

Implementation of Hagedorn States into UrQMD

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HGS-HIRe *for FAIR*
Helmholtz Graduate School for Hadron and Ion Research

- 1 Motivation
 - Hagedorn States
 - UrQMD
- 2 Model
 - Pal-Danielewicz-approach
 - The Hagedorn spectrum
 - Decay width
- 3 Conclusion

Intention

- full integration of Hagedorn States (HS) into UrQMD
- study of the impacts on UrQMD due to HS on:
 - $\frac{\eta}{s}$
 - particle multiplicities
 - chemical equilibration times
 - ...

Hagedorn States

- resonance properties become increasingly uncertain with increasing resonance mass
- Rolf Hagedorn proposed in 1965 that particle number density must grow exponentially with the mass

$$\rho(m) \sim \exp\left(\frac{m}{T_H}\right) \quad (1)$$

- according to the Statistical-Bootstrap-Model near $T_H \approx 0.17$ GeV heavy resonances called Hagedorn States (*HS*) may become dramatically important
- HS were successfully applied in estimating chemical equilibration times for hyperons at RHIC
- HS are able to explain $\frac{\eta}{s}$ and the speed of sound near T_c

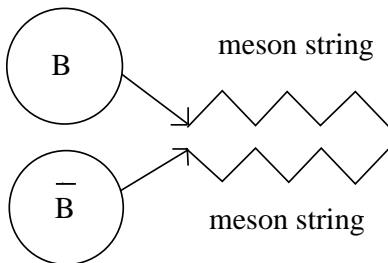
UrQMD

- microscopic Hadron-String Transport model simulating p+p, p+N and A+A collisions in the energy range from Bevalac and SIS up to AGS, SPS and RHIC
- detailed balance is enforced for following processes: meson-baryon, meson-meson, resonance-nucleon and resonance-resonance interaction
- but apart from particle production also string production is possible if total energy of collision participants is high enough
- string cross section is given by

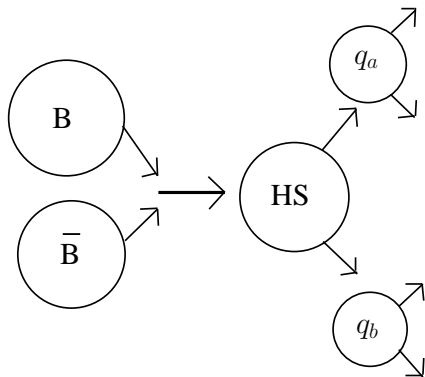
$$\sigma_{string} = \sigma_{tot} - \sigma_{el} - \sum_i \sigma_{inel}^i \quad (2)$$

First aspired changes

strings in UrQMD now



our project: strings \Rightarrow HS



Pal-Danielewicz-approach

- calculation of HS decay width via Pal-Danielewicz-approach
- statistical model for decay and formation of heavy hadronic resonances
- in analogy to low-energy nuclear physics Weisskopf's compound-nucleus model is employed

Modified Hagedorn spectrum

- modified Hagedorn spectrum with charges

$$\rho(m, \vec{q}) = A \frac{\exp\left[\frac{m - m_g(\vec{q}) f(m - m_g(\vec{q}))}{T_H}\right]}{\left[\{m - m_g(\vec{q}) f(m - m_g(\vec{q}))\}^2 + m_r^2\right]^\alpha} \quad (3)$$

$$(\vec{q} = (B, S, I))$$

- spectrum bottom m_g

$$m_g(\vec{q}) = \alpha_Q \max(|3B + S|, 2I) + \alpha_S |S| \quad (4)$$

- suppression function decreases the effect of m_g at high m

$$f(m - m_g(\vec{q})) = \frac{1}{1 + \left[\frac{m - m_g(\vec{q})}{m_c}\right]^n} \quad (5)$$

Fitting the Hagedorn spectrum

- all known discrete states of specie i with spin J_i and mass m_i

$$N_{ex}(m) = \sum_i (2J_i + 1) \Theta(m - m_i) \quad (6)$$

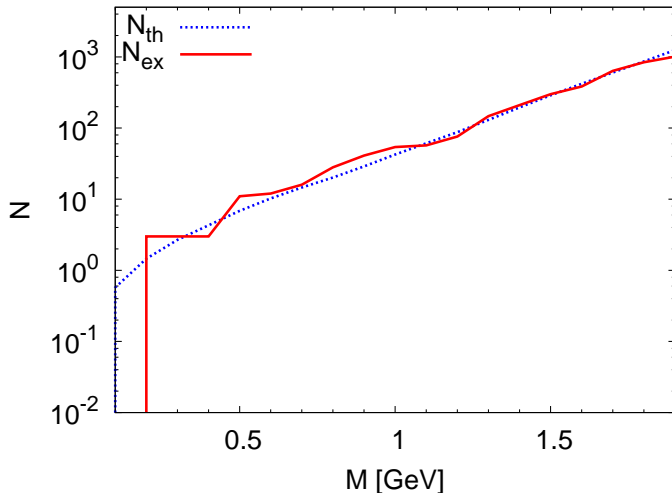
- Hagedorn hypothesis of a continuous particle density at different \vec{q}

$$N_{th}(m) = \sum_q \int_{m_g(\vec{q})}^m dm' \rho(m', \vec{q}) \quad (7)$$

- minimizing of $\chi^2(N_{ex}, N_{th})$ (*MINUIT*) gives A and α
- final parameter set used reads:

- | | | |
|--------------------------|--------------------|-------------------|
| • $\alpha_Q = 0.387$ GeV | • $m_c = 1.8$ GeV | • $\alpha = 1.95$ |
| • $\alpha_S = 0.459$ GeV | • $T_H = 0.17$ GeV | • $A = 0.29$ |

Theoretical and experimental cumulants



Cross sections

- formation cross section of a resonance q (h or HS) by following interactions: $q_a + q_b$, $q_a + q_b$, $q_a + q_b$ ($q_i \cong h$, $q_i \cong HS$, $q_i \cong h$ or HS)
- we enforce detailed balance between decay and formation of the resonance q so $|M_{q_a+q_b \rightarrow q}|^2 = |M_{q \rightarrow q_a+q_b}|^2$
- general cross section

$$\sigma(q_a + q_b \rightarrow q) = \frac{2\pi m_a m_b}{m p^* (m_a, m_b)} \rho(m, \vec{q}) |M_{q_a+q_b \rightarrow q}|^2 \quad (8)$$

- geometrical cross section

$$\sigma(q_a + q_b \rightarrow q) = \langle I^a I_3^a I^b I_3^b | I I_3 \rangle^2 \pi R^2 \quad (9)$$

- radius R of HS with $r_0=1$ fm and $m_d=1$ GeV is defined as:

$$R(m) = r_0 \left(\frac{m}{m_d} \right)^{\frac{1}{3}} \quad (10)$$

General expression and decay mode (i)

- general expression for partial decay width of resonance q into the daughters q_a and q_b

$$\Gamma(q \rightarrow q_a + q_b) = \int \frac{d\vec{p}}{(2\pi)^3} \int dm'_a \frac{\rho(m'_a, \vec{q}_a)}{e_a(\vec{p})} \int dm'_b \frac{\rho(m'_b, \vec{q}_b)}{e_b(\vec{p})} \times \quad (11)$$

$$|M_{q \rightarrow q_a + q_b}|^2 2\pi \delta(e_a(\vec{p}) + e_b(\vec{p}) - m)$$

- first binary decay mode (i) concerns a HS-decay into two hadrons with particle density $\rho_i(m'_i) = (2J_i + 1) \delta(m'_i - m_i)$:

$$\Gamma^{(i)}(q \rightarrow q_a + q_b) = \langle I^a I_3^a I^b I_3^b | I I_3 \rangle^2 \times \quad (12)$$

$$\frac{(2J_a + 1)(2J_b + 1) p^*(m_a, m_b) R^2}{2\pi \rho(m, \vec{q})}$$

Decay modes (ii) and (iii)

- second binary decay mode (ii) concerns a HS-decay into a hadron q_a and a HS q_b

$$\Gamma^{(ii)}(q \rightarrow q_a + q_b) = \langle I^a I_3^a I^b I_3^b | I I_3 \rangle^2 \times \quad (13)$$

$$\frac{(2J_a + 1) m R^2}{2\pi \rho(m, \vec{q})} \int_0^{p^*(m_a, m_b)} dp p^3 \frac{\rho\left(\sqrt{m^2 + m_b^2 - 2me_a}, \vec{q}_a\right)}{e_a \sqrt{m^2 + m_a^2 - 2me_a}}$$

- third binary decay mode (iii) concerns a HS-decay into two HSs.

$$\Gamma^{(iii)}(q \rightarrow q_a + q_b) = \langle I^a I_3^a I^b I_3^b | I I_3 \rangle^2 \times \quad (14)$$

$$\frac{R^2}{2\pi \rho(m, \vec{q})} \int_{m_a^c}^{m - m_b^c} dm'_a \int_{m_b^c}^{m - m_a} dm'_b p^*(m'_a, m'_b) \rho(m'_a, \vec{q}_a) \rho(m'_b, \vec{q}_b)$$

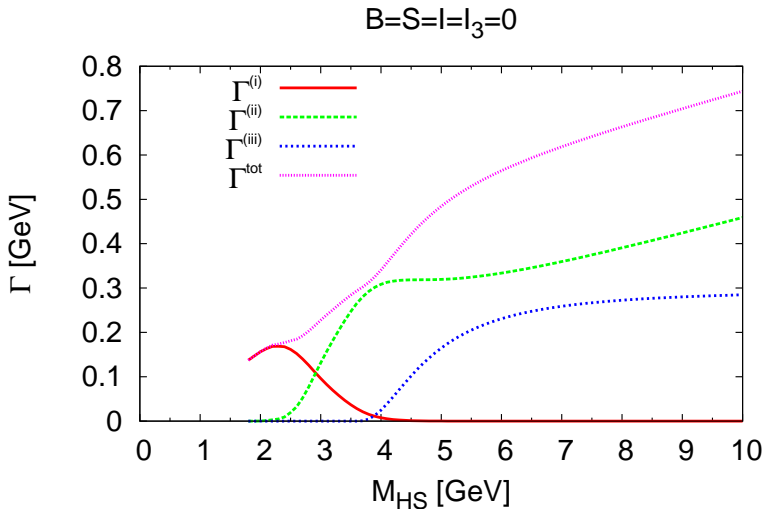
$$(m_i^c = \max[m_c, m_g(\vec{q}_i)])$$

Total decay width

- total decay width $\Gamma(m, \vec{q})$ is the sum over all decay channels where B,S and I are conserved exactly

$$\Gamma(m, \vec{q}) = \sum_{\vec{q}_a, \vec{q}_b} \Gamma^{(i)}(q \rightarrow q_a + q_b) + \sum_{\vec{q}_a, \vec{q}_b} \Gamma^{(ii)}(q \rightarrow q_a + q_b) + \sum_{\vec{q}_a, \vec{q}_b} \Gamma^{(iii)}(q \rightarrow q_a + q_b) \quad (15)$$

Decay widths for charge-neutral HS



Conclusion

- replacement of strings in UrQMD by HS
- HS creations in secondary collisions like $b\bar{b}$ as first step
- presentation of a decay width model of HS

THANK YOU FOR YOUR ATTENTION!

Appendix

