

# Matching Stages In Heavy Ion Collision Models

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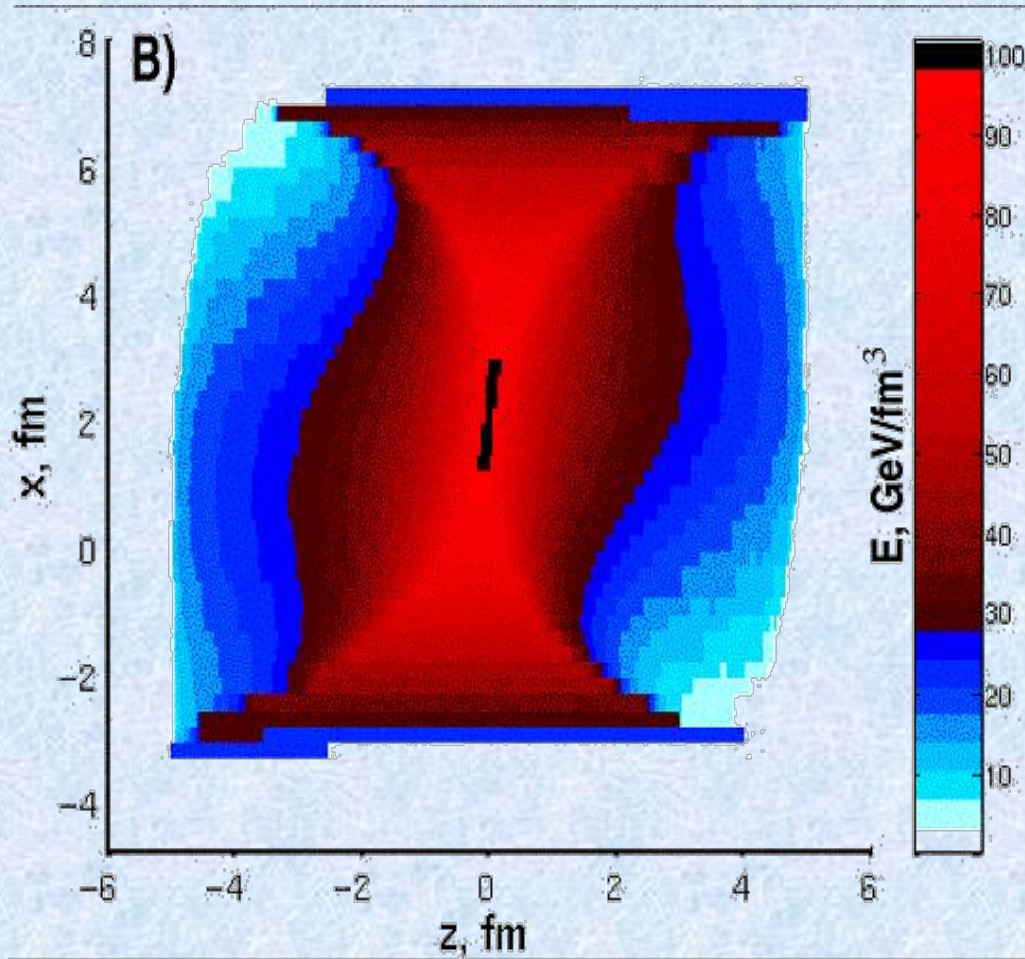
L.P. Csernai

# Outline

- Introduction → Multi-Module model
- Numerical Extraction of FO hyper-surface
- Cooper-Frye Description (space-like FO ?)
- Anisotropic distribution (cut-Juttner, cancelling Juttner )
- Simple covariant solution
- Summary

[Y. Cheng, et al.,  
Phy. Rev. C81, 064910 (2010)]

# Relativistic fluid dynamical model

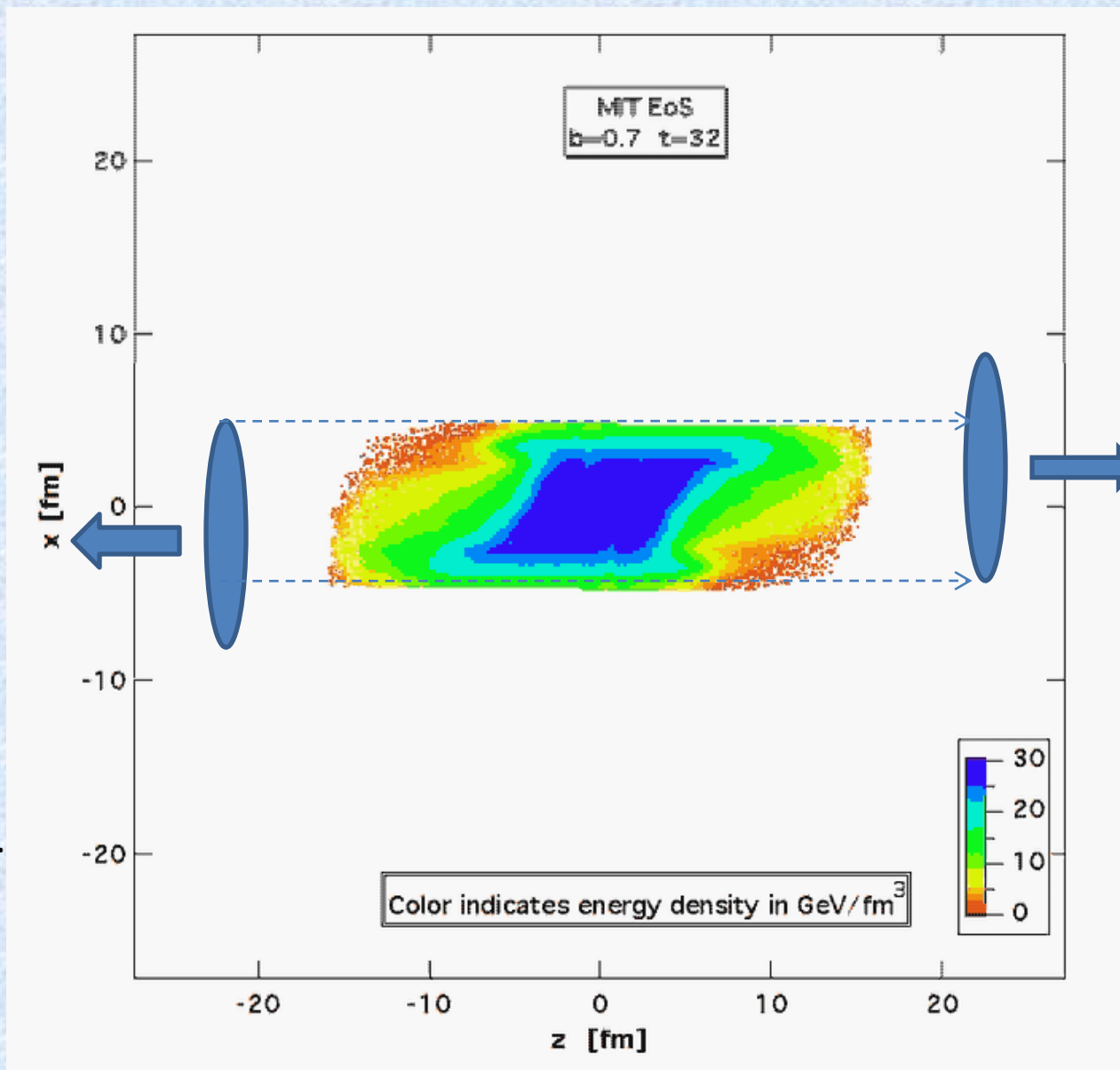


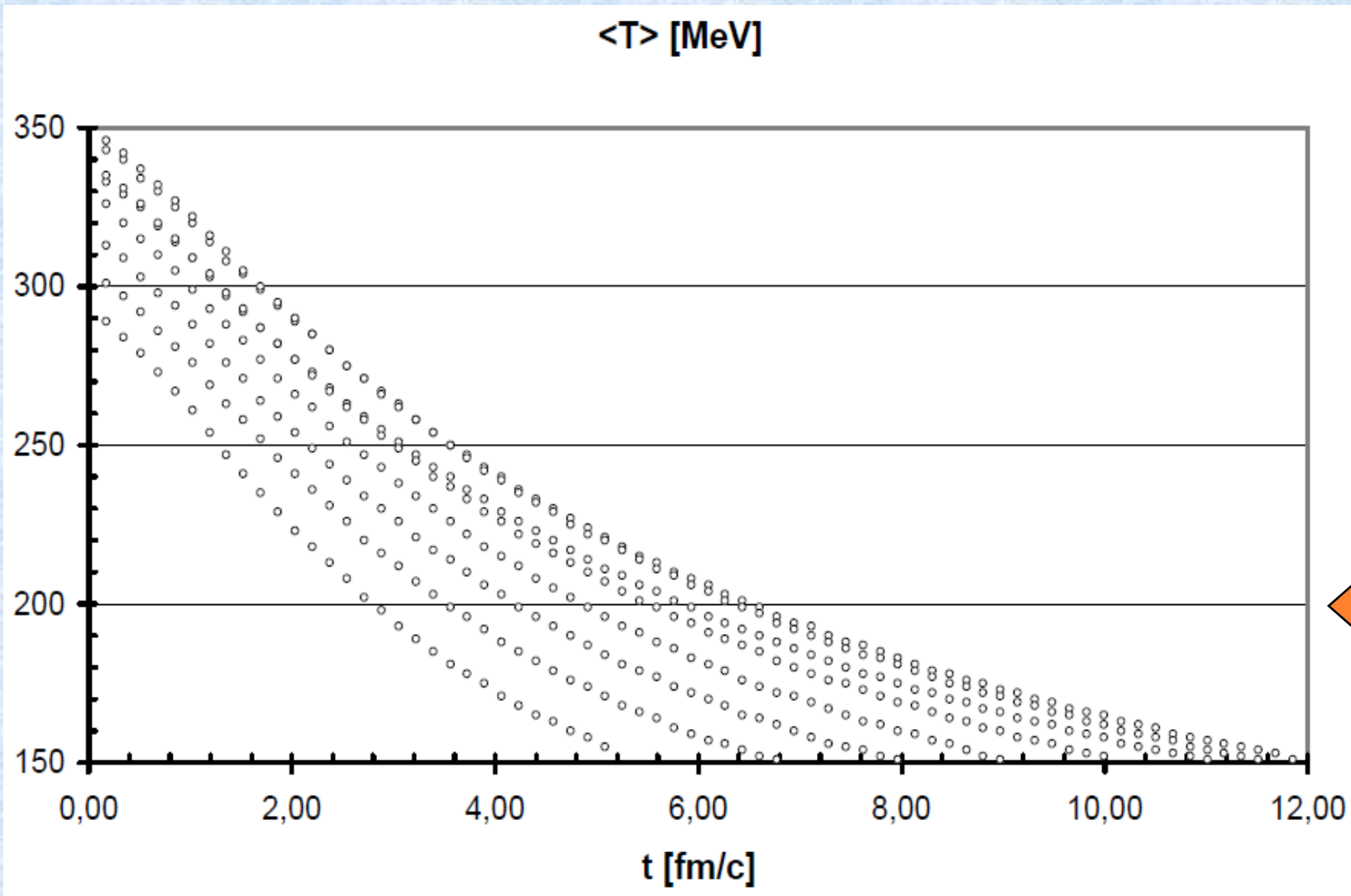
## Multi Module Modeling

- Initial state. YM flux tube
- Fluid dynamics
- Freeze out or Transition to MD
- NCQ scaling

Initial state by V. Magas, L.P. Csernai and D. Strottman  
Phys. Rev. C64 (01) 014901

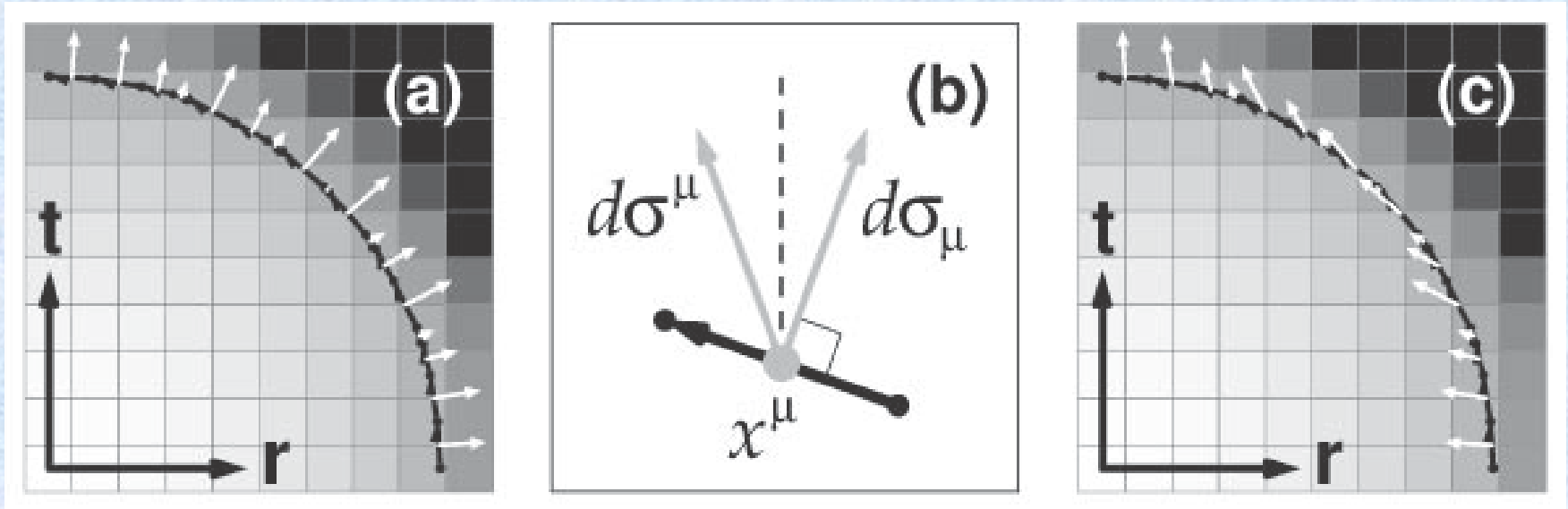
- Au+Au 65+65 A GeV,  $b = 70\%$  of  $b_{\text{max}}$
- Lagrangian fluid cells, moving,  $\sim 5$  mill.
- MIT Bag model EoS
- FO at  $T \sim 200$  MeV, but calculated much longer, until pressure is zero for 90% of the cells.
- Structure and asymmetries of initial state are maintained in nearly perfect expansion.
- Spatially tilted at FO, 3<sup>rd</sup> Flow component!





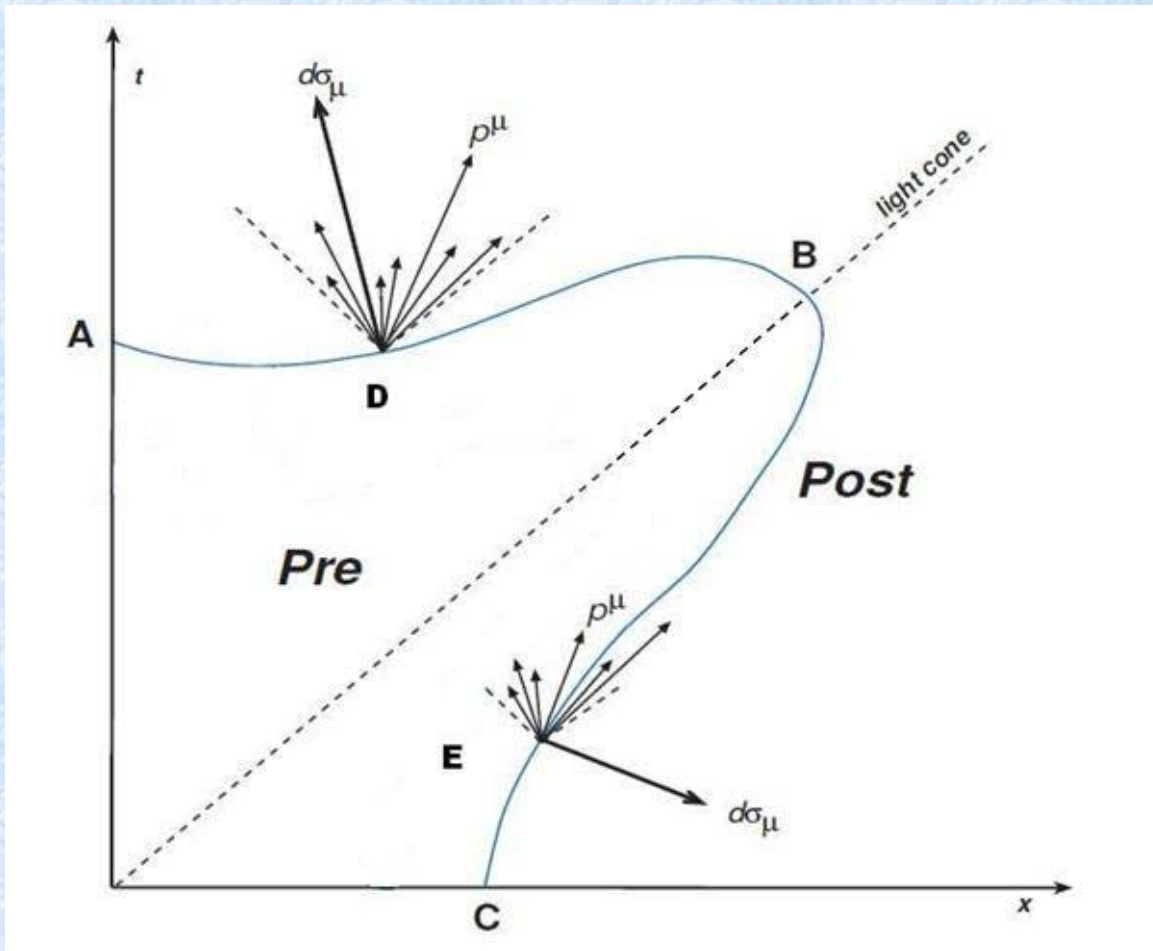
**Average temperature versus time in Au+Au collisions at 65+65 AGeV, for impact parameters,  $b = 0, 0.1, 0.2, \dots, 0.7 b_{\max}$  from the top (0.0) down (0.7).**

# Numerical Extraction of hyper-surface



- (a) gray-level  $\rightarrow$  temperature evolution, contra-variant FO contour vector  $x^\mu$  (black) and covariant normal vector  $d\sigma_\mu$  (white)
- (b) Covariant and contra-variant normal vectors  $d\sigma_\mu$  and  $d\sigma^\mu$
- (c) Contra-variant FO contour vector  $x^\mu$  (black), and contra-variant normal vector  $d\sigma^\mu$  (white)  $\rightarrow$  Not all contra-variant normal vectors point to the exterior direction.

# Conservation laws across hyper-surface



$$[N^\mu d\sigma_\mu] = 0;$$

$$[T^{\mu\nu} d\sigma_\mu] = 0;$$

$$[S^\mu d\sigma_\mu] \geq 0,$$

Equilibrate states:

Pre FO  $\rightarrow$  Post FO

$$j = N^\mu d\sigma_\mu$$

$$A^\mu = T^{\mu\nu} d\sigma_\nu$$

# FD to MD Transition

- From Fluid dynamical model to Molecular dynamical model

Cooper-Frye formula:

$$E \frac{dN}{d^3 p} = \int_{\sigma} f(x, p) p^{\mu} d\sigma_{\mu}$$

Generating the phase space distribution of non-interacting particles

In local equilibrium

When time-like FO surface  $p^{\mu} d\sigma_{\mu} > 0$

Need an anisotropic distribution :

$$f(x, p) \longrightarrow \Theta(p^{\nu} \Lambda_{\nu}) f(p) > 0$$

? When space-like FO surface  $p^{\mu} d\sigma_{\mu} > 0$

Enforcing the conservation laws



# Anisotropic distribution in post FO

- Juttner distribution

RFG (Rest Frame of Gas)

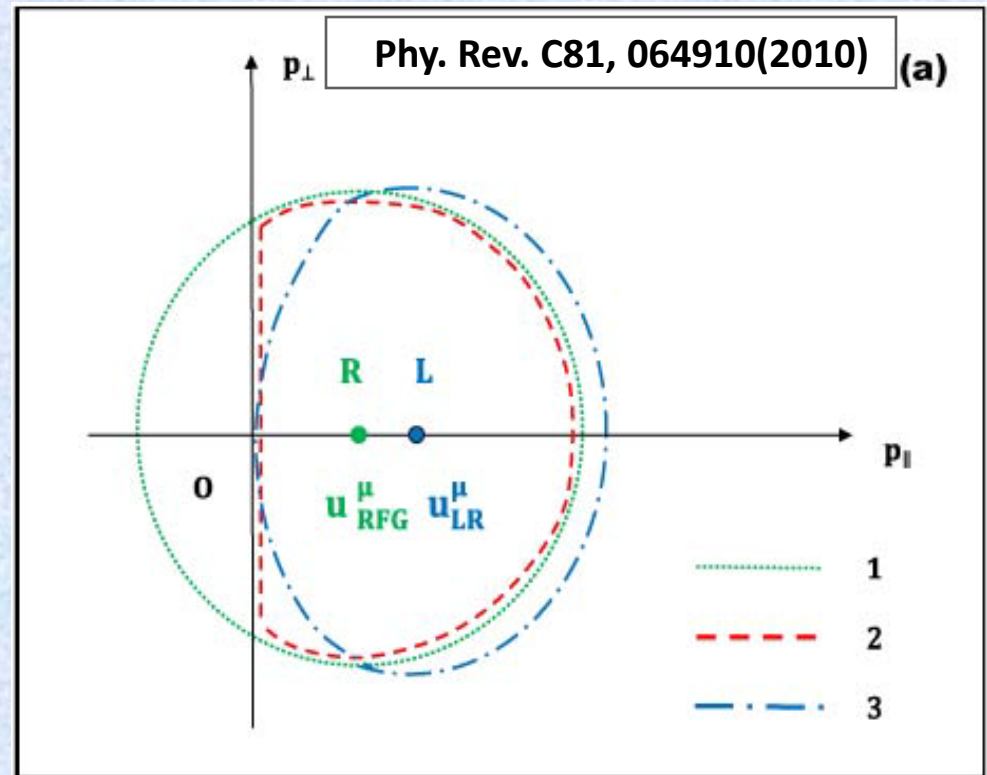
→ starting → “R”  $u_{RFG}^\mu$

- Cut /Canceling Juttner distribution

[Tamosiunas ,2004]

LR (Local Rest Frame)

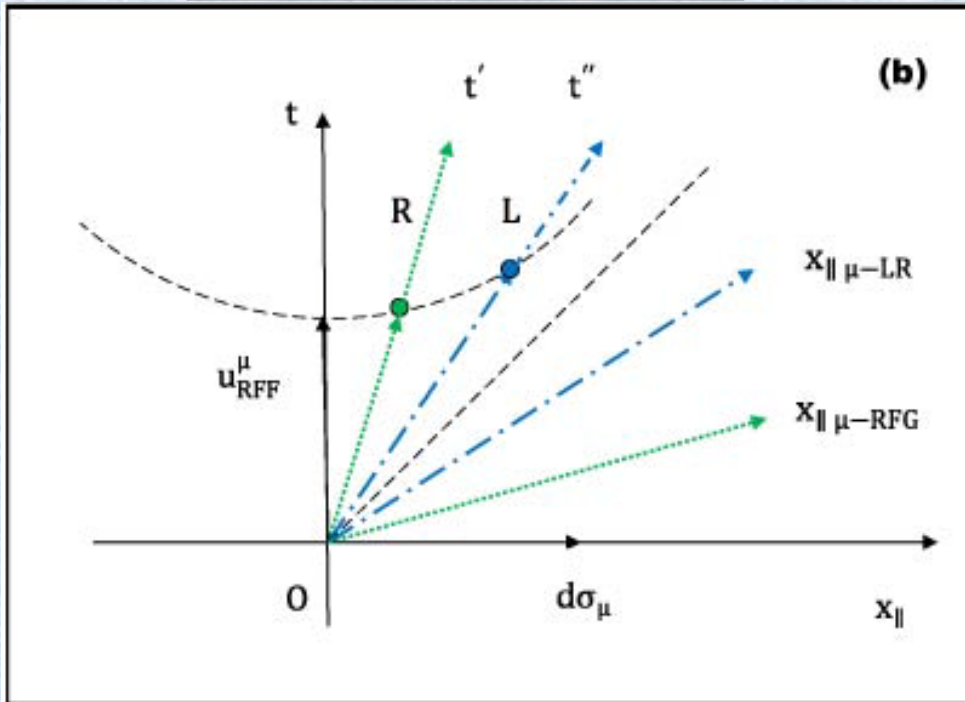
→ final → “L”  $u_{LR}^\mu$



Boost in the orthogonal direction of  $v_{RFG}$  leads to the spatial anisotropic distribution

# Anisotropic distribution in post FO

Phy. Rev. C81, 064910(2010)



## Space-time coordinate

Rest Frame of Front (RFF)  $\rightarrow$  front of cut/canceling  $\rightarrow$  orthogonal to the normal surface  $d\sigma_{\mu}$

Rest Frame of Gas (RFG)  $\rightarrow$  "R"  $\rightarrow$  Starting Isotropic/uncut  $\rightarrow$  "t' "

Local Rest Frame (LR)  $\rightarrow$  "L"  $\rightarrow$

Final cut/canceling  $\rightarrow$  "t" "

# Basic conventions

- Energy momentum tensor, for perfect fluid:

$$T^{\mu\nu} = (e + P)u^\mu u^\nu - P g^{\mu\nu}$$

- FO hyper-surface normal vector:

$$\left. \begin{aligned} d\sigma^\mu &= (\sigma_t, \sigma_x, \sigma_y, \sigma_z) = (\sigma_t, \sigma) \quad \text{and} \quad d\sigma_\mu = (\sigma_t, -\sigma_x, -\sigma_y, -\sigma_z) = (\sigma_t, -\sigma) \\ d\hat{\sigma}^\mu &= \gamma_\sigma(1, s_x, s_y, s_z) = \gamma_\sigma(1, \mathbf{s}) \quad \text{and} \quad d\hat{\sigma}_\mu = \gamma_\sigma(1, -s_x, -s_y, -s_z) = \gamma_\sigma(1, -\mathbf{s}) \end{aligned} \right\} \begin{aligned} d\hat{\sigma}^\mu &\equiv d\sigma^\mu / D \\ \gamma_\sigma^2 &= \pm \frac{1}{1 - \mathbf{s}^2} \end{aligned}$$

- Conserved energy momentum current, baryon current:

$$A^\mu = T^{\mu\nu} d\sigma_\nu = w u^\mu u^\nu d\sigma_\nu - P g^{\mu\nu} d\sigma_\nu = w \gamma^2 (\sigma_t - (\mathbf{v} \cdot \sigma)) \mathbf{v} - P \sigma$$

$$j \equiv N^\mu d\sigma_\mu = n u^\mu d\sigma_\mu = n \gamma (\sigma_t - (\mathbf{v} \cdot \sigma))$$

- $A^\mu \rightarrow$  Taking its norm, taking its projection to normal vector:

↓

$$A^\mu A_\mu$$

↓

$$A^\mu d\sigma_\mu$$

# Simple solution for equilibrated final state

Simple solution for the final state in equilibrium:

$$\begin{aligned}
 \mathbf{A}^\mu \mathbf{A}_\mu &= w^2 (u^\mu d\sigma_\mu)^2 + P^2 (d\sigma^\mu d\sigma_\mu) - 2Pw (u^\mu d\sigma_\mu) u_\mu g^{\mu\nu} d\sigma_\nu \\
 &= w(e - P) (u^\mu d\sigma_\mu)^2 + P^2 (d\sigma^\mu d\sigma_\mu) , \\
 \mathbf{A}^\mu d\sigma_\mu &= w (u^\mu d\sigma_\mu)^2 - P (d\sigma^\mu d\sigma_\mu) .
 \end{aligned}$$



$$\mathbf{A}^\mu \mathbf{A}_\mu = (e - P) \mathbf{A}^\mu d\sigma_\mu + e P (d\sigma^\mu d\sigma_\mu)$$

+ EoS

Special case for baryon  
free matter with EoS as:

$$P = e/3$$

$$d\hat{\sigma}^\mu d\hat{\sigma}_\mu e^2 + 2 a^\mu d\hat{\sigma}_\mu e - 3 a^\mu a_\mu = 0$$

$$u^\mu = d\hat{\sigma}^\mu \qquad a^\mu = \mathbf{A}^\mu / D$$

# Simple solution for equilibrated final state

Case for finite baryon current, and with an EoS as:  $P=P(n,e)$

$$A^\mu A_\mu = (e - P)A^\mu d\sigma_\mu + e P (d\sigma^\mu d\sigma_\mu)$$

$$j = n(u^\mu d\sigma_\mu)$$

$$A^\mu d\sigma_\mu = w(u^\mu d\sigma_\mu)^2 - P(d\sigma^\mu d\sigma_\mu)$$

$$\frac{e + P}{n^2} = \frac{1}{j^2} [A^\mu d\sigma_\mu + P(d\sigma^\mu d\sigma_\mu)]$$

This equation is the same as the generalized Rayleigh line, [Taub, 1949 & Csernai 1987] (and Taub adiabat):

$$j^2 = \frac{[P](d\sigma^\mu d\sigma_\mu)}{[X]}$$

$$[P] = \frac{[(e + P)X]}{X_1 + X_0}$$

$$X = \frac{e + P}{n^2}$$

Space-like FO is weak at RHIC /LHC, but important at FAIR!

# Simple solution for anisotropic final state

- Anisotropic distribution  $\rightarrow$  anisotropic pressure tensor  $\rightarrow$  pressure components are non-identical

$$T^{\mu\nu} = \text{diag}(e, P, P, P)|_{LR} \longrightarrow T^{\mu\nu} = \text{diag}(e, P_{\parallel}, P_{\perp}, P_{\perp})|_{LR}$$

Isotropic momentum dist.  $\rightarrow$  anisotropic momentum dist.

Perfect  $T^{\mu\nu}$   $\longrightarrow$  
$$T^{\mu\nu} = e u_{LR}^{\mu} u_{LR}^{\nu} - P_{\perp} \Delta_{LR}^{\mu\nu} + (P_{\parallel} - P_{\perp}) \hat{F}^{\mu} \hat{F}^{\nu}$$

Orthogonal projector to  $u_{LR}^{\mu}$

$$\Delta_{LR}^{\mu\nu} = g^{\mu\nu} - u^{\mu} u^{\nu}$$

Unit vector projection of  $d\hat{\sigma}^{\mu}$   
Orthogonal to  $u_{LR}^{\mu}$

$$\hat{F}^{\mu} = C \Delta^{\mu\nu} d\hat{\sigma}_{\nu}$$

# Simple solution for anisotropic final state

$$T^{\mu\nu} = e u_{\text{LR}}^\mu u_{\text{LR}}^\nu - P_\perp \Delta_{\text{LR}}^{\mu\nu} + (P_\parallel - P_\perp) \hat{F}^\mu \hat{F}^\nu$$

$$\begin{aligned} A_0^\mu A_{0\mu} &= w_\perp (e - P_\perp) (u^\mu d\sigma_\mu)^2 + (P_\perp)^2 (d\sigma^\mu d\sigma_\mu) \\ &\quad - (P_\parallel - P_\perp) (P_\parallel + P_\perp) (\hat{F}^\nu d\sigma_\nu)^2, \\ A_0^\mu d\sigma_\mu &= w_\perp (u^\mu d\sigma_\mu)^2 - P_\perp (d\sigma^\mu d\sigma_\mu) \\ &\quad + (P_\parallel - P_\perp) (\hat{F}^\nu d\sigma_\nu)^2. \end{aligned}$$

Escape probability relative to degree of anisotropic

$$\begin{aligned} A_0^\mu A_{0\mu} &= (e - P_\perp) A_0^\mu d\sigma_\mu + e P_\perp (d\sigma^\mu d\sigma_\mu) \\ &\quad - (e + P_\parallel) (P_\parallel - P_\perp) (\hat{F}^\nu d\sigma_\nu)^2. \end{aligned}$$

Simple covariant solution

← Eos +

# Summary

- Initial state is important
- Anisotropic distribution, more realistic especially for space-like transitions
- Enforcing the conservation laws
- Transparent and covariant method
- Connecting PIC-hydro and PACIAE is in progress



Thanks for your attention!