

Kinetics of Charm in Medium

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Stiftung / Foundation



Outline

Heavy-quark rescattering in QGP

Non-photonic e^\pm Observables: v_2 and R_{AA}

Conclusions and Outlook

Motivation

- ▶ Measured p_T spectra and v_2 of non-photonic single electrons
- ▶ coalescence model describes data under assumption of thermalized c quarks, flowing with the bulk medium

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- ▶ coalescence model describes data under assumption of thermalized c quarks, flowing with the bulk medium
- ▶ What is the underlying microscopic mechanism for thermalization?
 - ▶ pQCD elastic HQ scattering: need unrealistically large α_s [Moore, Teaney '04]
 - ▶ Gluon-radiative energy loss: need to enhance transport coefficient \hat{q} by large factor [Armesto et al '05] or enhanced gluon density [Djordjevic, Gyulassi et al '05]
 - ▶ including pQCD elastic scattering: still not enough equilibration of heavy quarks [Wicks et al '05]

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 - ▶ including pQCD elastic scattering: still not enough equilibration of heavy quarks [Wicks et al '05]
- ▶ Assumption: survival of D - and B -meson resonances in the sQGP
- ▶ facilitates elastic heavy-quark rescattering

Free Lagrangian: Particle Content

- ▶ Chiral symmetry $SU_V(2) \otimes SU_A(2)$ in light-quark sector of QCD

$$\mathcal{L}_D^{(0)} = \sum_{i=1}^2 [(\partial_\mu \Phi_i^\dagger)(\partial^\mu \Phi_i) - m_D^2 \Phi_i^\dagger \Phi_i] + \text{massive (pseudo-)vectors } D^*$$

- ▶ Φ_i : two doublets: pseudo-scalar $\sim \begin{pmatrix} D^0 \\ D^- \end{pmatrix}$ and scalar
- ▶ Φ_i^* : two doublets: vector $\sim \begin{pmatrix} D^{0*} \\ D^{-*} \end{pmatrix}$ and pseudo-vector

$$\mathcal{L}_{qc}^{(0)} = \bar{q} i \not{\partial} q + \bar{c} (i \not{\partial} - m_c) c$$

- ▶ q : light-quark doublet $\sim \begin{pmatrix} u \\ d \end{pmatrix}$
- ▶ c : singlet

Interactions

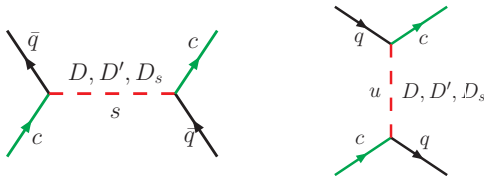
- ▶ Interactions determined by **chiral** symmetry
- ▶ For transversality of vector mesons:
heavy-quark effective theory vertices

$$\begin{aligned} \mathcal{L}_{\text{int}} = & -G_S \left(\bar{q} \frac{1 + \not{v}}{2} \Phi_1 c_v + \bar{q} \frac{1 + \not{v}}{2} i\gamma^5 \Phi_2 c_v + h.c. \right) \\ & -G_V \left(\bar{q} \frac{1 + \not{v}}{2} \gamma^\mu \Phi_{1\mu}^* c_v + \bar{q} \frac{1 + \not{v}}{2} i\gamma^\mu \gamma^5 \Phi_{2\mu}^* c_v + h.c. \right) \end{aligned}$$

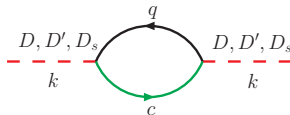
- ▶ v : four velocity of heavy quark
- ▶ in **HQET**: spin symmetry $\Rightarrow G_S = G_V$

Resonance Scattering

- ▶ elastic heavy-light-(anti-)quark scattering



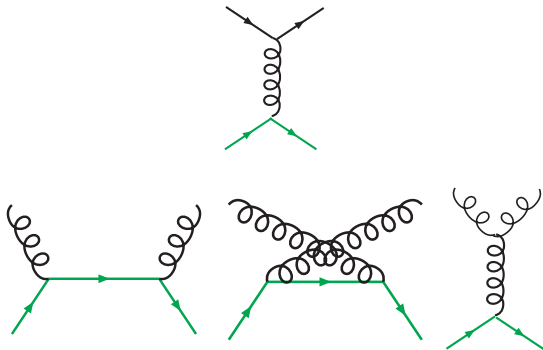
- ▶ D - and B -meson like resonances in sQGP



- ▶ parameters
 - ▶ $m_D = 2 \text{ GeV}$, $\Gamma_D = 0.4 \dots 0.75 \text{ GeV}$
 - ▶ $m_B = 5 \text{ GeV}$, $\Gamma_B = 0.4 \dots 0.75 \text{ GeV}$

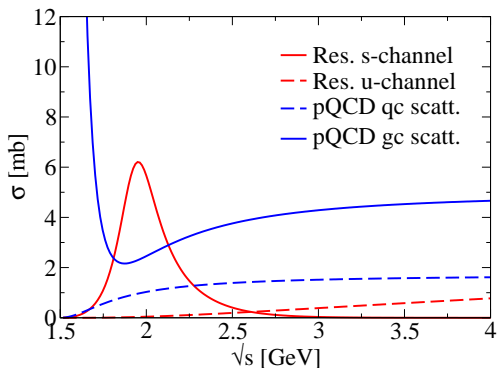
Contributions from pQCD

- ▶ Lowest-order matrix elements (Cambridge '79)



- ▶ In-medium **Debye-screening mass** for t -channel gluon exchange:
 $\mu_g = gT$, $\alpha_s = 0.4$

Cross sections



- ▶ total pQCD and resonance cross sections: comparable in size
- ▶ BUT pQCD forward peaked \leftrightarrow resonance isotropic
- ▶ resonance scattering more effective for friction and diffusion

The Fokker-Planck Equation

- ▶ Neglect **flow** of the medium (only for the moment!)
- ▶ heavy particle (**c,b quarks**) in a **heat bath** of light particles (QGP)

$$\frac{\partial f(t, \vec{p})}{\partial t} = \frac{\partial}{\partial p_i} \left[p_i A(t, p) + \frac{\partial}{\partial p_j} B_{ij}(t, \vec{p}) \right] f(t, \vec{p})$$

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- ▶ A and B_{ij} given by **averages** with matrix elements (cross sections) from **resonance model**

$$\langle X(\vec{p}') \rangle = \frac{1}{\gamma_c} \frac{1}{2E_p} \int \frac{d^3\vec{q}}{(2\pi)^3 2E_q} \int \frac{d^3\vec{q}'}{(2\pi)^3 2E_{q'}} \int \frac{d^3\vec{p}''}{(2\pi)^3 2E_{p''}} \sum |\mathcal{M}|^2 (2\pi)^4 \delta^{(4)}(p + q - p' - q') \hat{f}(\vec{q}) X(\vec{p}'')$$

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- ▶ $A(t, \vec{p})$ friction (drag) coefficient = $1/\tau_{\text{eq}}$
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- ▶ B_{ij} : time scale for momentum fluctuations

$$\begin{aligned} B_{ij}(t, \vec{p}) &= \frac{1}{2} \langle (p_i - p'_i)(p_j - p'_j) \rangle \\ &= B_0(t, p) \left(\delta_{ij} - \frac{p_i p_j}{p^2} \right) + B_1(t, p) \frac{p_i p_j}{p^2} \end{aligned}$$

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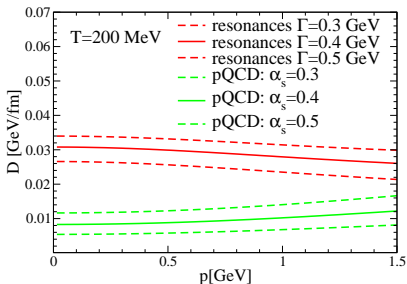
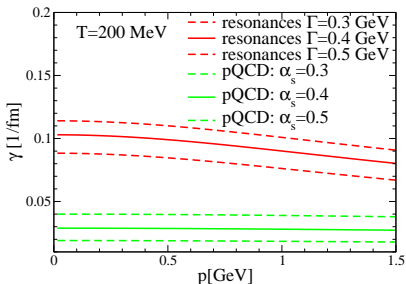
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- ▶ to ensure correct equilibrium limit: $B_1(t, p) = T(t) E_p A(t, p)$
(Einstein dissipation-fluctuation relation)

Drag and Diffusion: pQCD vs. resonance scattering

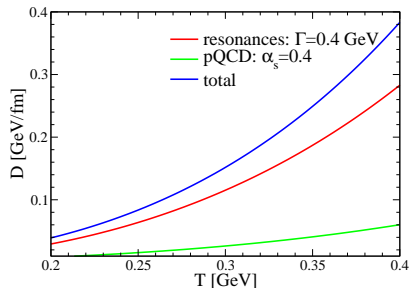
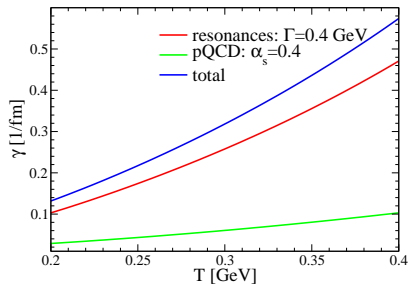
- ▶ three-momentum dependence



- ▶ resonance contributions factor $\sim 2 \dots 3$ higher than pQCD!

The Coefficients: pQCD vs. resonance scattering

► Temperature dependence



Time evolution of the fire ball

- ▶ Elliptic **fire-ball** parameterization
fitted to hydrodynamical flow pattern [Kolb '00]

$$V(t) = \pi(z_0 + v_z t)a(t)b(t), \quad a, b: \text{half-axes of ellipse,}$$

$$v_{a,b} = v_\infty[1 - \exp(-\alpha t)] \mp \Delta v[1 - \exp(-\beta t)]$$

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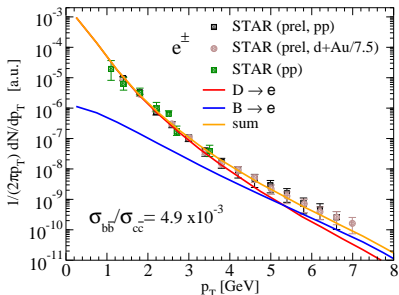
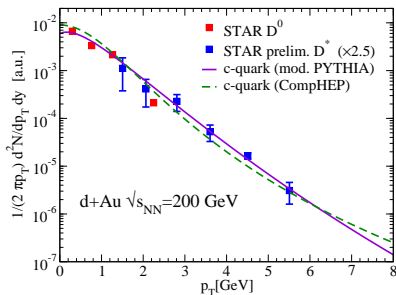
- ▶ **Isentropic expansion**: $S = \text{const}$ (fixed from N_{ch})
- ▶ **QGP Equation of state**:

$$s = \frac{S}{V(t)} = \frac{4\pi^2}{90} T^3 (16 + 10.5n_f^*), \quad n_f^* = 2.5$$

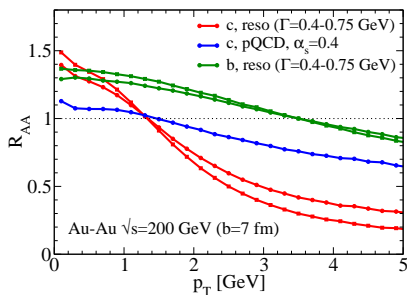
- ▶ obtain $T(t) \Rightarrow A(t, p)$, $B_0(t, p)$ and $B_1 = TEA$
- ▶ for semicentral collisions ($b = 7$ fm): $T_0 = 340$ MeV,
 QGP lifetime $\simeq 5$ fm/ c .
- ▶ simulate FP equation as **relativistic Langevin process**

Initial conditions

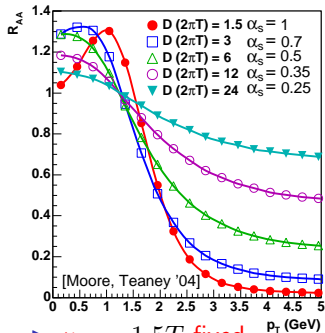
- ▶ need initial p_T -spectra of **charm** and **bottom** quarks
 - ▶ (modified) PYTHIA to describe exp. **D** meson spectra, assuming **δ -function fragmentation**
 - ▶ exp. **non-photonic single- e^\pm spectra**: Fix bottom/charm ratio



Spectra and elliptic flow for heavy quarks

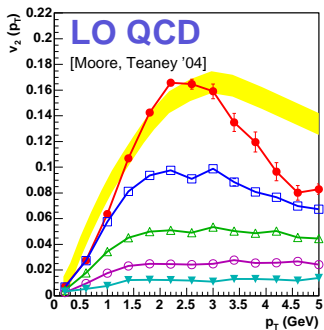
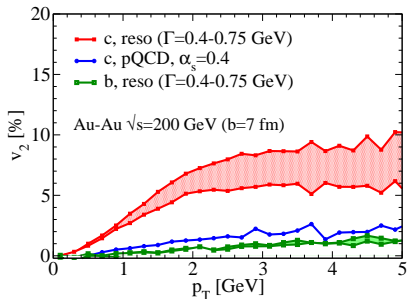


- ▶ $\mu_D = gT$, $\alpha_s = g^2/(4\pi) = 0.4$
- ▶ resonances \Rightarrow c-quark thermalization **without upscaling of cross sections**
- ▶ Fireball parametrization consistent with hydro



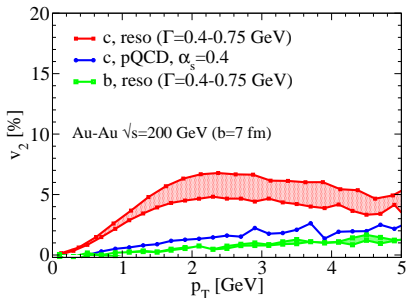
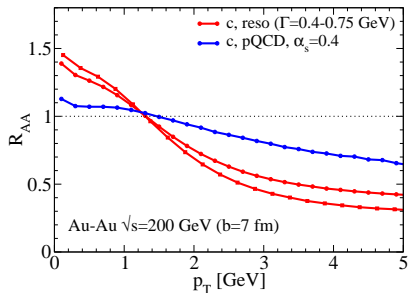
- ▶ $\mu_D = 1.5T$ fixed
- ▶ $2\pi TD \simeq \frac{3}{2\alpha_s^2}$

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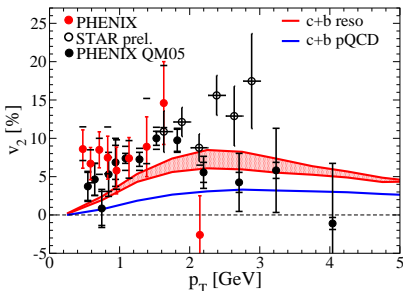
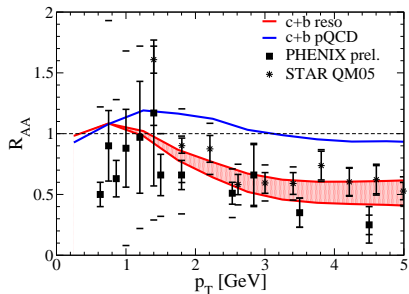
Spectra and elliptic flow for heavy quarks

With **form-factor vertices** instead of point vertices ($\Lambda = 1$ GeV)



Observables: p_T -spectra (R_{AA}), v_2

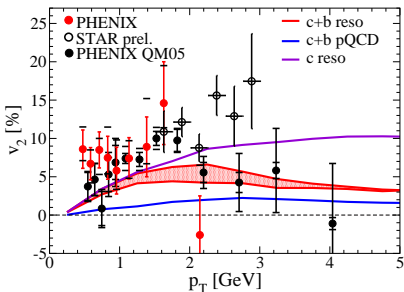
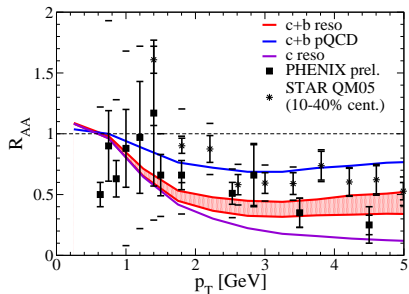
- ▶ **Hadronization: Coalescence** with light quarks (fixed before [Greco et al 03]) + **fragmentation** ($c\bar{c}$, $b\bar{b}$ conserved)
- ▶ single electrons from decay of D - and B -mesons



- ▶ Without further adjustments: data quite well described

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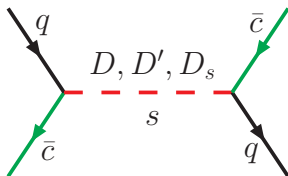


How to check resonance assumption?

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- ▶ dominant channel: quark-**anti- c -quark** s channel



- ▶ **CBM@FAIR**: quark dominated $\Rightarrow \bar{c}$ quarks most affected
- ▶ thermalization effects more pronounced for \bar{D} (D^-) than for D (D^+) mesons!

Conclusions and Outlook

- ▶ Assumption: survival of **resonances** in the (s)QGP
- ▶ **nonperturbative re-interactions** of heavy quarks in QGP
- ▶ **Observables** via Langevin approach and coalescence+fragmentation
 - ▶ **Elastic resonance scattering** $\Rightarrow R_{AA}^{(c)} \simeq 0.2, v_2^{(c)} \simeq 0.1$
without upscaling of cross sections
 - ▶ small effects on **bottom quarks**
 - ▶ **Heavy-light quark coalescence** enhances $v_2^{(e)}$ and R_{AA} for
 $p_T \simeq 2$ GeV
 - ▶ **bottom** dominates for $p_T > 3.5$ GeV \Rightarrow reduced suppression, $v_2^{(e)}$

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- ▶ Further investigations
 - ▶ improved (softer) fragmentation
 - ▶ implementation of gluon-radiation processes
 - ▶ quantitative consequences for quarkonia