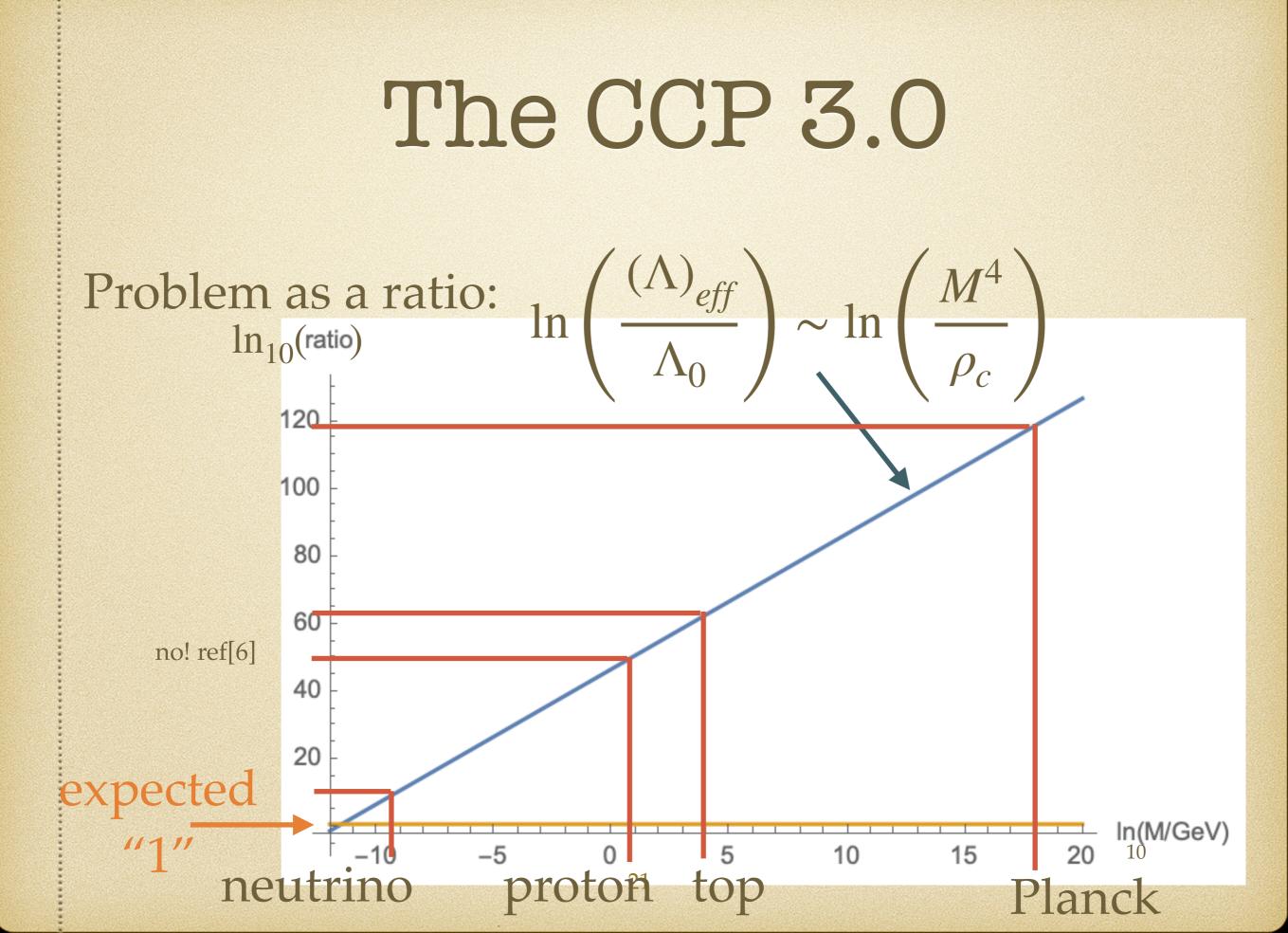


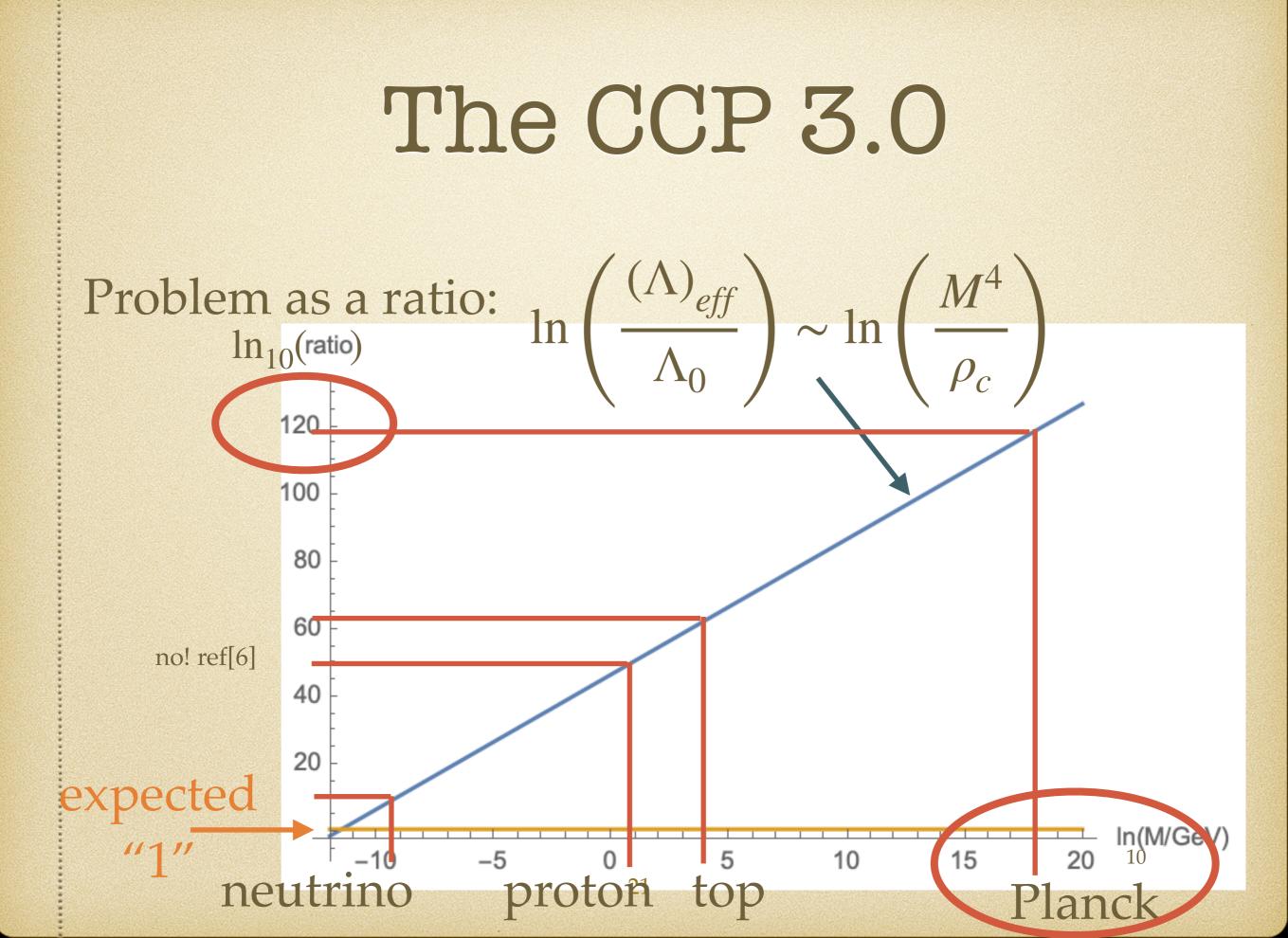












Problem as a ratio:

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$$\frac{(\Lambda)_{eff}}{\Lambda_0} \sim \frac{1}{G_N \cdot \Lambda_0} \sim \frac{M_P^4}{\rho_c} \approx 10^{120}$$

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we try to address this problem

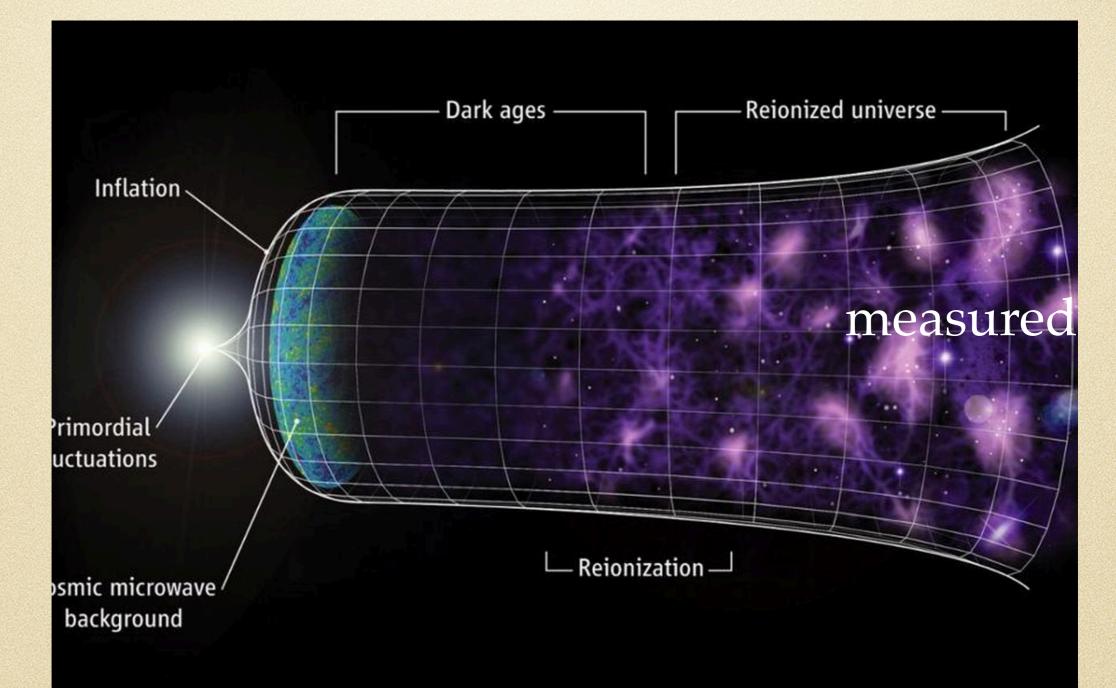
Problem as a ratio:

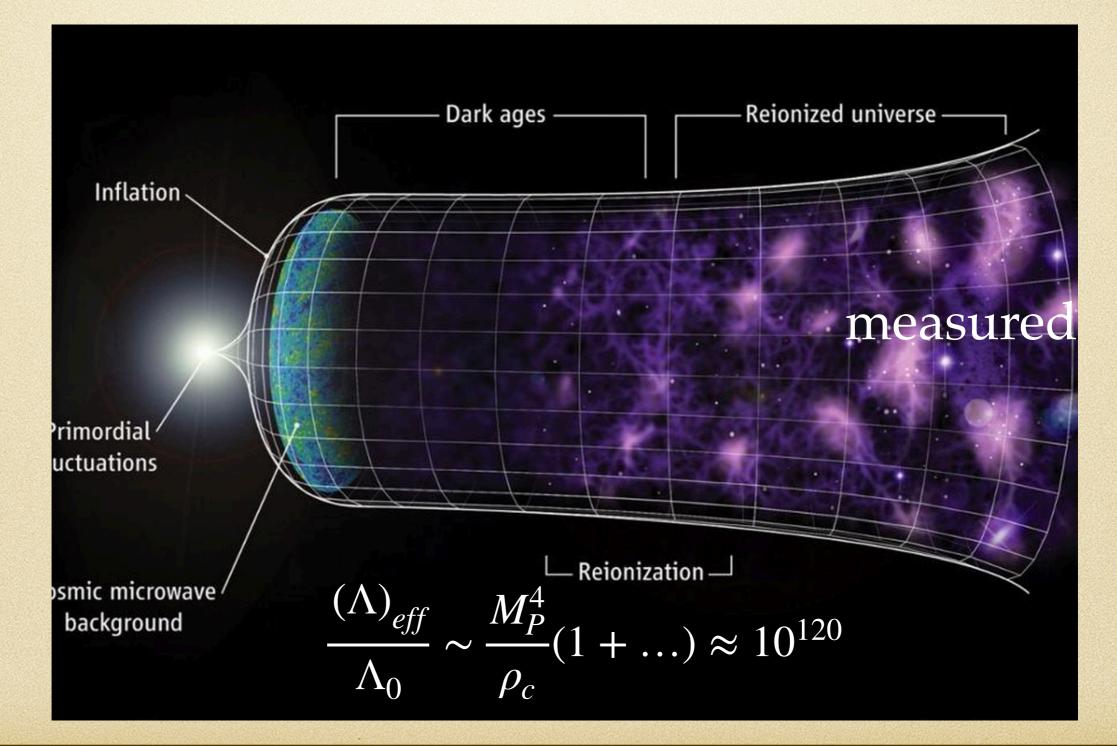
$$\frac{(\Lambda)_{eff}}{\Lambda_0} \sim \frac{1}{G_N \cdot \Lambda_0} \sim \frac{M_P^4}{\rho_c} \approx 10^{120}$$

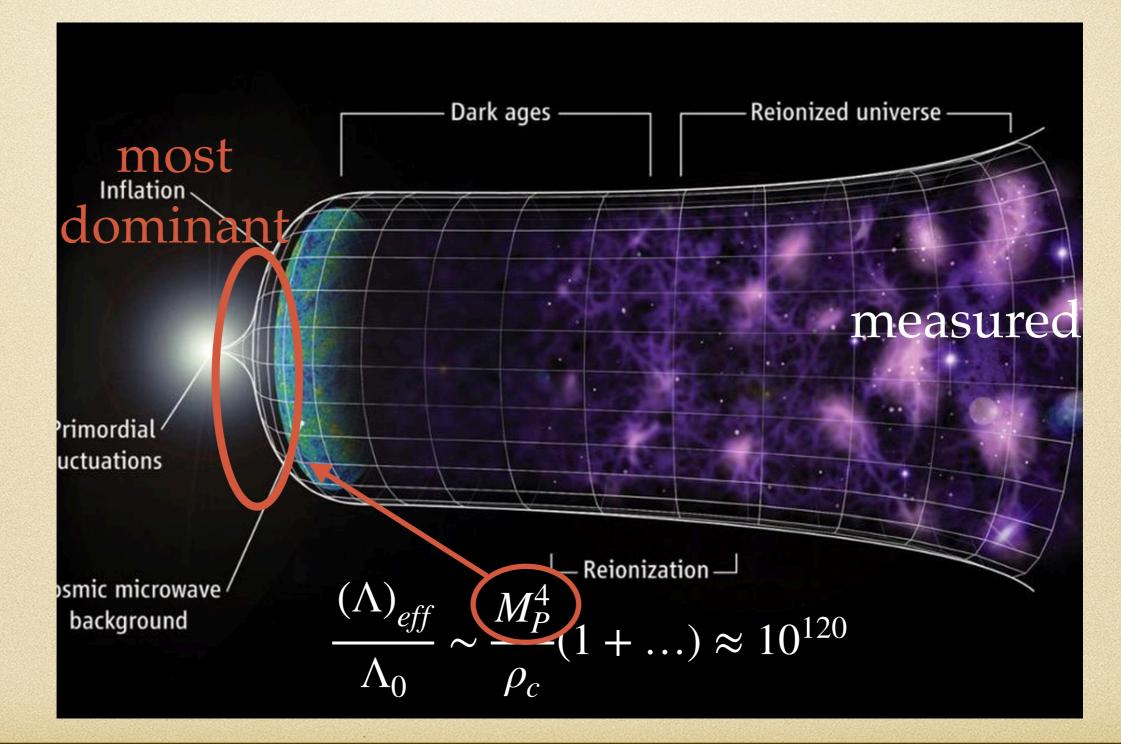
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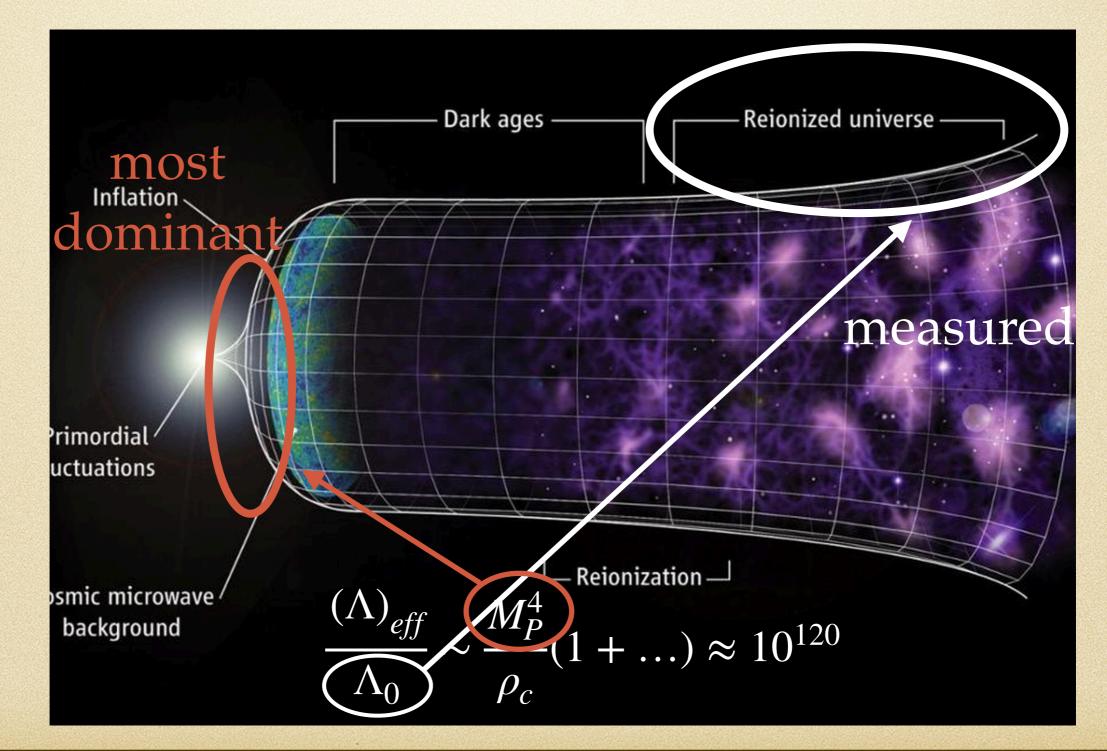
assuming there are quantum fluctuations of gravity associated to the Planck scale

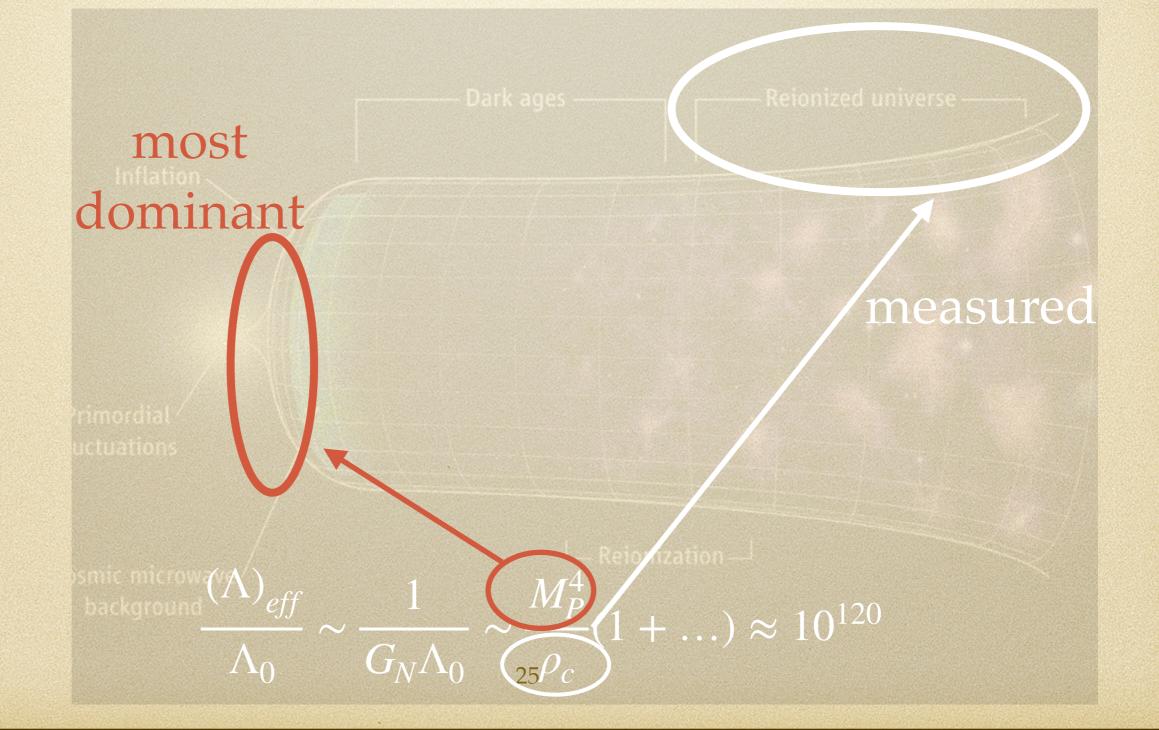
Evolving Universe Issue

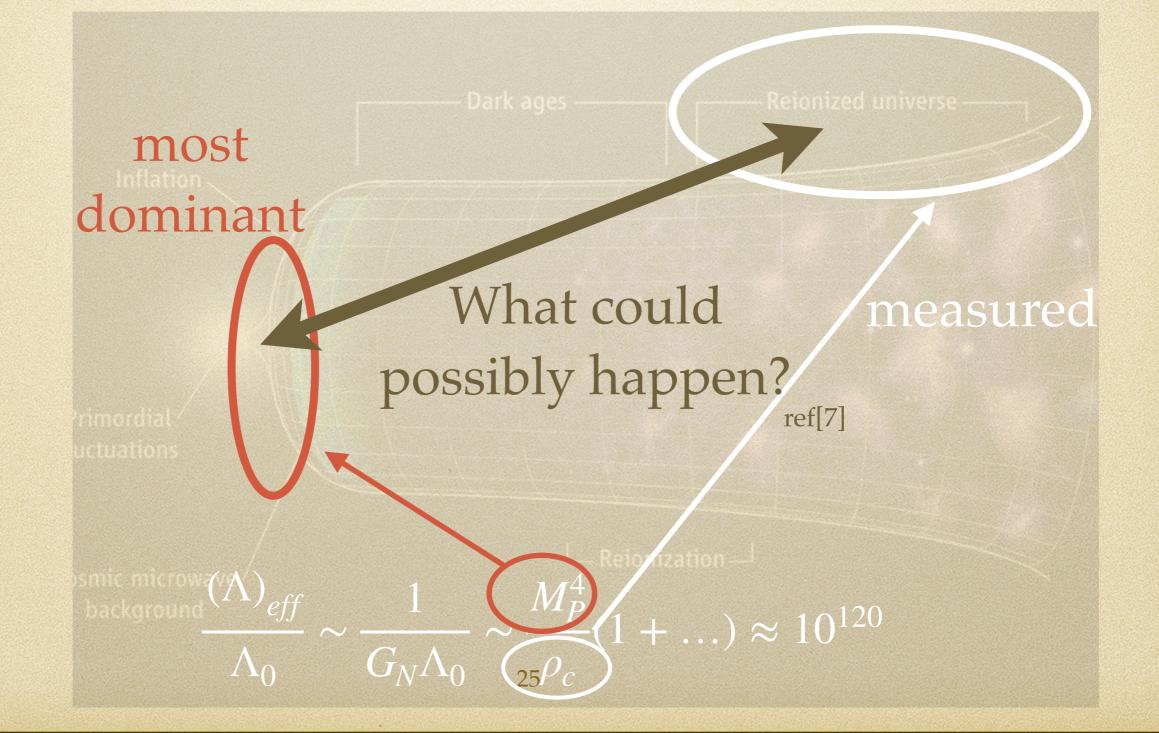












Gravity as classical theory

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$$S = \int d^4x \sqrt{-g} \left(\frac{R}{G_N} - 2\frac{\Lambda_0}{G_N} \right)$$

Gravity as classical theory

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Gravity as effective QFT

Gravity as classical theory

$$S = \int d^4x \sqrt{-g} \left(\frac{R}{G_N} - 2\frac{\Lambda_0}{G_N}\right)$$

Gravity as effective QFT

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

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Non renormalizable?

Gravity as effective QFT

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

Non renormalizable? Yes, but ...

Gravity as effective QFT

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

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Yes, but ...



Gravity as effective QFT

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

Non renormalizable? Yes, but ... Could still be predictive QFT (Asymptotic Safety)



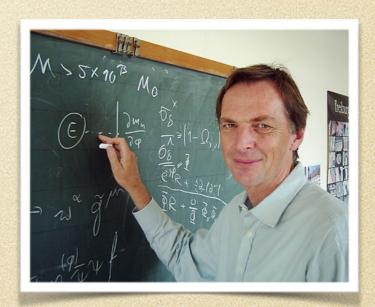
Asymptotic Safety in a nutshell

• Idea: works if non trivial UV-fixed points for finite number of couplings (S.W)



Asymptotic Safety
in a nutshell
$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2 \frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

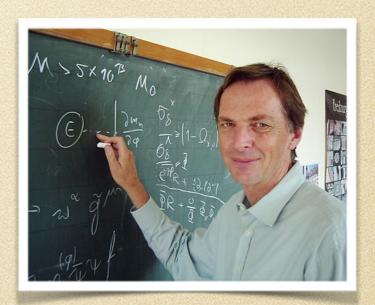
Tool: Functional renormalization group equation



Asymptotic Safety
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$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2 \frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

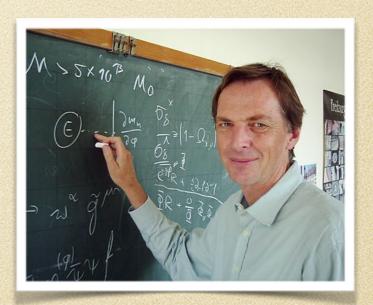
Tool: Functional renormalization group equation

$$\partial_k \Gamma_k = \frac{1}{2} Tr \left(\frac{\partial_k R_k}{\Gamma_k^{(2)} + R_k} \right)$$



Tool: Functional renormalization group equation

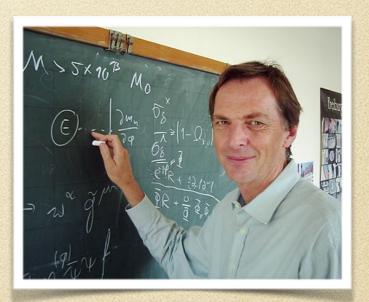
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Tool: Functional renormalization group equation

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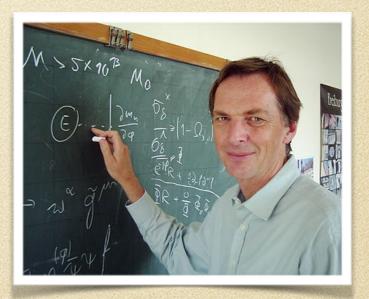
two point
function



Tool: Functional renormalization group equation

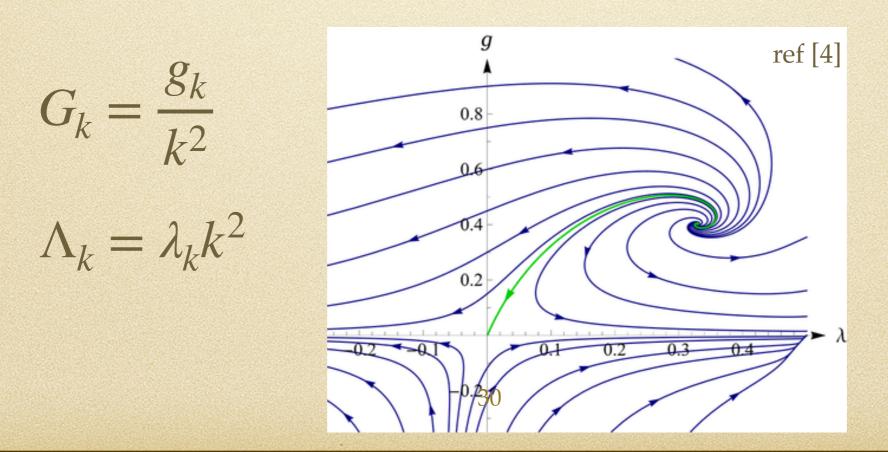
$$\partial_k \Gamma_k = \frac{1}{2} Tr \left(\frac{\partial_k R_k}{\Gamma_k^{(2)} + R_k} \right)$$

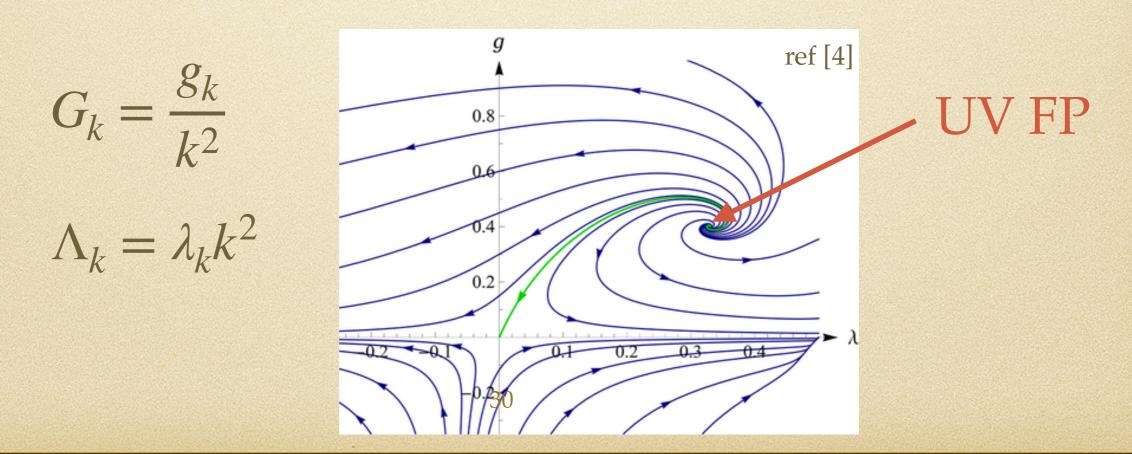
two point regulator
function

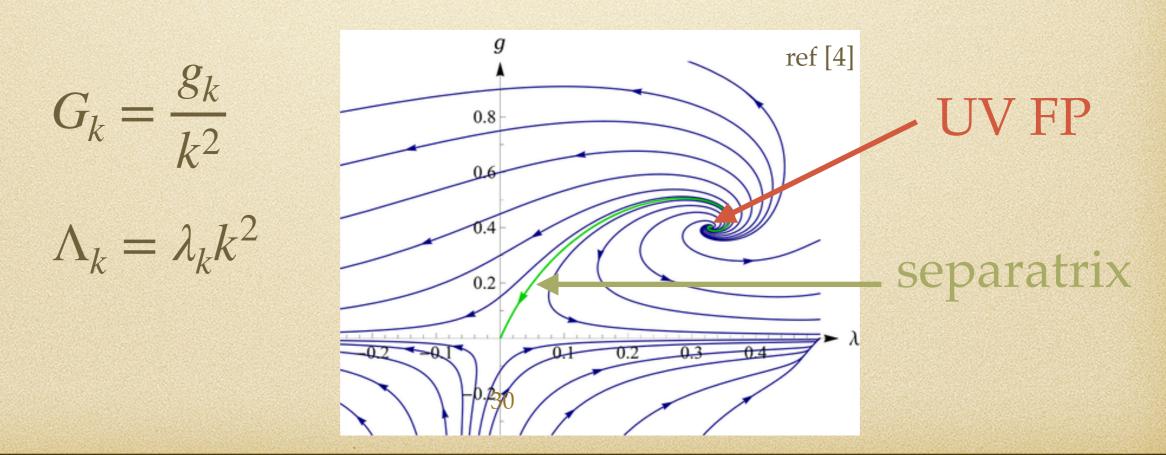


Asymptotic Safety
in a nutshell
$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2 \frac{\Lambda_{k}}{G_{k}} \right) + \dots$$

$$G_k = \frac{g_k}{k^2}$$
$$\Lambda_k = \lambda_k k^2$$







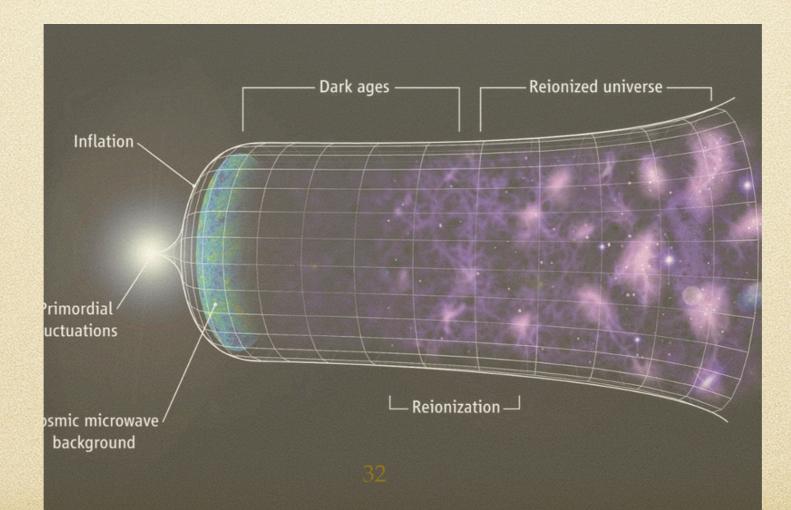
Comment by Giddings:

> Even far below M_{Pl} : Effect of non local operators

Scale Dependent
Framework
K, assume
$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2 \frac{\Lambda_{k}}{G_{k}} \right)$$

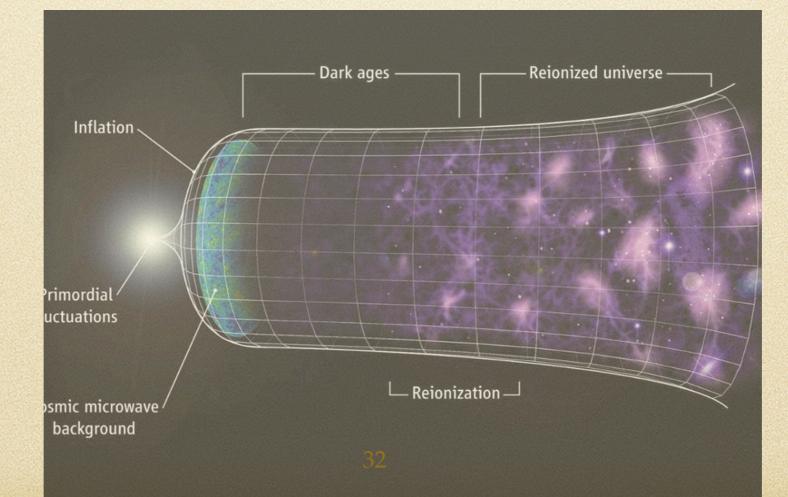
Scale DependentFrameworkFrameworkOK, assume $\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ Implications for CCP 3.0?

Scale DependentFrameworkOK, assume $\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ Implications for CCP 3.0?



Scale DependentFrameworkOK, assume $\Gamma_k = \int d^4 x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ Implications for CCP 3.0?

Maybe:



Scale Dependent
Framework
$$GK$$
, assume $\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ $Mr = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ $Maybe:$ $\frac{1}{G_k \cdot \Lambda_k} \approx 1$ $\int \frac{1}{G_k \cdot \Lambda_k} \approx 1$

Scale Dependent
Eracus
$$(M, n)$$
 $F_k = \int d^4 x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ (M, n) $F_k = \int d^4 x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right)$ $(Maybe)$ $1 + \frac{1}{G_k \cdot \Lambda_k} \approx 1 + \frac{1}{G_k \cdot \Lambda_k} + \frac{1}{G_k \cdot \Lambda_k} \neq 1$

Scale Dependent
Framework
$$(M, n) \in \mathbb{R}$$
 $(M, n) \in \mathbb{R}$ (M, n)

Scale Dependent
Framework
$$(M, n) \in \mathbb{R}$$
 $(M, n) \in \mathbb{R}$ $($

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right)$$

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Need to solve gap equations

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Need to solve gap equations

$$G_{\mu\nu} = -\Lambda_k g_{\mu\nu} - \Delta t_{\mu\nu}$$

with

$$\Delta t_{\mu\nu} = G_k \left(g_{\mu\nu} \nabla^{\alpha} \nabla_{\alpha} - \nabla_{\mu} \nabla_{\nu} \right) \frac{1}{G_k}$$

$$\Gamma_{k} = \int d^{4}x \sqrt{-g} \left(\frac{R}{G_{k}} - 2\frac{\Lambda_{k}}{G_{k}} \right)$$

Need to solve gap equations

with

$$G_{\mu\nu} = -\Lambda_k g_{\mu\nu} - \Delta t_{\mu\nu}$$
not constant!
$$\Delta t_{\mu\nu} = G_k \left(g_{\mu\nu} \nabla^{\alpha} \nabla_{\alpha} - \nabla_{\mu} \nabla_{\nu} \right) \frac{1}{G_k}$$

Assume homogenous background

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$$ds^{2} = -dt^{2} + a(t)\left(\frac{1}{1 - \kappa r^{2}}dr^{2} + r^{2}d\Omega_{2}^{2}\right)$$

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Gap equations

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Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda_k}{3} = \frac{1}{3}\rho_{SD}$$
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda_k = -p_{SD}$$

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$$ds^{2} = -dt^{2} + a(t)\left(\frac{1}{1 - \kappa r^{2}}dr^{2} + r^{2}d\Omega_{2}^{2}\right)$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \qquad \text{scale}$$

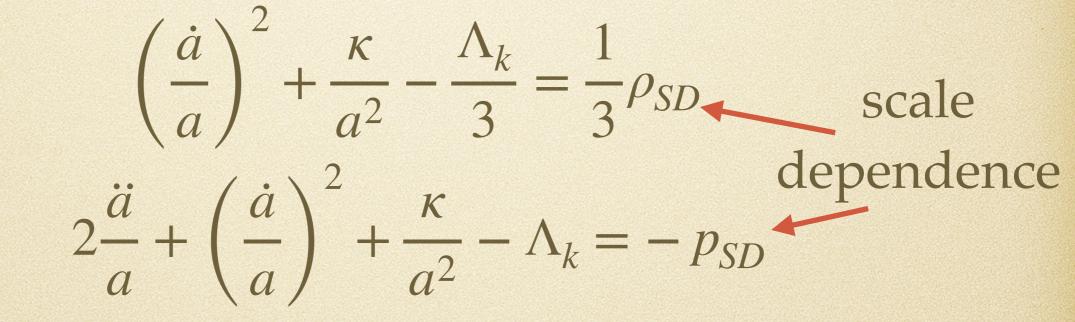
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \qquad \text{scale}$$

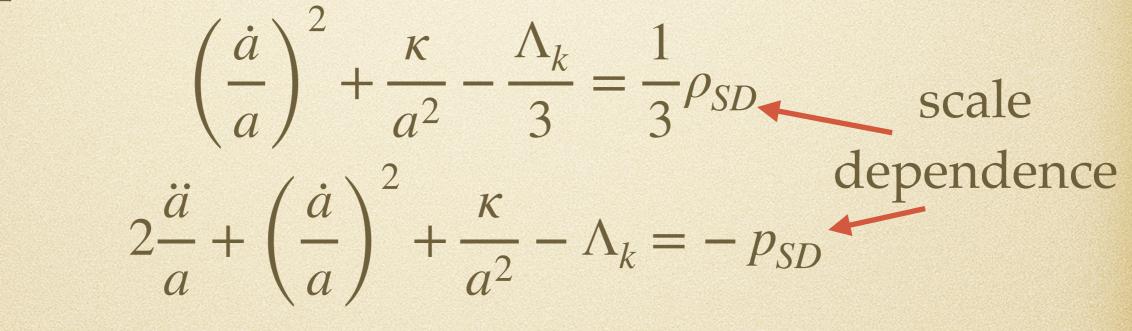
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$$
dependence

Gap equations



Since k=k(t)

Gap equations



Since $k=k(t) \Rightarrow$

Gap equations

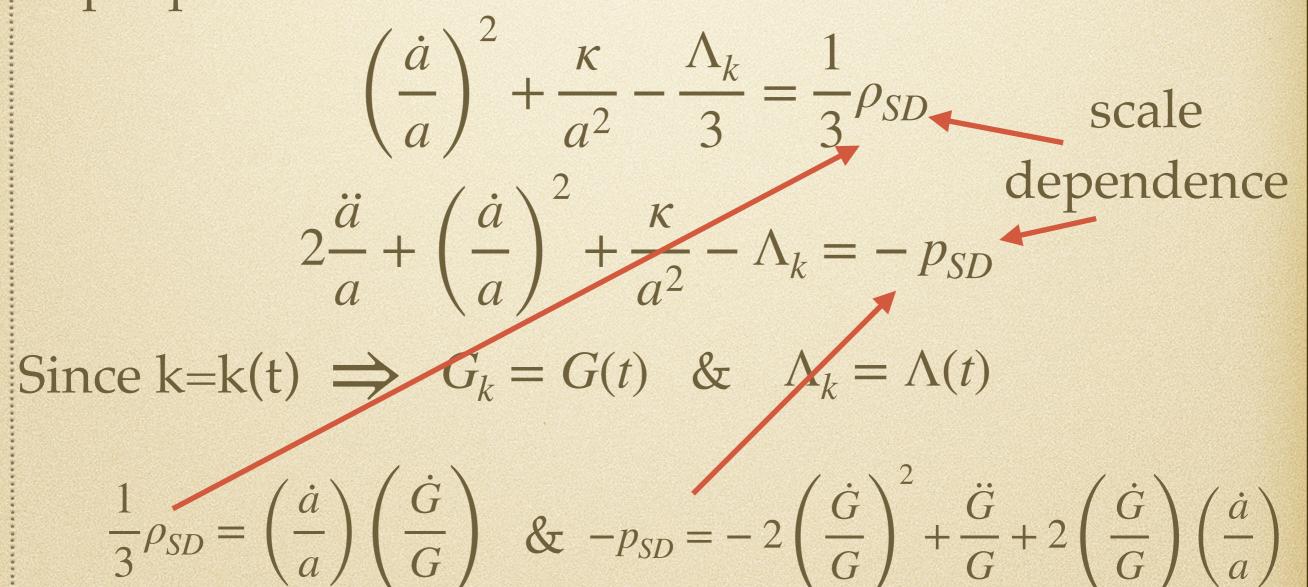
 $\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \qquad \text{scale}$ $2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$ Gince k=k(t) $\Rightarrow G_{k} = G(t) \& \Lambda_{k} = \Lambda(t)$

Gap equations

 $\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \quad \text{scale}$ $2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$ dependence $Since \ k=k(t) \implies G_{k} = G(t) \quad \& \quad \Lambda_{k} = \Lambda(t)$

$$\frac{1}{3}\rho_{SD} = \left(\frac{\dot{a}}{a}\right)\left(\frac{\dot{G}}{G}\right)$$

Gap equations

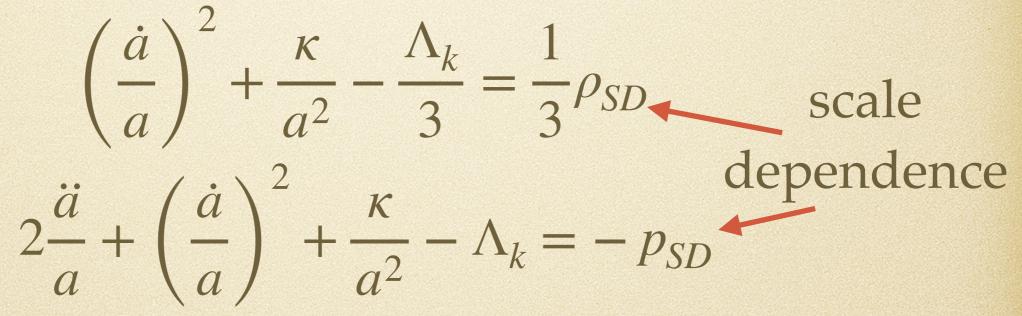


Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \qquad \text{scale}$$

$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$$
dependence

Gap equations



Note:

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda_{k}}{3} = \frac{1}{3}\rho_{SD} \quad \text{scale}$$

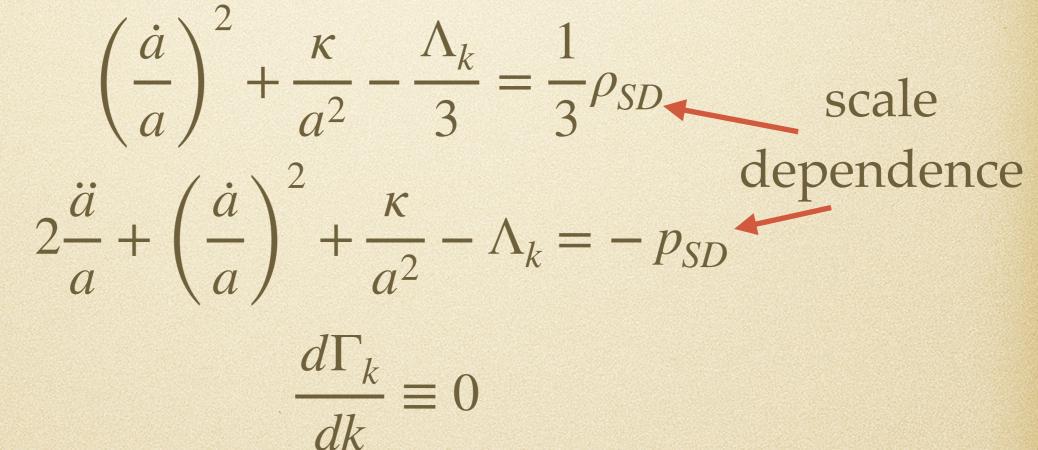
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda_{k} = -p_{SD}$$

$$\frac{d\Gamma_{k}}{dk} \equiv 0$$

Note:

Gap equations

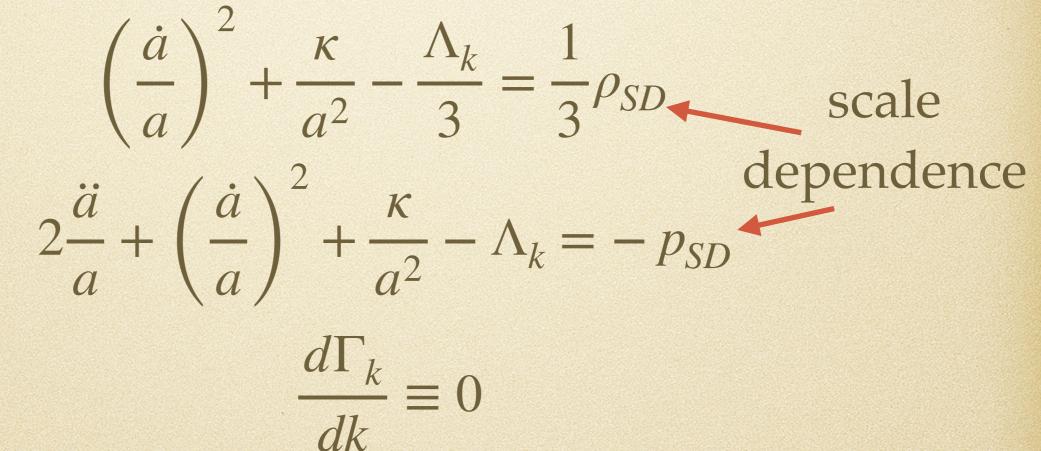
Note:



Ensures Bianchi identities & diff. inv.

Gap equations

Note:



Ensures Bianchi identities & diff. inv. $^{ref [8]}$ Back to G(t) ...

Gap equations

 $\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$ $2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$$
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Problem:

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right)\left(\frac{\dot{G}}{G}\right)$$
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Problem: 2 equations

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$$
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$

Problem: 2 equations

3 unknown functions: a(t), G(t), $\Lambda(t)$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$$
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$

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Solution:

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$$
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Problem: 2 equations

3 unknown functions: a(t), G(t), $\Lambda(t)$

Solution: Impose energy condition!

Null Energy Condition (NEC):

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 $\Delta t_{\mu\nu}\ell^{\mu}\ell^{\nu}=0$

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Null Energy Condition (NEC):

$$\Delta t_{\mu\nu}\ell^{\mu}\ell^{\nu}=0$$

$$\frac{d\ell^{\mu}}{dt} + \Gamma^{\mu}_{\alpha\beta}\ell^{\alpha}\ell^{\beta} = 0 \quad \longrightarrow \quad \ell^{\mu} = c_0 \frac{1}{a} \left(1, \frac{1}{\sqrt{1 - \kappa r^2}} \frac{1}{a}, 0, 0 \right)$$

Null Energy Condition (NEC):

 $\Delta t_{\mu\nu}\ell^{\mu}\ell^{\nu}=0$

where

$$\frac{d\ell^{\mu}}{dt} + \Gamma^{\mu}_{\alpha\beta}\ell^{\alpha}\ell^{\beta} = 0 \quad \longrightarrow \quad \ell^{\mu} = c_0 \frac{1}{a} \left(1, \frac{1}{\sqrt{1 - \kappa r^2}}, \frac{1}{a}, 0, 0 \right)$$

thus

$$-2\left(\frac{\dot{G}}{G}\right)^{2} + \left(\frac{\ddot{G}}{G}\right) - \left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right) = 0$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right)$$
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^{2} + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$
$$+ \left(\frac{\ddot{G}}{G}\right) - \left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right) = 2\left(\frac{\dot{G}}{G}\right)^{2}$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right) \qquad 2 \text{ gap}$$

$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^{2} + \frac{\kappa}{a^{2}} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^{2} + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$

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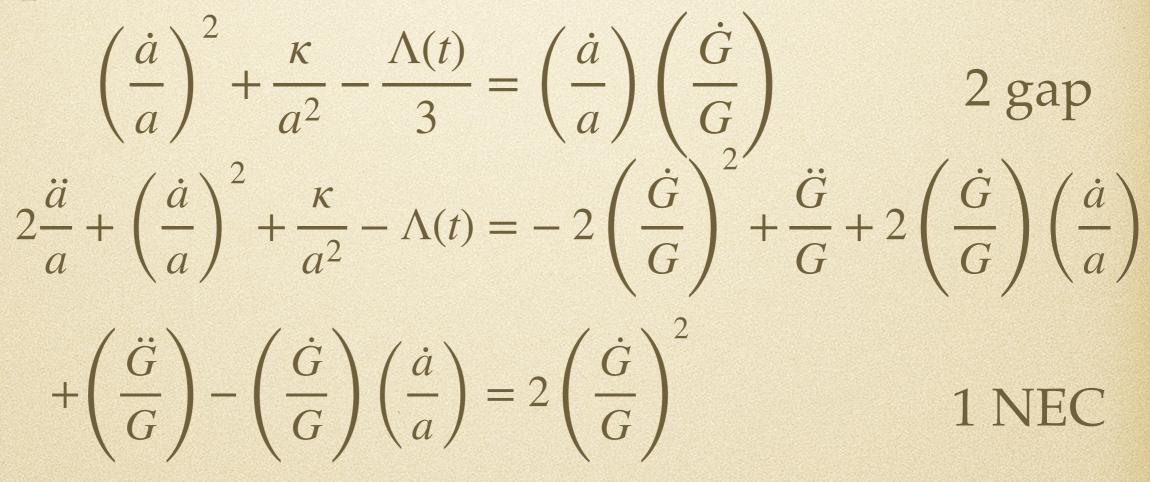
Gap equations

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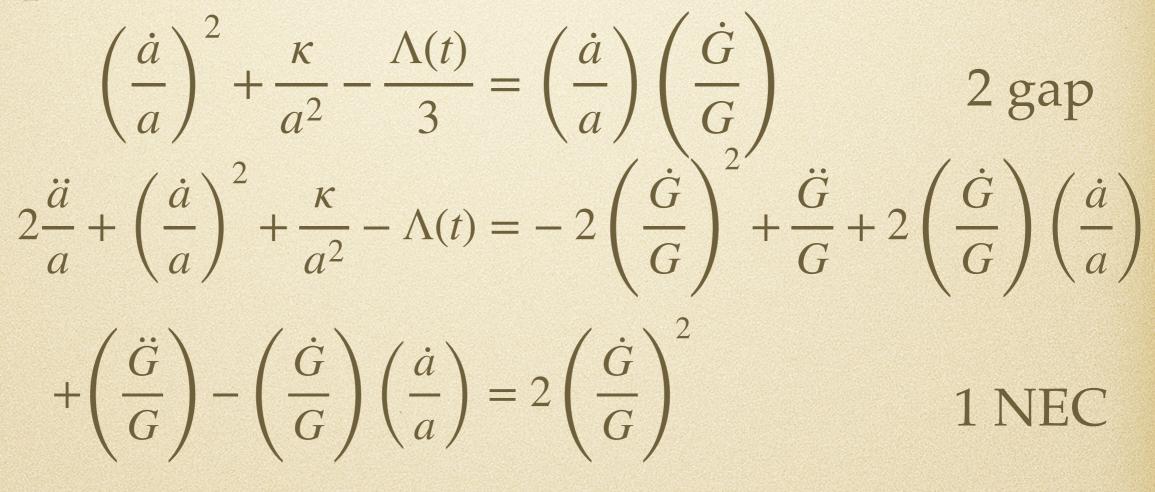
$$+ \left(\frac{\ddot{G}}{G}\right) - \left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right) = 2\left(\frac{\dot{G}}{G}\right)^{2} \qquad 1 \text{ NEC}$$

Gap equations

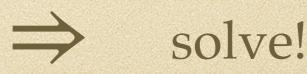


3 unknowns, 3 equations

Gap equations



3 unknowns, 3 equations



Solution:

$$a(t) = a_i e^{\frac{t}{\sqrt{\Lambda_0/3}}}$$
$$G(t) = \frac{G_0}{1 + \xi a(t)}$$

$$\Lambda(t) = \Lambda_0 \left[\frac{1 + 2\xi a(t)}{1 + \xi a(t)} \right]$$

Solution:

 $a(t) = a_i e^{\frac{t}{\sqrt{\Lambda_0/3}}}$ $G(t) = \frac{G_0}{1 + \xi a(t)}$

still inflation

$$\Lambda(t) = \Lambda_0 \left[\frac{1 + 2\xi a(t)}{1 + \xi a(t)} \right]$$

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3 integration constants:

 G_0, Λ_0, ξ

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3 integration constants:

41

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still inflation

$$\Lambda(t) = \Lambda_0 \left[\frac{1 + 2\xi a(t)}{1 + \xi a(t)} \right]$$

3 integration constants:

 G_0, Λ_0, ξ controls SD

Solution:

 $\lim_{\xi \to 0} G(t) = G_0$

 $\xi \rightarrow 0$

 $\lim a(t) = a_i e^{\sqrt{\Lambda_0/3}}$

 $\lim_{\xi \to 0} \Lambda(t) = \Lambda_0$

 G_0, Λ_0, ξ

controls SD

3 integration constants:

Solution:

$$a(t) = a_i e^{\frac{t}{\sqrt{\Lambda_0/3}}}$$
$$G(t) = \frac{G_0}{1 + \xi a(t)}$$

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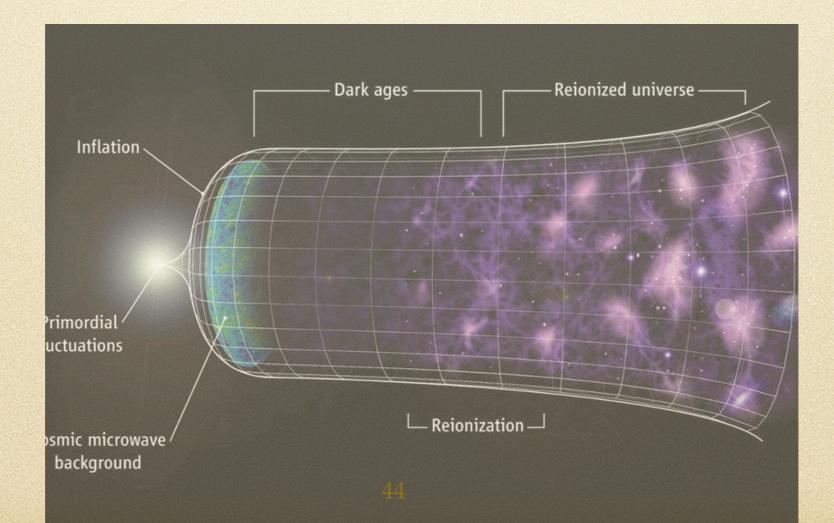
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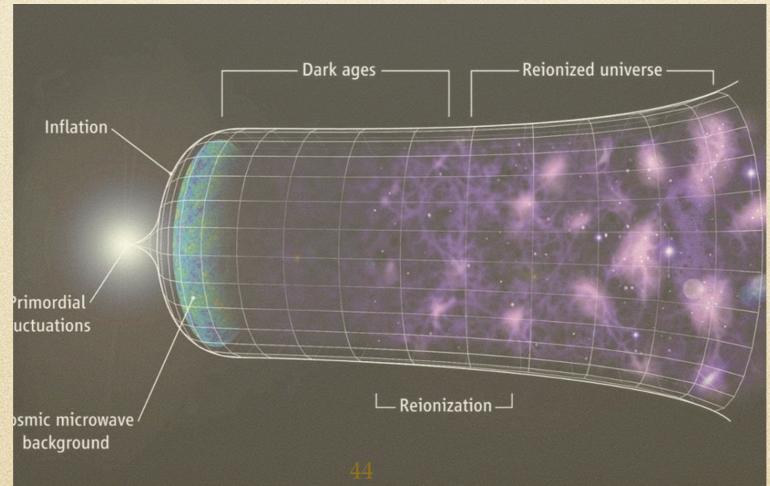
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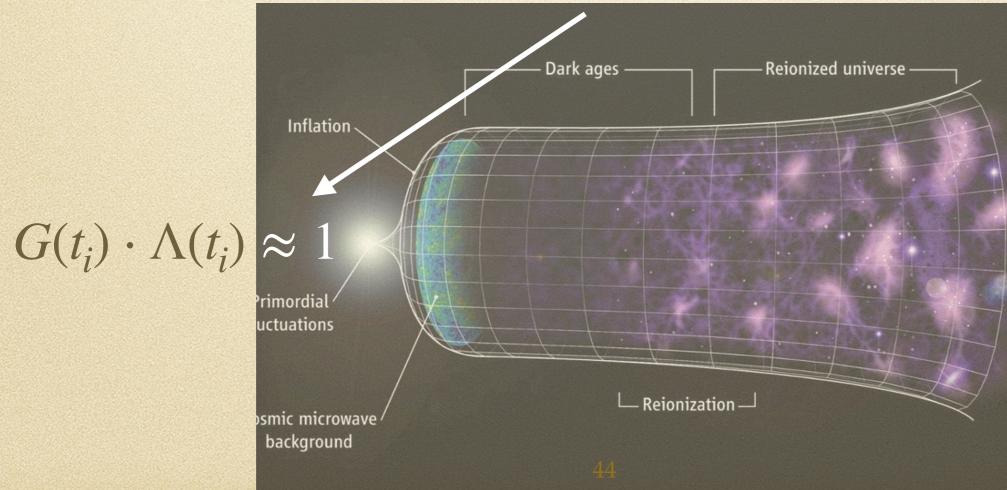
What does this mean for the CCP?

$G_k \cdot \Lambda_k = G(t) \cdot \Lambda(t)$



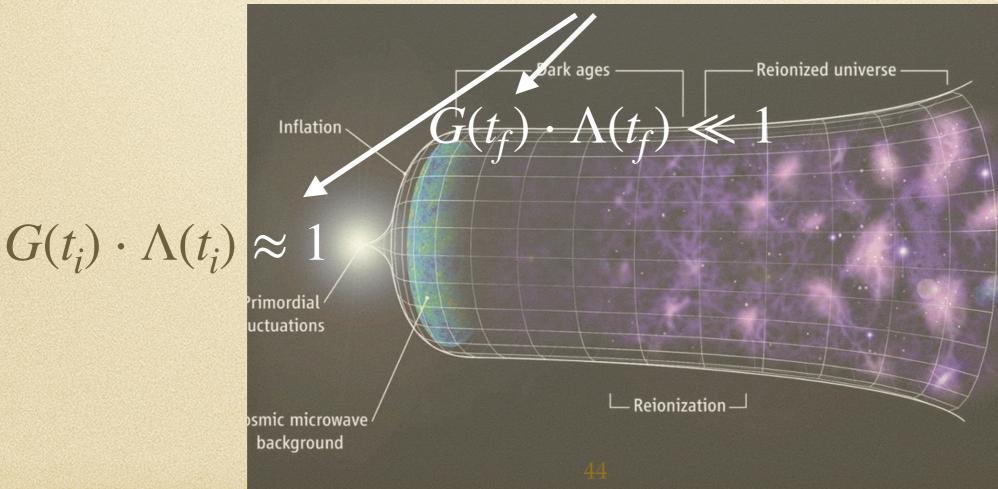
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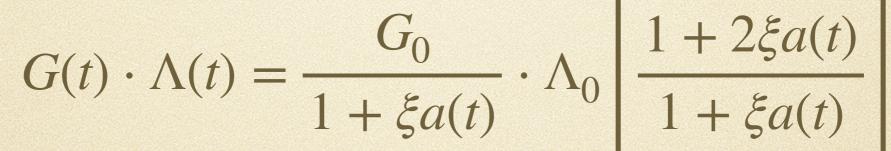


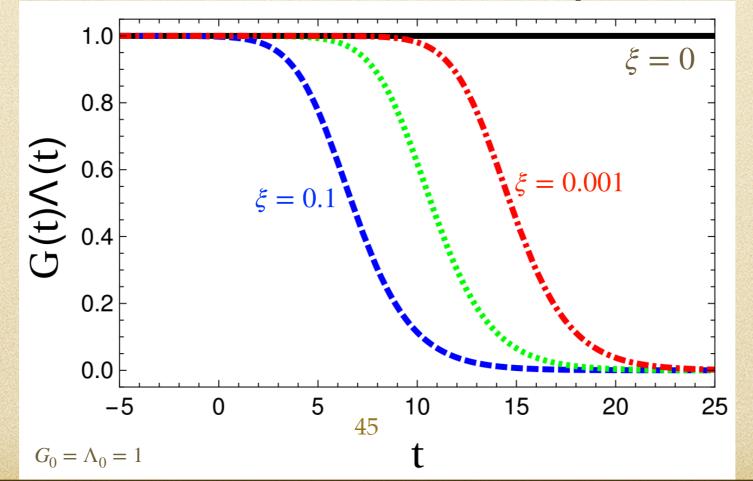
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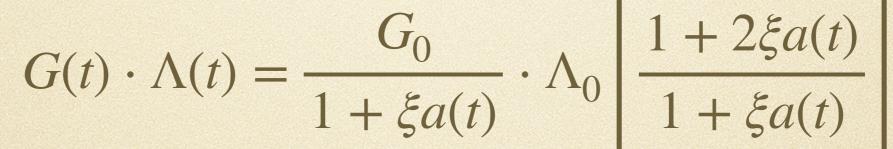


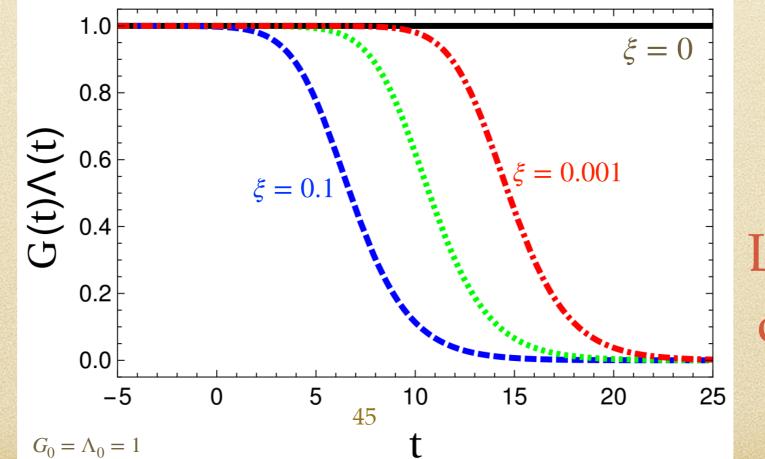
What does this mean for the CCP? $G(t) \cdot \Lambda(t) = \frac{G_0}{1 + \xi a(t)} \cdot \Lambda_0 \left[\frac{1 + 2\xi a(t)}{1 + \xi a(t)} \right]$





What does this mean for the CCP?____





Looks good, conditions?

CCP conditions on parameters

CCP conditions on parameters

- Initial a
- Initial CCP
- Final G
- Final CCP
- Flatness

CCP conditions on parameters

• Initial a

$$a(t_i) = 1$$

- Initial CCP
- Final G
- Final CCP
- Flatness

CCP conditions on parameters

• Initial a

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Initial CCP

 $\Lambda(t_i) \cdot G(t_i) = 1$

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CCP conditions on parameters

• Initial a

 $a(t_i) = 1$

Initial CCP

- $\Lambda(t_i) \cdot G(t_i) = 1$
- Final G

 $G(t_f) = G_N$

- Final CCP
- Flatness

CCP conditions on parameters

• Initial a

 $a(t_i) = 1$

Initial CCP

 $\Lambda(t_i) \cdot G(t_i) = 1$

- Final G
- Final CCP

 $G(t_f) = G_N$

 $G(t_f) \cdot \Lambda(t_f) = 10^{-(120\pm 5)}$

Flatness

CCP conditions on parameters

• Initial a

 $a(t_i) = 1$

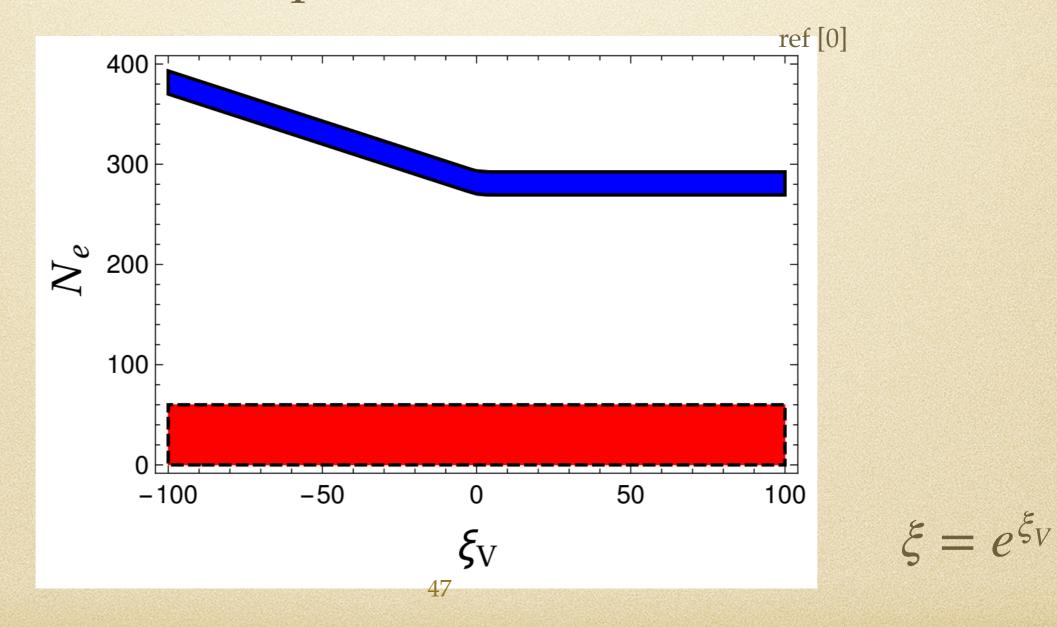
 $\Lambda(t_i) \cdot G(t_i) = 1$

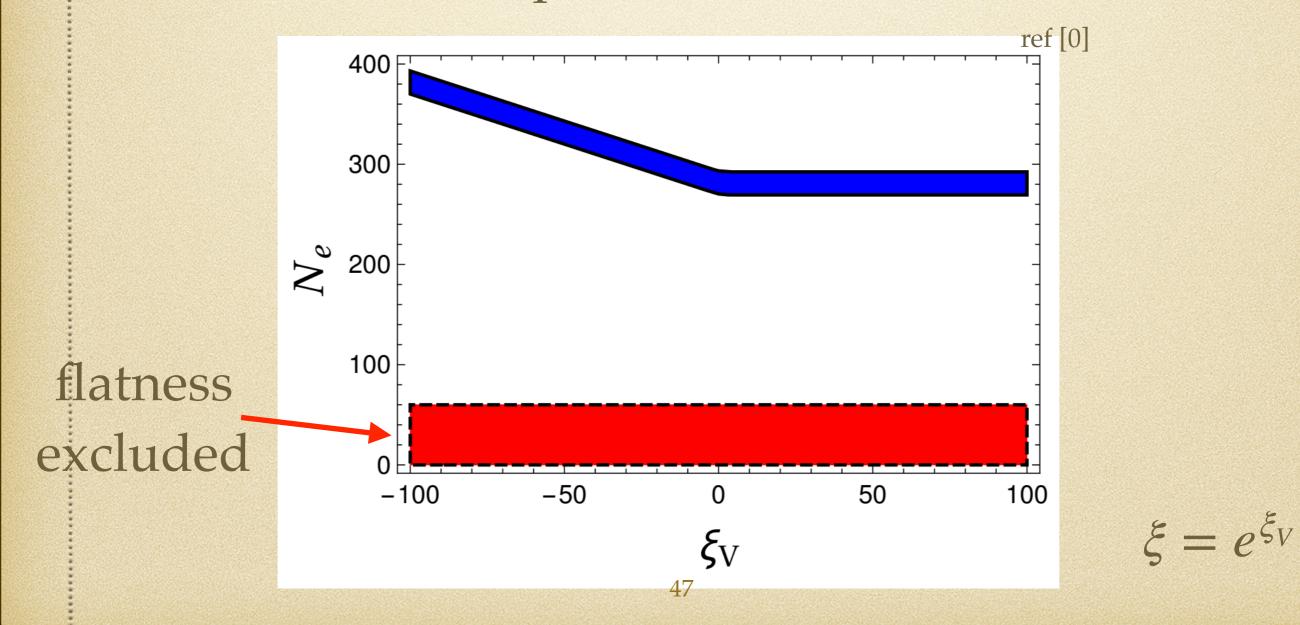
Initial CCP

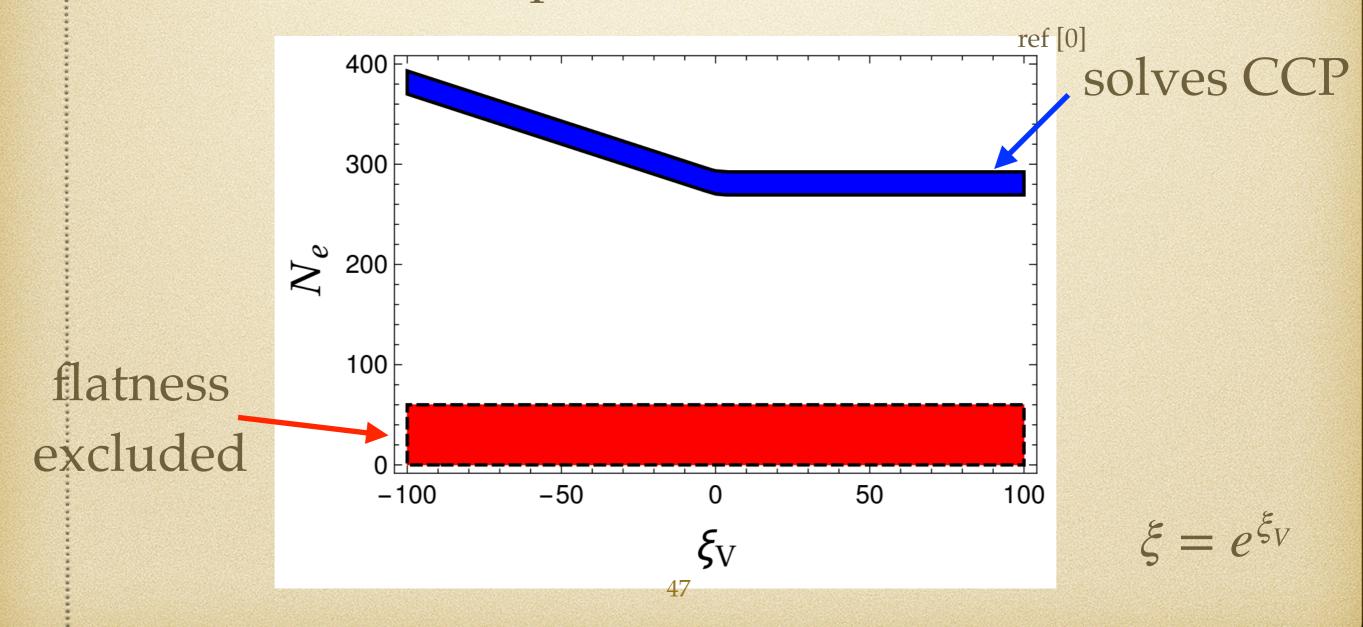
• Final G

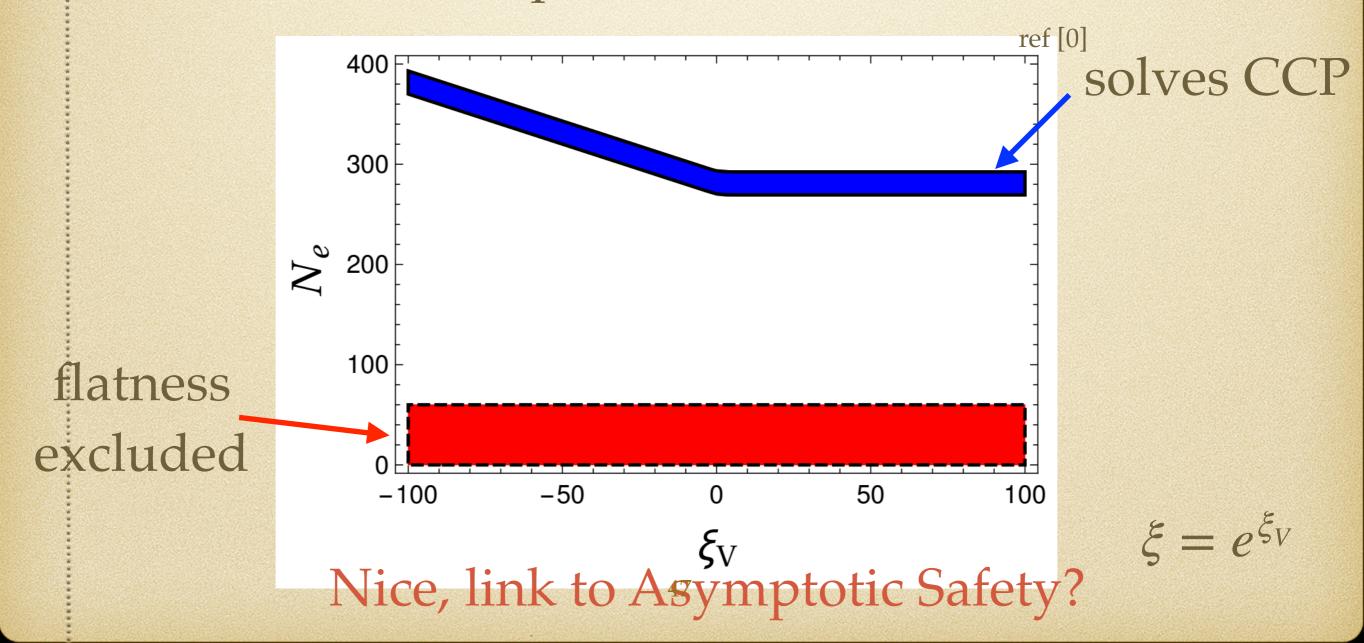
- Final CCP
- Flatness

 $G(t_f) = G_N$ $G(t_f) \cdot \Lambda(t_f) = 10^{-(120\pm 5)}$ $N_e \ge 60; \quad t_f - t_i = N_e \sqrt{\Lambda_0/3}$









Similar thoughts In the late Universe

Similar thoughts In the late Universe

Graviton fluctuations erase the cosmological constant

C. Wetterich^{*}

Universität Heidelberg, Institut für Theoretische Physik, Philosophenweg 16, D-69120 Heidelberg

arXiv:1704.08040v2

Similar thoughts In the late Universe

Graviton fluctuations erase the cosmological constant

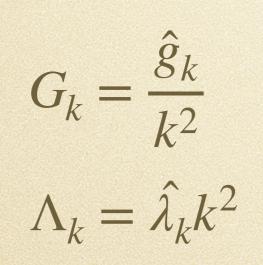
C. Wetterich^{*}

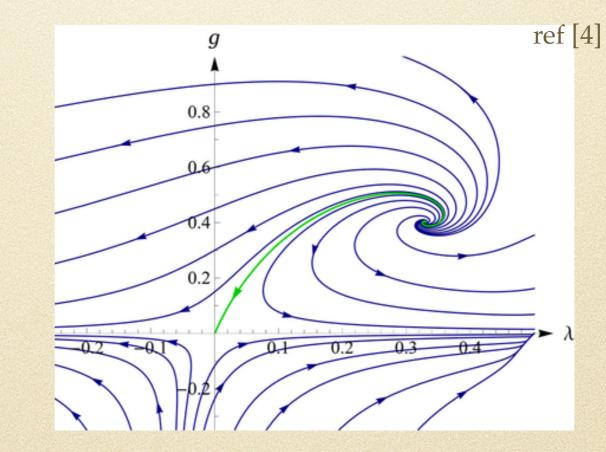
Universität Heidelberg, Institut für Theoretische Physik, Philosophenweg 16, D-69120 Heidelberg

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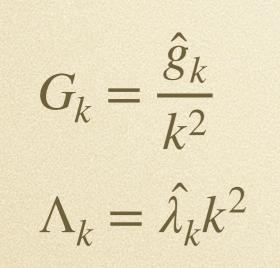
Now our idea coming from the UV:

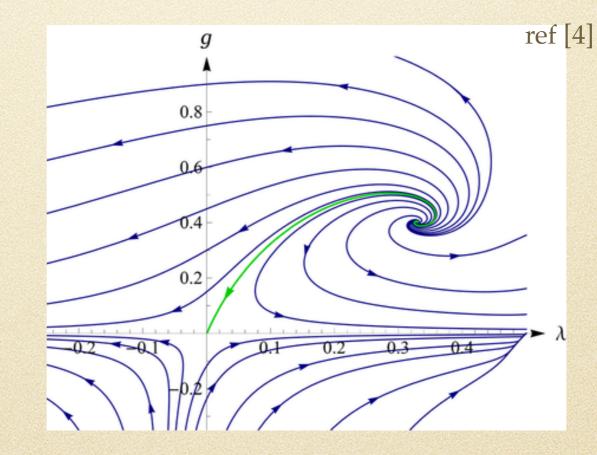
$$G_k = \frac{\hat{g}_k}{k^2}$$
$$\Lambda_k = \hat{\lambda}_k k^2$$



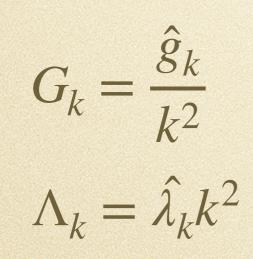


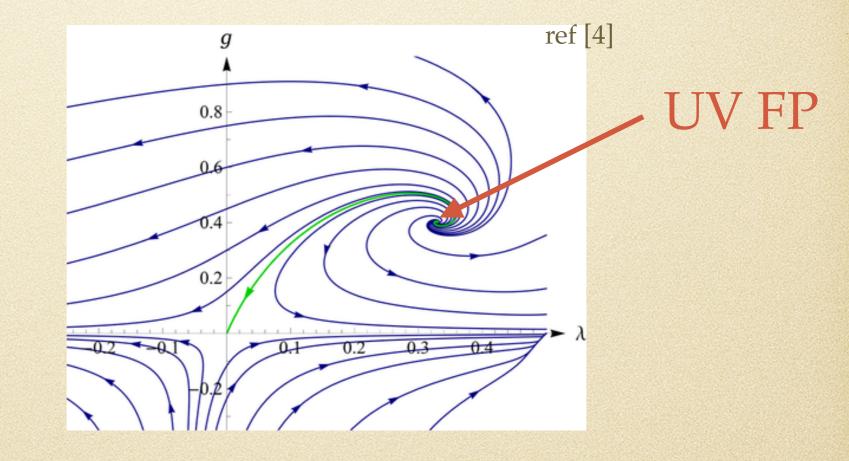
Remember:

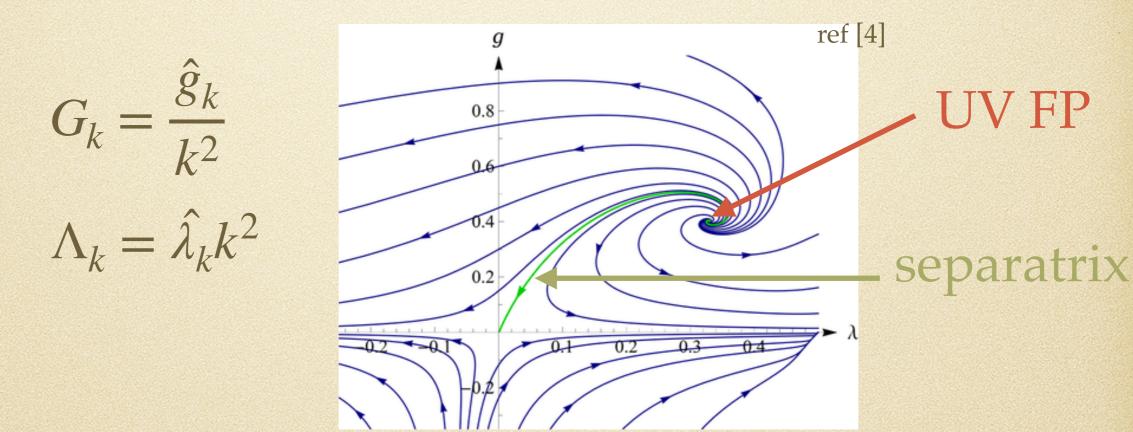


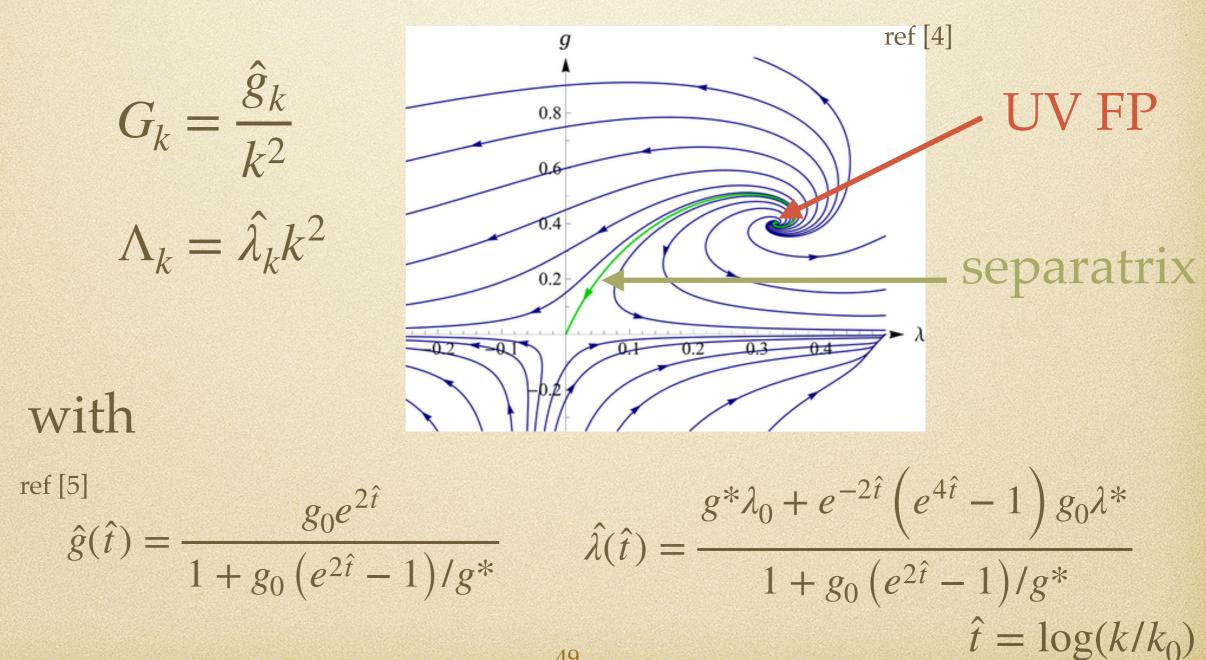


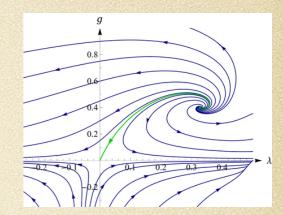
UV FP



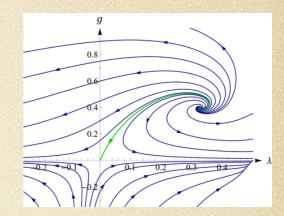






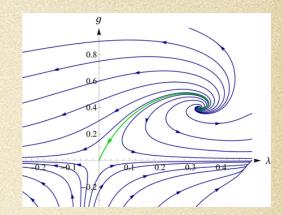


For CCP need:



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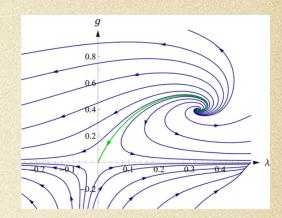
$$G_k \cdot \Lambda_k = \frac{\hat{g}_k}{k^2} k^2 \hat{\lambda}_k = \hat{g}_k \cdot \hat{\lambda}_k$$



For CCP need:

$$G_k \cdot \Lambda_k = \frac{\hat{g}_k}{k^2} k^2 \hat{\lambda}_k = \hat{g}_k \cdot \hat{\lambda}_k$$

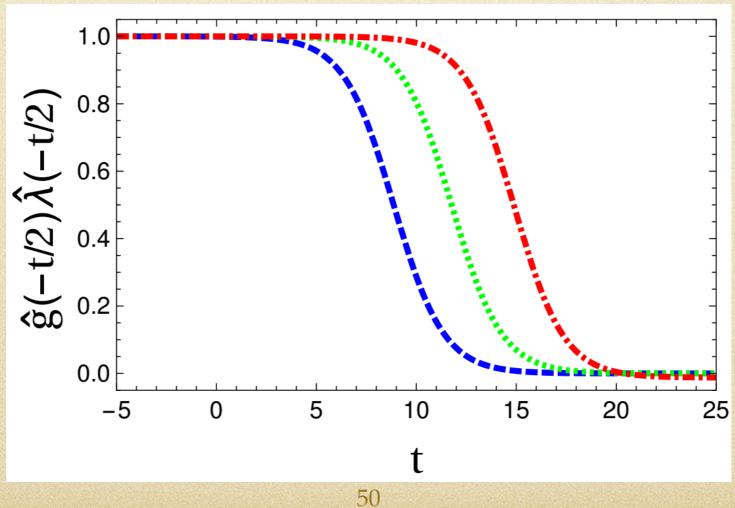
insert & plot

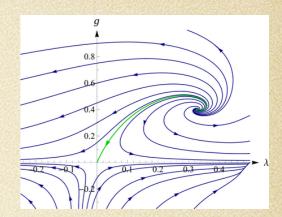


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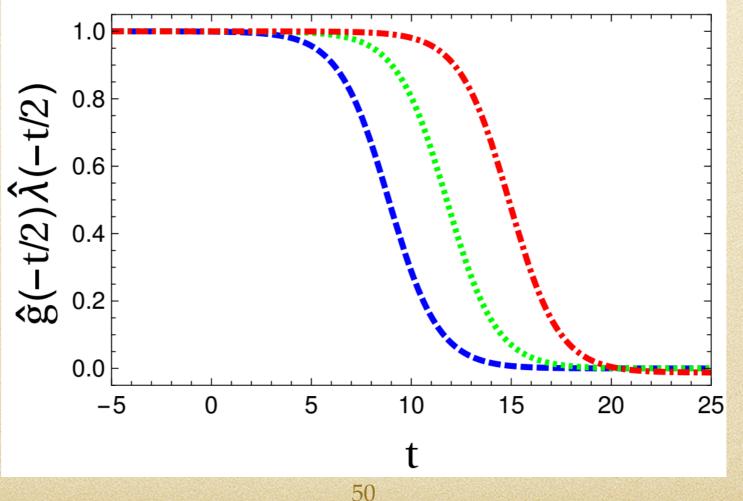




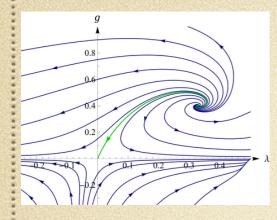
For CCP need:

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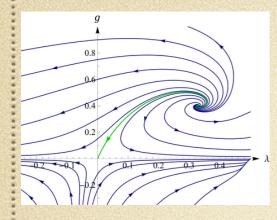
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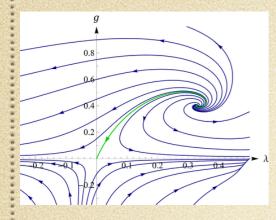
looks familiar?



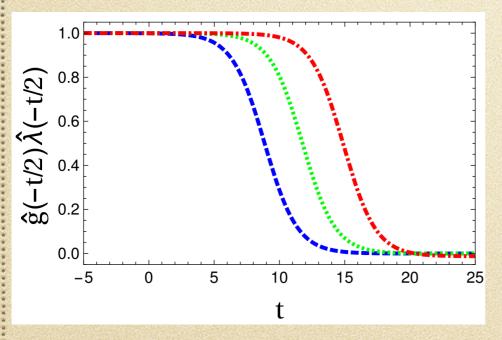
AS renormalization flow $G_k \cdot \Lambda_k = \frac{\hat{g}_k}{k^2} k^2 \hat{\lambda}_k = \hat{g}_k \cdot \hat{\lambda}_k$

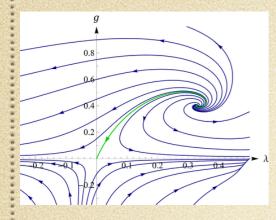


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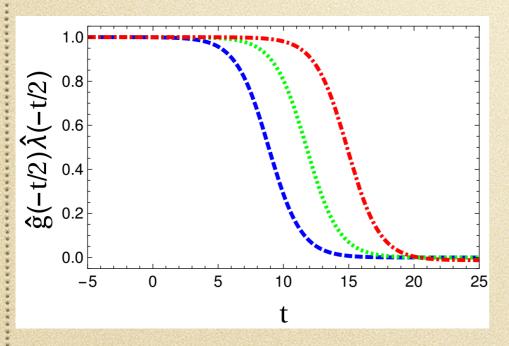


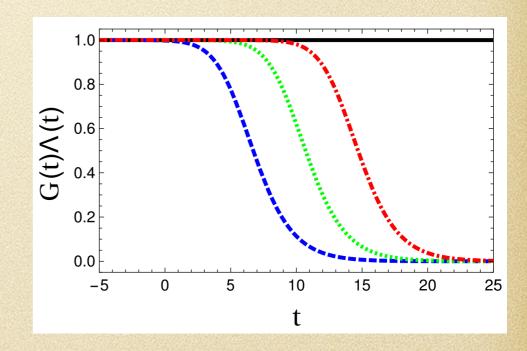
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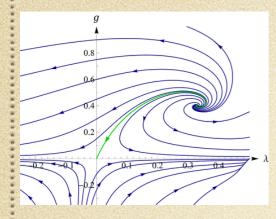




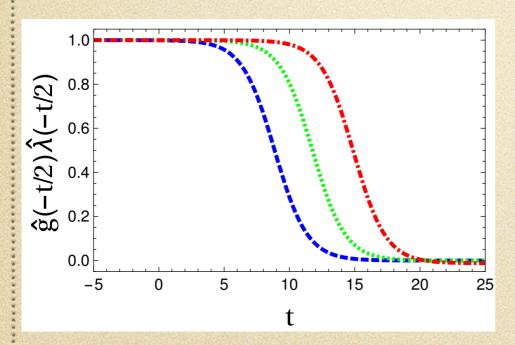
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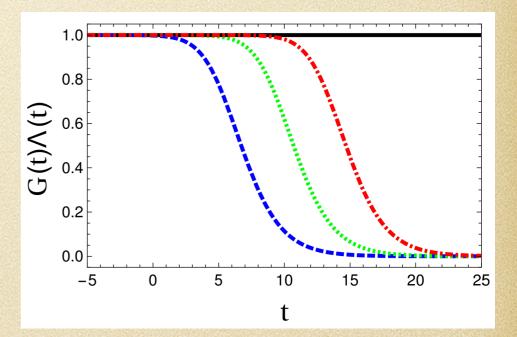


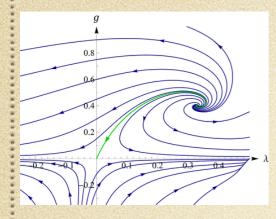


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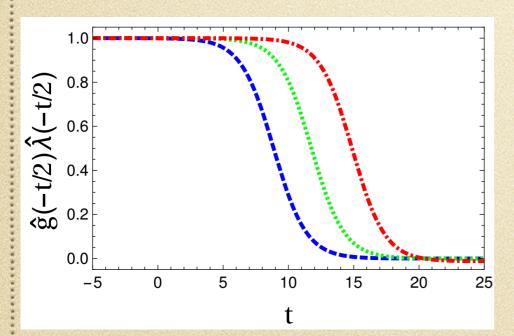


looks familiar!

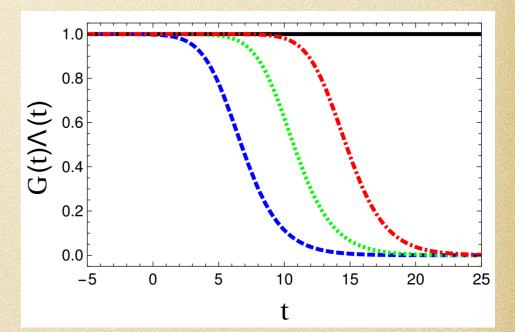




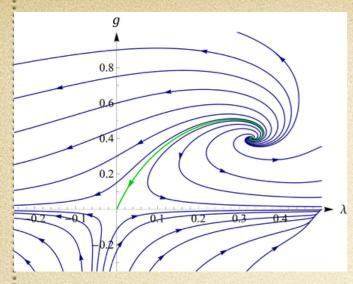
AS renormalization flow $G_k \cdot \Lambda_k = \frac{\hat{g}_k}{k^2} k^2 \hat{\lambda}_k = \hat{g}_k \cdot \hat{\lambda}_k$



looks familiar! SD & NEC $G(t) \cdot \Lambda(t)$

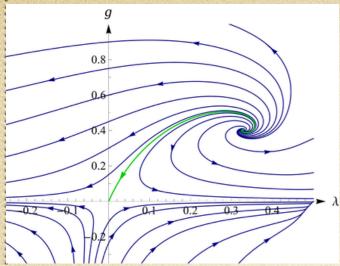


why, how?

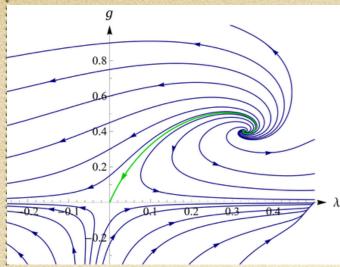


Remember:

52

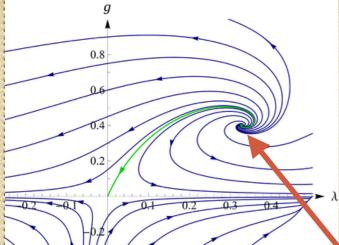


$$\hat{g}(\hat{t}) \cdot \hat{\lambda}(\hat{t}) = \frac{g_0 e^{2\hat{t}}}{1 + g_0 \left(e^{2\hat{t}} - 1\right)/g^*} \cdot \frac{g^* \lambda_0 + e^{-2\hat{t}} \left(e^{4\hat{t}} - 1\right) g_0 \lambda^*}{1 + g_0 \left(e^{2\hat{t}} - 1\right)/g^*}$$



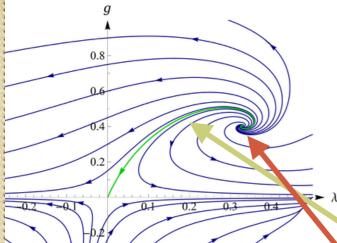
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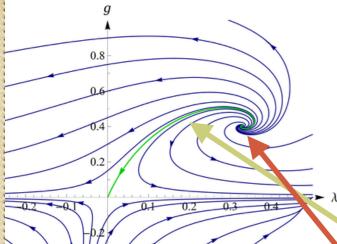
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Remember:

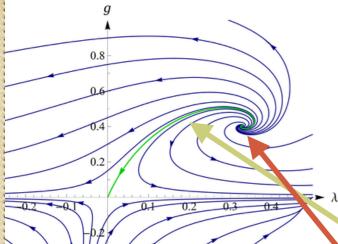
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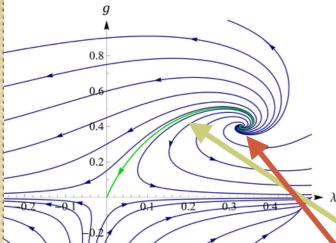
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Approximate to UV FP & separatrix

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For

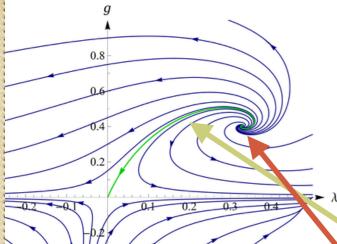


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$$g^*\lambda^* \to G_0\Lambda_0$$
$$g_0 \to G_0/(a_i\xi)$$
$$\hat{t} \to -t/(2\tau)$$



Remember:

For

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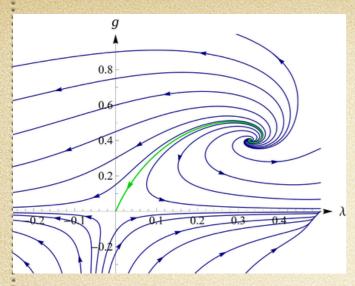
Approximate to UV FP & separatrix

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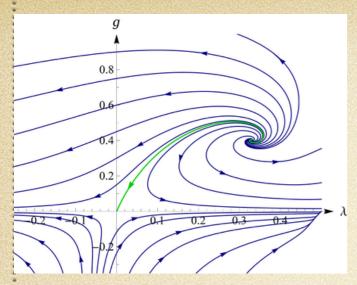
$$g^*\lambda^* \to G_0\Lambda_0$$

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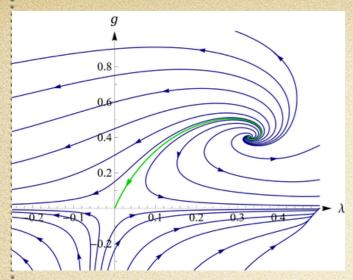
$$\hat{t} \to -t/(2\tau)$$



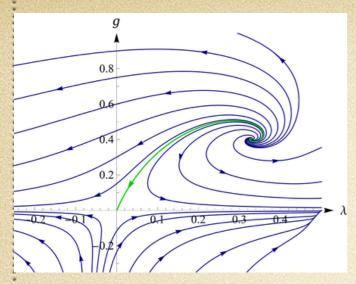
Comments on matching:



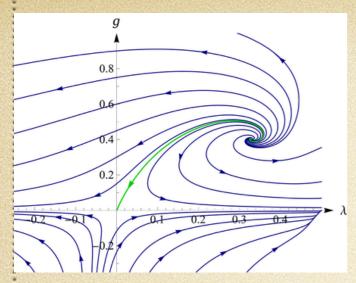
Comments on matching: $\hat{g}(\hat{t})\hat{\lambda}(\hat{t}) \equiv G(t) \cdot \Lambda(t)$



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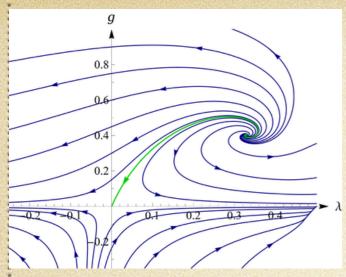
Comments on matching: $\hat{g}(\hat{t})\hat{\lambda}(\hat{t}) \equiv G(t) \cdot \Lambda(t)$ AS RG $g^*\lambda^* \to G_0\Lambda_0$ $g_0 \to G_0/(a_i\xi)$ $\hat{t} \to -t/(2\tau)$



Comments on matching: $\hat{g}(\hat{t})\hat{\lambda}(\hat{t}) \equiv G(t) \cdot \Lambda(t)$ $g^{*\lambda^*} \rightarrow G_0\Lambda_0$ $g_0 \rightarrow G_0/(a_i\xi)$ S

SD & NEC

 $\hat{t} \rightarrow -t/(2\tau)$



Comments on matching: $\hat{g}(\hat{t})\hat{\lambda}(\hat{t}) \equiv G(t) \cdot \Lambda(t)$

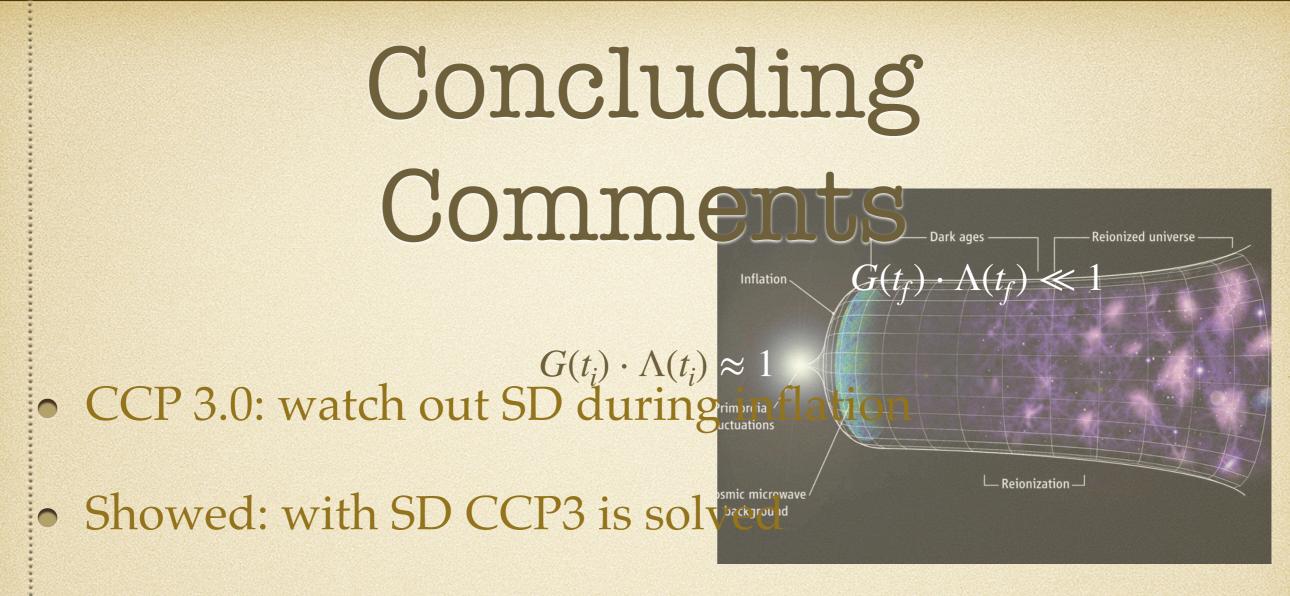
AS RG

 $g^* \lambda^* \to G_0 \Lambda_0$ $g_0 \to G_0 / (a_i \xi)$ $\hat{t} \to -t / (2\tau)$

SD & NEC

- Non trivial "coincidence"
- Works for many flow truncations
- UV FP @ inflation makes sense
- Separatrix special flow trajectory
- scale setting makes sense $\frac{k}{k_0} = e^{-t/(2\tau)}$

- CCP 3.0: watch out SD during inflation
- Showed: with SD CCP3 is solved
- Beautiful matching between AS & SD
- Outlook: Post inflation?



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solve

0

 $\xi_{
m V}$

50

100

-50

100

-100

- CCP 3.0: watch out SD during inflation
- Showed: with SD CCP3 is solved [≥]²⁰⁰
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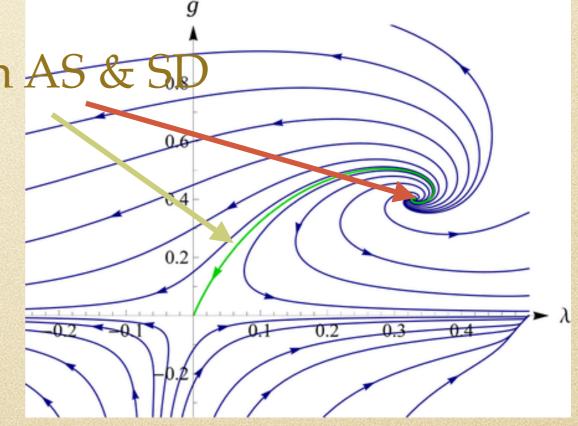
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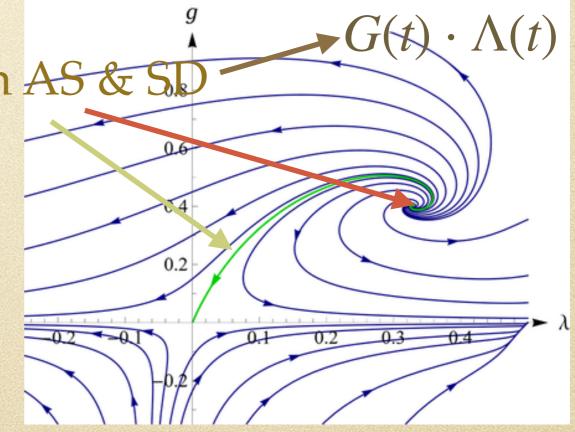


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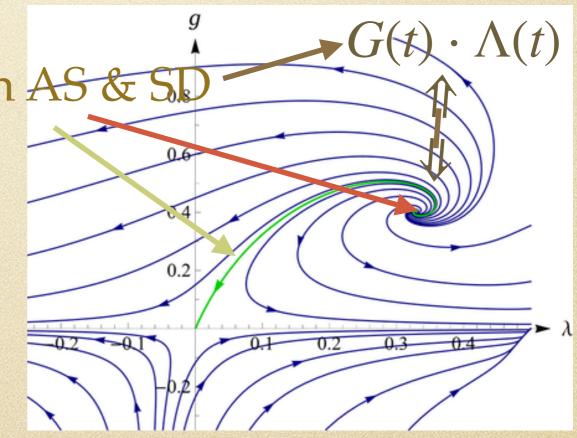
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CCP 3.0: watch out SD during inflation

Showed: with SD CCP3 is solved

Beautiful matching between AS & SD

Outlook: Post inflation?



Thank You!



Post inflation:

More couplings, more complicated

Post inflation:

More couplings, more complicated

Simplifying assumption:

Post inflation:

More couplings, more complicated Simplifying assumption:

Matter couplings don't run (so much)

Post inflation:

$$H_0^2 \Omega_{\Lambda}(t) + H_0^2 \Omega_{rad,0} a^{-4} g(t) + H_0^2 \Omega_{mat,0} a^{-3} g(t) = H^2(t) - H(t) \frac{g}{g},$$

$$-\frac{H_0^2 \Omega_{rad,0} g}{a^4} + 3H_0^2 \Omega_{\Lambda}(t) - H^2 + \frac{2H\dot{g}}{g} - \frac{2\dot{g}^2}{g^2} + \frac{\ddot{g}}{g} = 2\frac{\ddot{a}}{a}$$

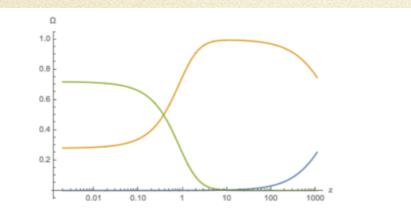
 $\dot{g}(g\dot{a} + 2a\dot{g}) = ag\ddot{g}$

Post inflation:

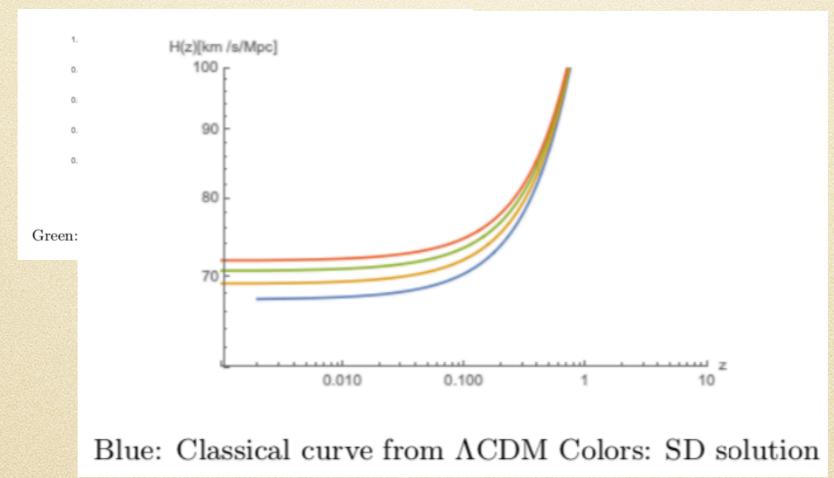
 $\begin{aligned} H_0^2 \Omega_{\Lambda}(t) + H_0^2 \Omega_{rad,0} a^{-4} g(t) + H_0^2 \Omega_{mat,0} a^{-3} g(t) &= H^2(t) - H(t) \frac{g}{g}, \\ - \frac{H_0^2 \Omega_{rad,0} g}{a^4} + 3H_0^2 \Omega_{\Lambda}(t) - H^2 + \frac{2H\dot{g}}{g} - \frac{2\dot{g}^2}{g^2} + \frac{\ddot{g}}{g} = 2\frac{\ddot{a}}{a} \\ \dot{g}(g\dot{a} + 2a\dot{g}) &= ag\ddot{g} \end{aligned}$

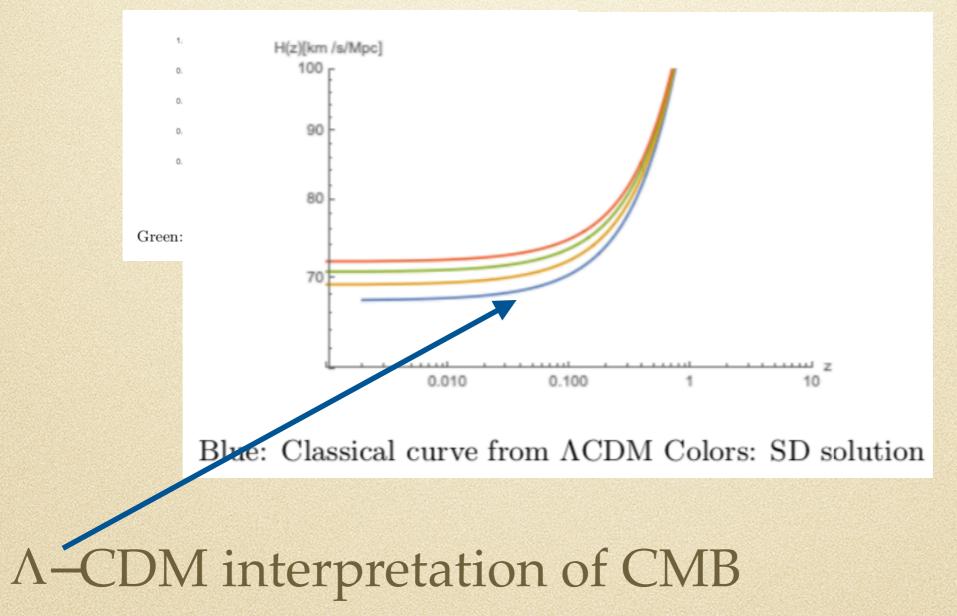
Solve Numerically and compare to NON-SD

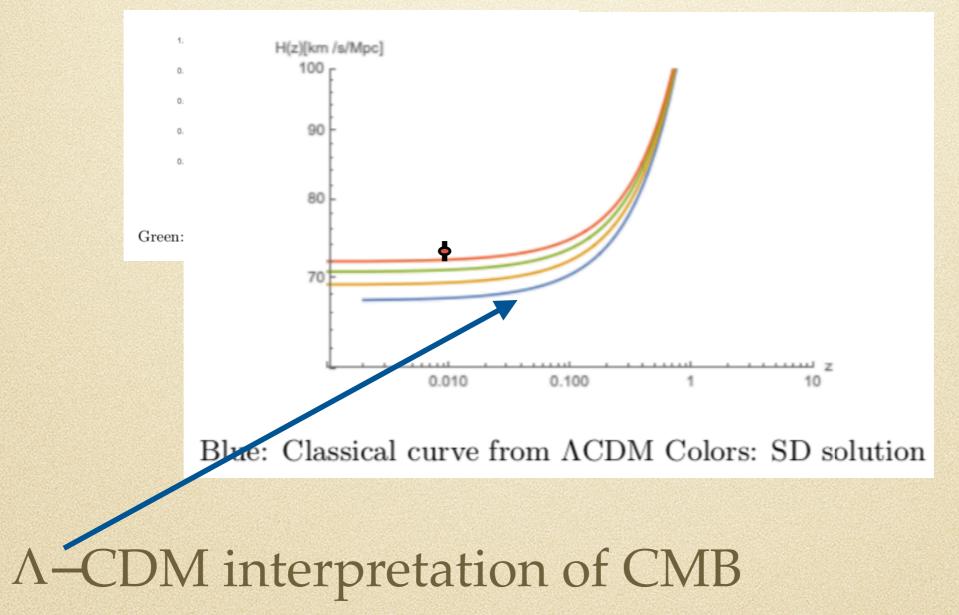
Post inflation:

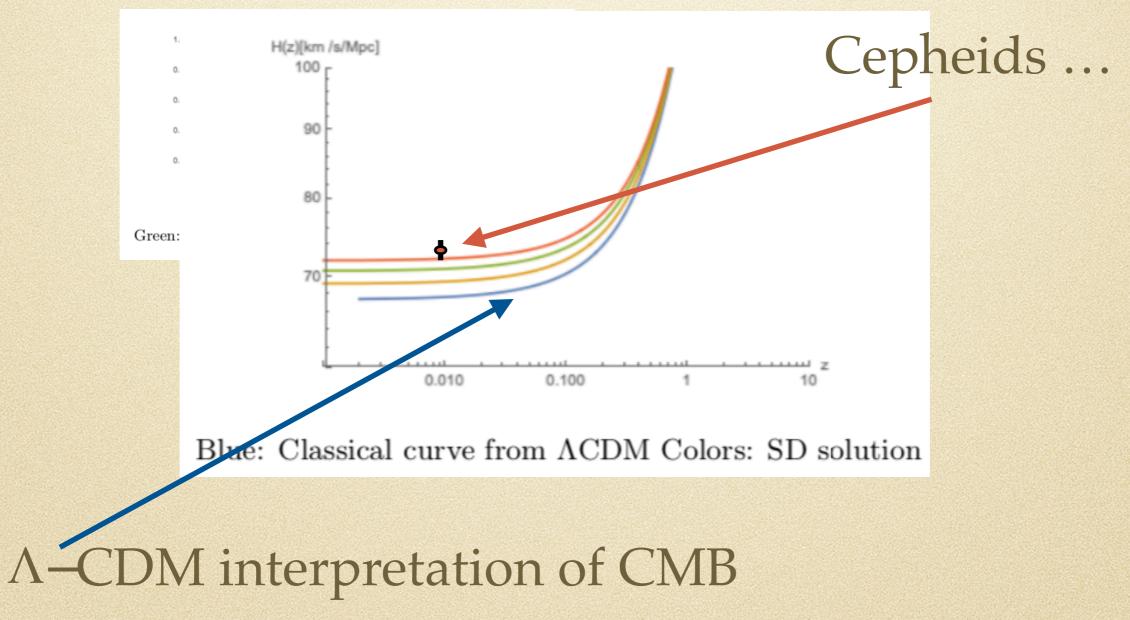


Green: Cosmological term, Orange: Matter, Blue: Radiation.









Literature

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