

Coupled-charge transport: current status of fluid dynamic modeling

CRC-Tr 211 Transport Meeting
Frankfurt

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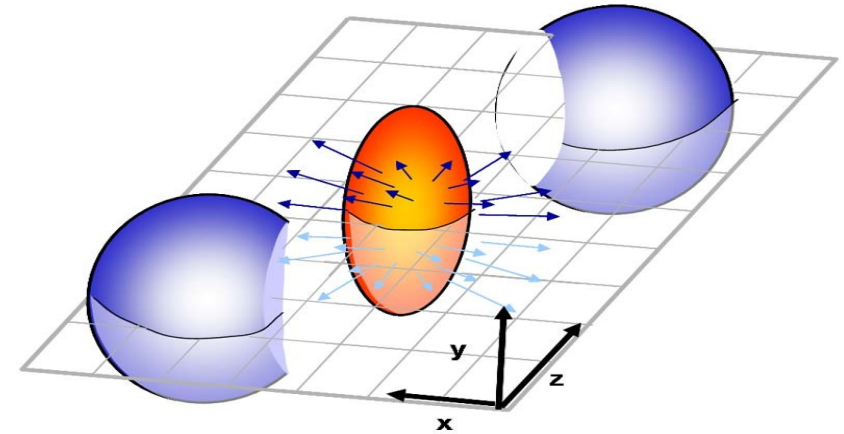
*Jakob Lohr, David Wagner, Semen Potesnov, Harri Niemi, Etele Molnár,
Dirk Rischke, Carsten Greiner*

The multi-component nature of nuclear matter

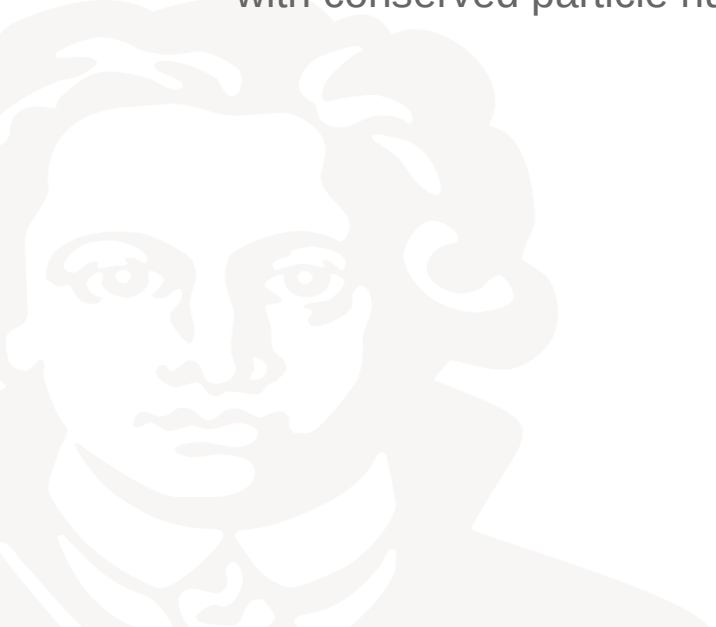
Traditionally:

Viewed as 'blob' of **one type of matter** (single component) with **one velocity field**

- usually 'blob' of energy
with conserved particle number



<https://www.quantumdiaries.org/wp-content/uploads/2011/02/FlowPr.jpg>

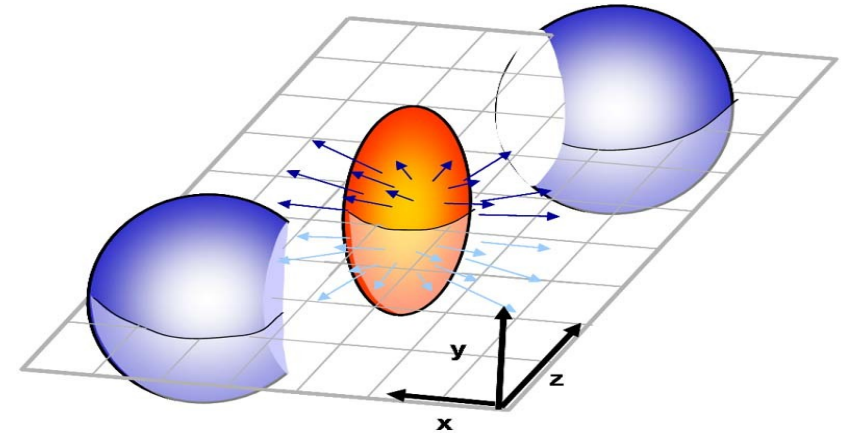


The multi-component nature of nuclear matter

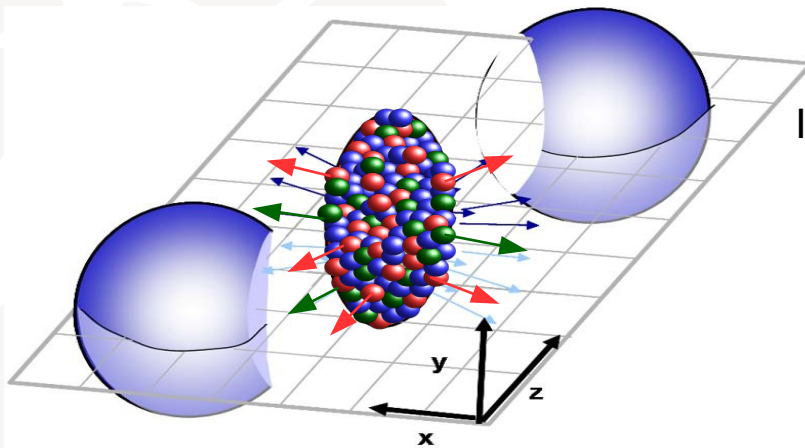
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In general:

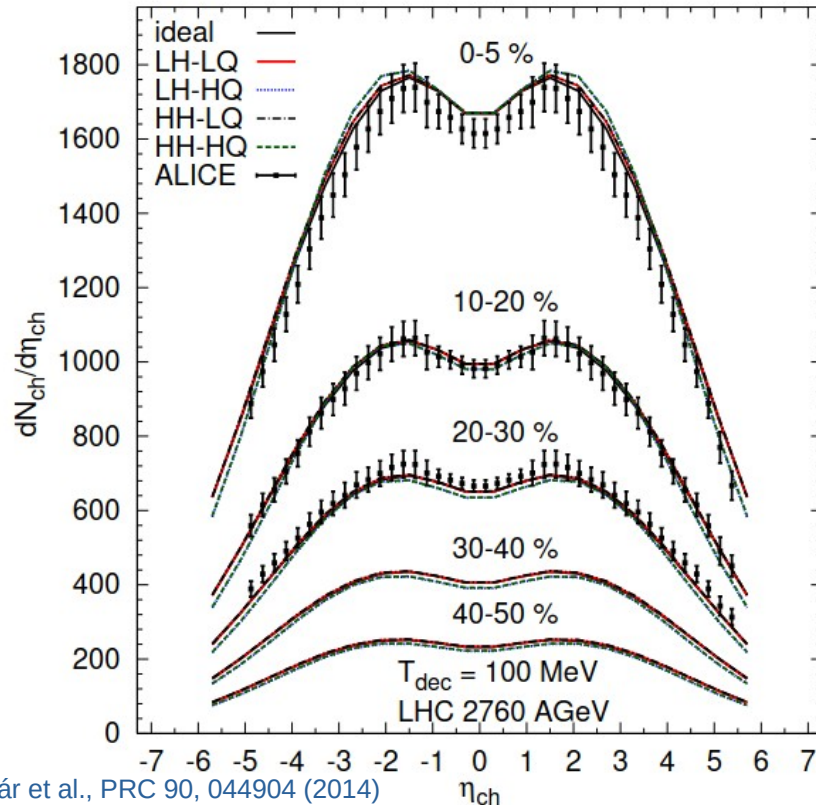
Consists of **multiple components** with various properties with **multiple velocity fields**

- with **multiple conserved quantities** (e.g. energy, electric charge, baryon number, strangeness, ...)

- mixed chemistry → **coupled charge currents!**

Hydrodynamics applied to heavy ion collisions

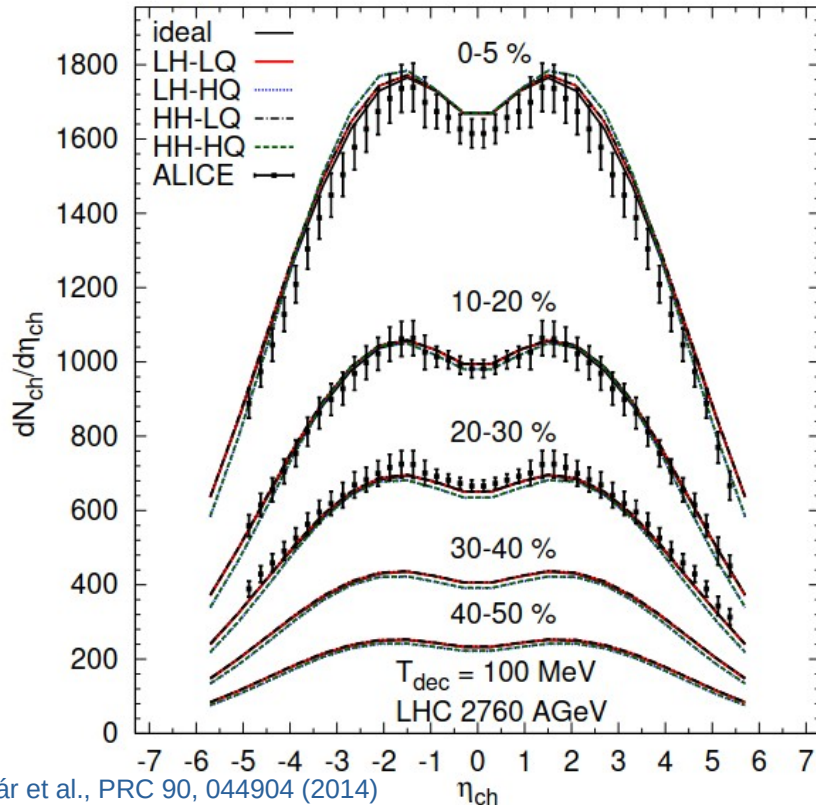
Hydro has been very successful in the **effective** description of the evolution of heavy-ion collisions ...



Molnár et al., PRC 90, 044904 (2014)

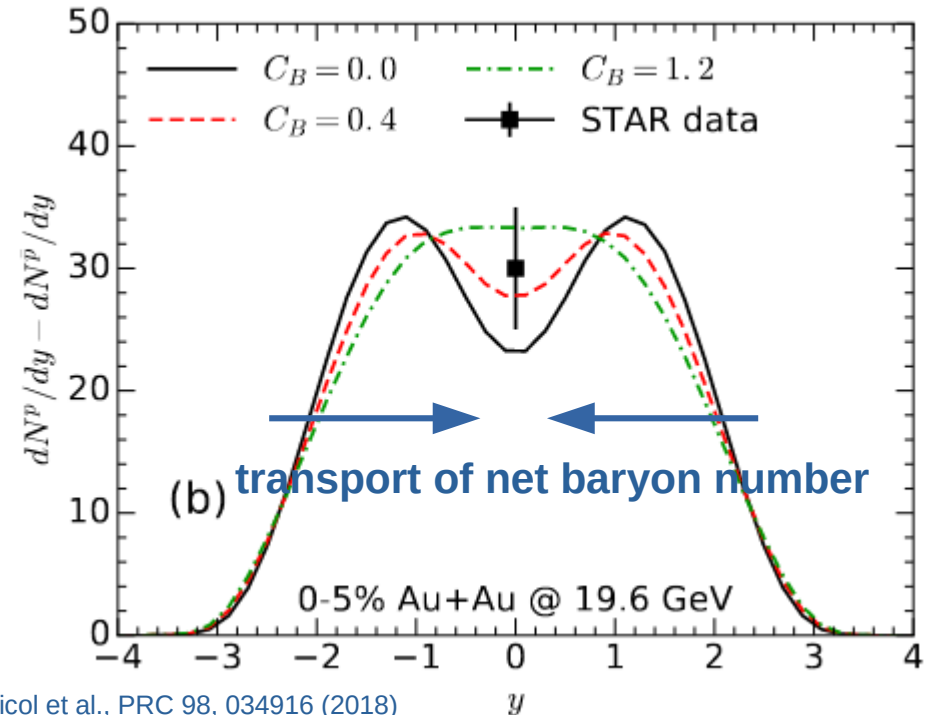
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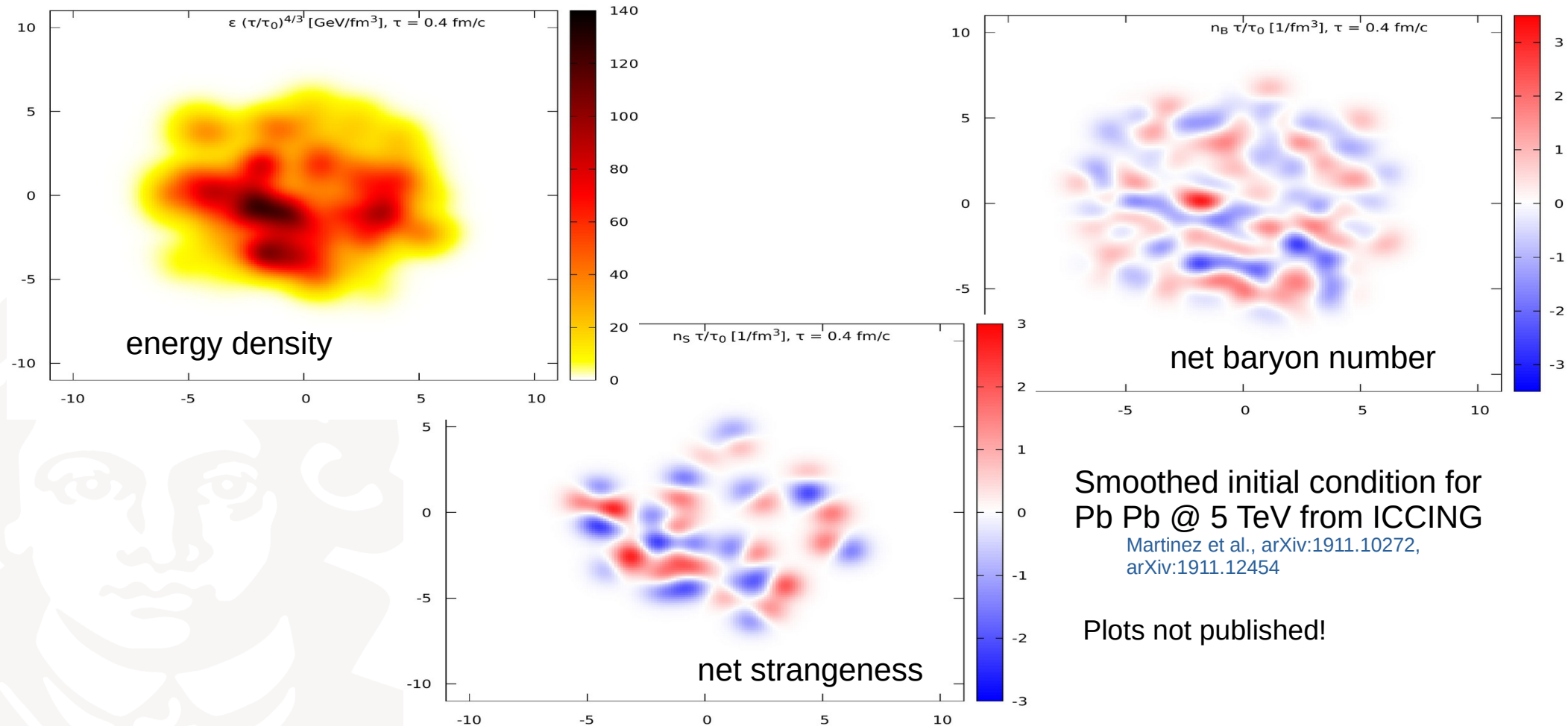
Molnár et al., PRC 90, 044904 (2014)

... but the theory **has to be extended** in order to explicitly account for effects and is dependent on the knowledge of the **underlying microscopic properties**.



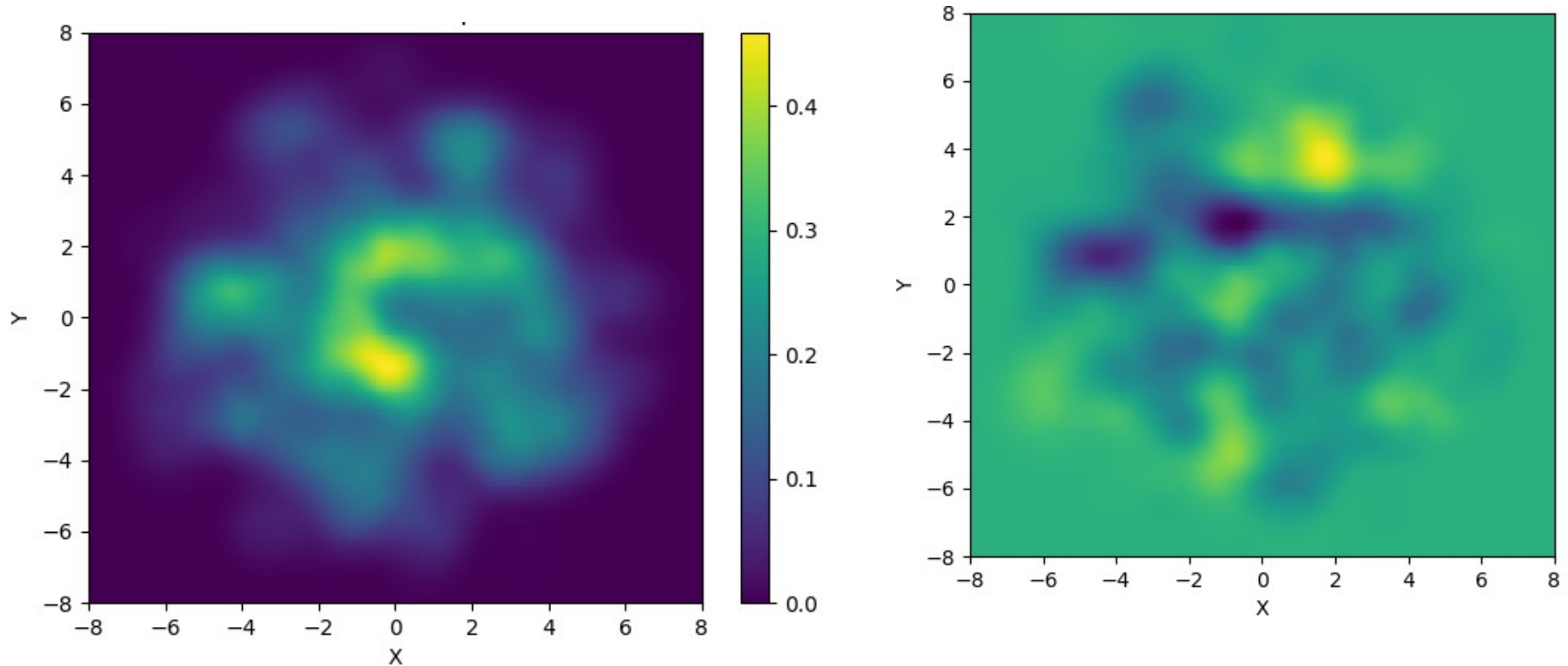
Denicol et al., PRC 98, 034916 (2018)

Initial state with multiple conserved charges



Initial state with multiple conserved charges

Collaborations with T. Dore, O. Garcia-Montero, S. Schlicht (**McDipper / KoMPoST**)
and H. Roch, N. Götz, H. Elfner (**SMASH**)



Fluid dynamics of multi-component systems

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

Conservation of charge: $\partial_\mu N_q^\mu = 0$



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$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

q-th conserved charge (eg. B,Q,S)

Conservation of charge: $\partial_\mu N_q^\mu = 0$

$$N_q^\mu = \sum_i q_i N_i^\mu = n_q u^\mu + V_q^\mu$$



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2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

$$\begin{aligned} \tau_\Pi \dot{\Pi} + \Pi &= S_\Pi \\ \sum_{q'} \tau_{qq'} \dot{V}_{q'}^{\langle\mu\rangle} + V_q^\mu &= S_q^\mu \\ \tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} &= S_\pi^{\mu\nu} \end{aligned}$$

Relaxation equations
(Israel-Stewart-type
causal theory)

Equations of motion with multiple conserved charges

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Relaxation equations
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$$\begin{aligned} S_q^{\mu} = & \sum_{q'} \kappa_{qq'} \nabla^{\mu} \alpha_{q'} - \sum_{q'} \tau_{qq'} V_{q',\nu} \omega^{\nu\mu} - \sum_{q'} \delta_{VV}^{(q,q')} V_{q'}^{\mu} \theta - \sum_{q'} \lambda_{VV}^{(q,q')} V_{q',\nu} \sigma^{\mu\nu} \\ & - \ell_{V\Pi}^{(q)} \nabla^{\mu} \Pi + \ell_{V\pi}^{(q)} \Delta^{\mu\nu} \nabla_{\lambda} \pi_{\nu}^{\lambda} + \tau_{V\Pi}^{(q)} \Pi \dot{u}^{\mu} - \tau_{V\pi}^{(q)} \pi^{\mu\nu} \dot{u}_{\nu} \\ & + \sum_{q'} \lambda_{V\Pi}^{(q,q')} \Pi \nabla^{\mu} \alpha_{q'} - \sum_{q'} \lambda_{V\pi}^{(q,q')} \pi^{\mu\nu} \nabla_{\nu} \alpha_{q'} \end{aligned}$$

Fotakis, Molnár, Niemi, Greiner, Rischke
PRD 106, 036009 (2022)

Mixed chemistry couples diffusion currents (coupled charge-transport); already present in 1st order term

2nd order terms: couples all currents to each other; depend on all gradients!

Explicit expressions for transport coefficients!

Equations of motion with multiple conserved charges

2nd-order (multi-component) hydro

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Diffusion
coefficients:
extensively
studied!

Fotakis, Molnár, Niemi, Greiner, Rischke
PRD 106, 036009 (2022)

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)
Fotakis, Soloveva et al, PRD 104, 034014 (2021)

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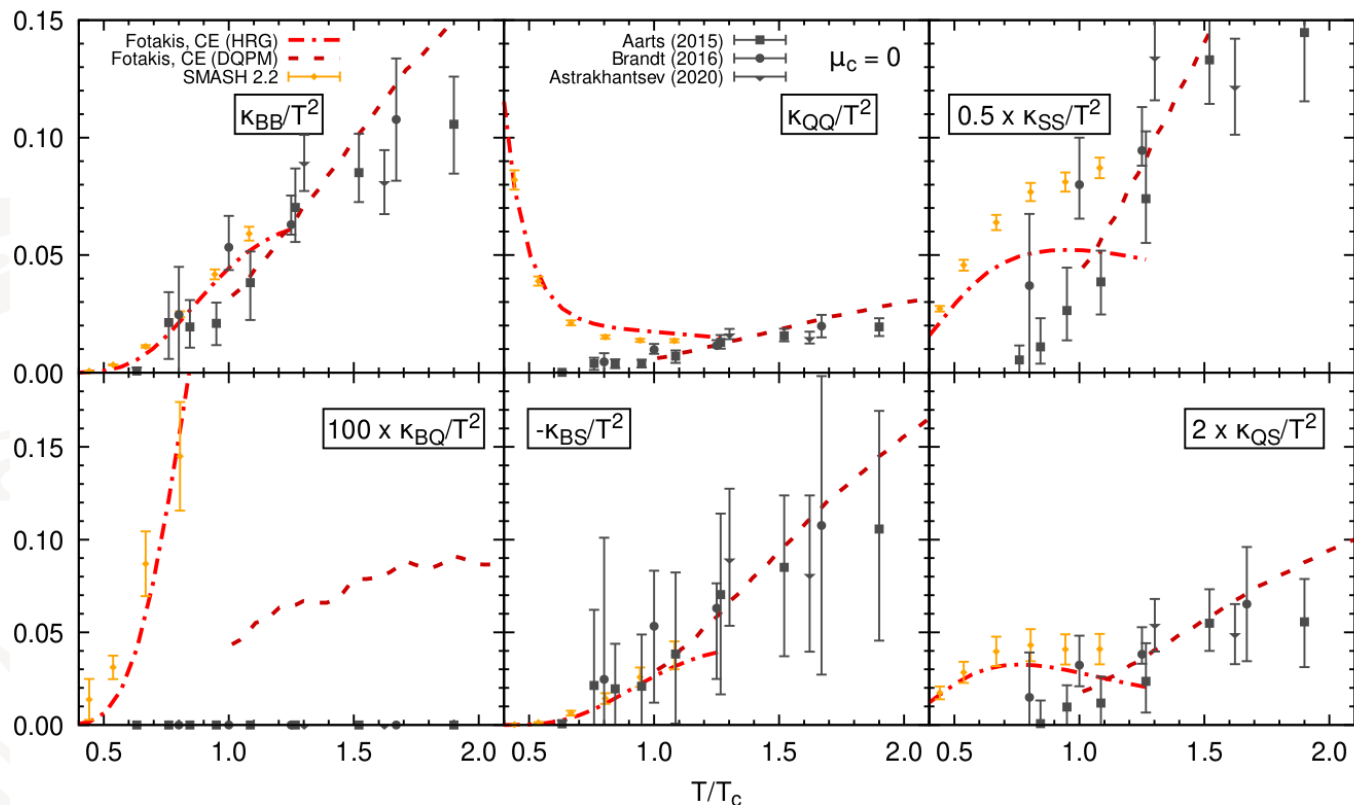
2nd order terms: couples all currents to each other; depend on all gradients!

Explicit expressions for transport coefficients!

Computation of transport coefficients (Example: diffusion coefficients)

$$\kappa_{qq'} = \sum_{n=0}^{\infty} \sum_{i,j=1}^{N_{\text{species}}} (c^{-1})_{ji,0n}^{(1)} q_j \left(q_i' \mathcal{J}_{n+1,1}^{(i)} - \frac{n_{q'}}{\epsilon + P_0} \mathcal{J}_{n+2,1}^{(i)} \right)$$

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)
Fotakis, Soloveva et al, PRD 104, 034014 (2021)
Hammelmann et al., arXiv:2307.15606 (2023)



Insights in **chemical composition** of nuclear matter?

Upcoming publication!
(Fotakis, Lohr, Greiner)

Upcoming project with T. Dore and S. Schlichting

Single-component vs. Multi-component system

What is with the second-order terms?

$$S_q^\mu = (\dots) + \ell_{V\pi}^{(q)} \Delta^{\mu\nu} \nabla_\lambda \pi_\nu^\lambda + (\dots)$$



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Ultrarelativistic, classical system with hard-sphere interactions:

Denicol et al., PRD 85, 114047 (2012)

TABLE I. The coefficients for the particle diffusion for a classical gas with constant cross section in the ultrarelativistic limit, in the 14-moment approximation. The transport coefficient $\tau_{n\pi}$ was incorrectly listed as being zero in Ref. [1]

κ	$\tau_n[\lambda_{\text{mfp}}]$	$\delta_{nn}[\tau_n]$	$\lambda_{nn}[\tau_n]$	$\lambda_{n\pi}[\tau_n]$	$\ell_{n\pi}[\tau_n]$	$\tau_{n\pi}[\tau_n]$
$3/(16\sigma)$	$9/4$	1	$3/5$	$\beta_0/20$	$\beta_0/20$	$\beta_0/80$

Used in simulations of
heavy-ion collisions!

$$\tau_n \dot{V}_q^{\langle\mu\rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'} - V_{q,\nu} \omega^{\nu\mu} - \tau_n V_{q'}^\mu \theta - \frac{3\tau_n}{5} V_{q,\nu} \sigma^{\mu\nu} + \frac{\tau_n}{20T} \Delta^{\mu\nu} \nabla_\lambda \pi_\nu^\lambda - \frac{\tau_n}{20T} \pi^{\mu\nu} \dot{u}_\nu - \frac{\tau_n}{20T} \pi^{\mu\nu} \nabla_\nu \alpha_q$$

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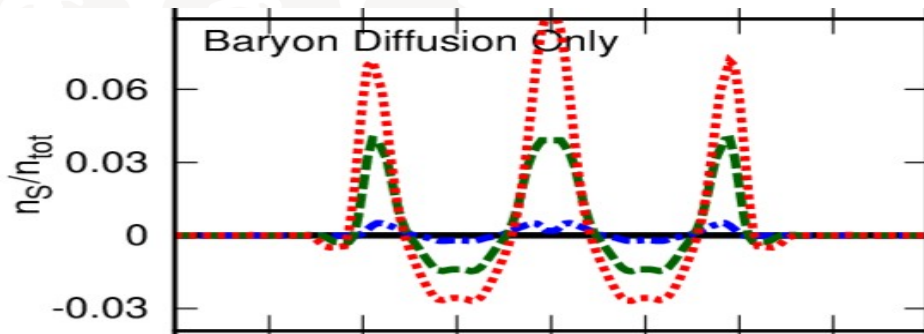
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Second-order transport coefficients
not consistent with assumed system

→ generation of unphysical charge currents

Consistency is important in charge transport!
Use multi-component expressions.

Single-component vs. Multi-component system

What is with the second-order terms?

$$S_q^\mu = (\dots) + \ell_{V\pi}^{(q)} \Delta^{\mu\nu} \nabla_\lambda \pi_\nu^\lambda + (\dots)$$

Ultrarelativistic, classical system with $h \gg \lambda_{\text{mfp}}$

$$\ell_{V\pi}^{(q')} = \frac{9}{80\sigma_{\text{tot}}P} \mathbf{c}_{q'} \xrightarrow{\text{single}} \frac{\beta}{20} \tau_q \quad [14047 (2012)]$$

TABLE I. The coefficients for the 14-moment approximation. [1]

κ	$\tau_n [\lambda_{\text{mfp}}]$
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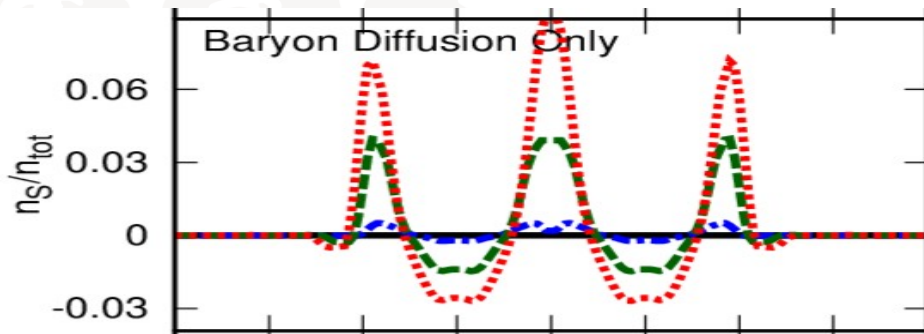
$$\tau_{V\pi}^{(q')} = \ell_{V\pi}^{(q')} \xrightarrow{\text{single}} \frac{\beta}{20} \tau_q,$$

TABLE II. The coefficients for the 14-moment approximation, in the ultrarelativistic limit, in the units of [1]

$\ell_{n\pi}[\tau_n]$	$\tau_{n\pi}[\tau_n]$
$\beta_0/20$	$\beta_0/80$

Used in simulations of heavy-ion collisions!

$$\tau_n \dot{V}_q^{\langle\mu\rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'} - V_{q,\nu} \omega^{\nu\mu} - \tau_n V_{q'}^\mu \theta - \frac{3\tau_n}{5} V_{q,\nu} \sigma^{\mu\nu} + \frac{\tau_n}{20T} \Delta^{\mu\nu} \nabla_\lambda \pi_\nu^\lambda - \frac{\tau_n}{20T} \pi^{\mu\nu} \dot{u}_\nu - \frac{\tau_n}{20T} \pi^{\mu\nu} \nabla_\nu \alpha_q$$



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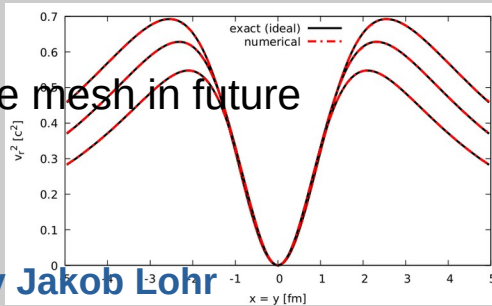
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The journey to a software package treating coupled-charged transport



- (3+1)D-hydro simulation
- multiple-conserved charges
- implemented multi-component theory
- user-defined EoS, transport coefficients, theories and geometries
- maybe adaptive mesh in future

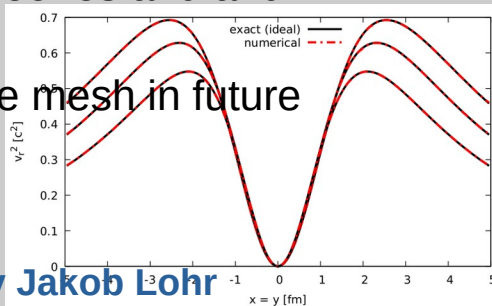


See recent talk by [Jakob Lohr](#)

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See recent talk by Jakob Lohr

TraCoLinR

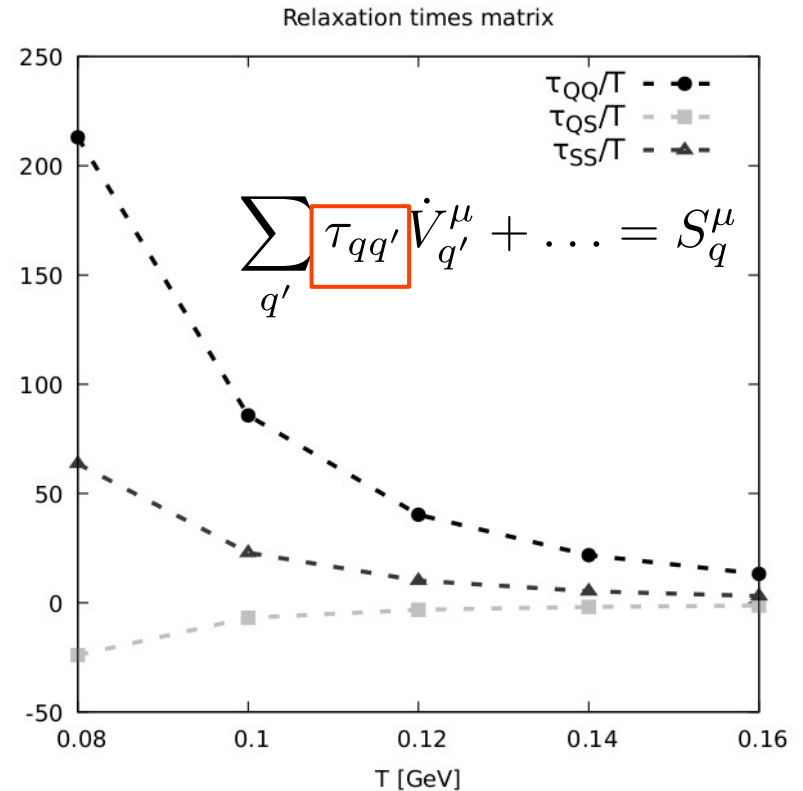
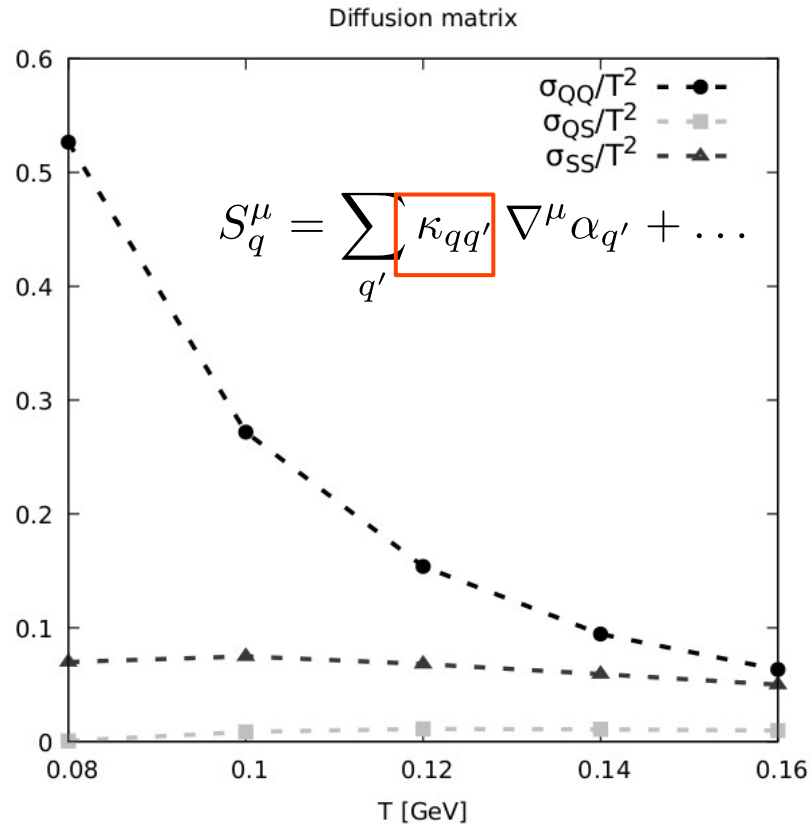
(TRANsport COefficients from LINear Response theory)

- Calculates all first- and second-order transport coefficients from the linearized Boltzmann equation
- Multi-component systems with isotropic (in-)elastic, s -dependent cross sections
- thermal masses and cross sections possible (e.g. DQPM)
- quantum corrections are planned

Upcoming publication!

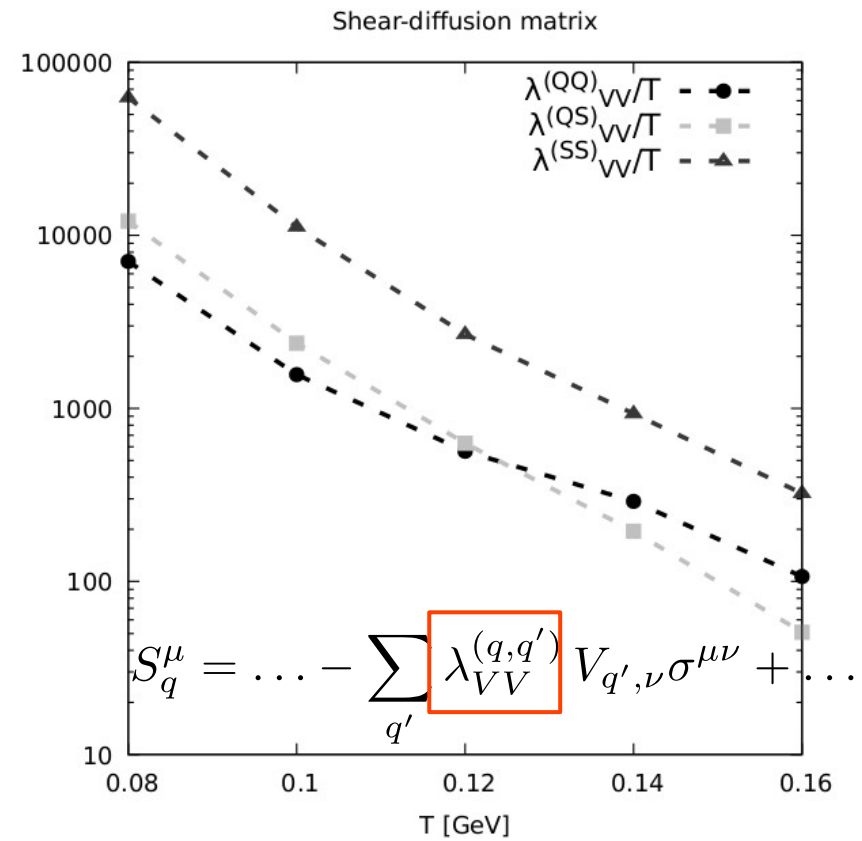
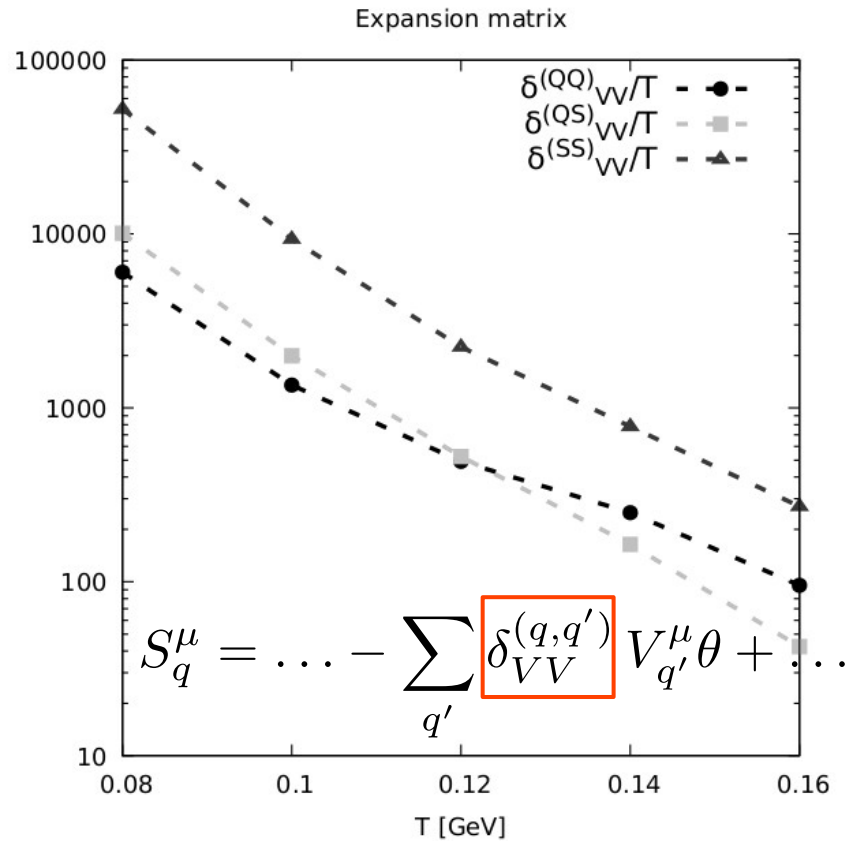
(Fotakis, Lohr, Wagner, Potesnov et al.)

Pion-Kaon gas with resonant cross sections (ρ , K^*) with hard-sphere “background”



Preliminary results from TraCoLinR

Pion-Kaon gas with resonant cross sections (ρ , K^*) with hard-sphere “background”



Conclusion

- Derived 2nd-order relativistic fluid dynamic theory for **multi-component systems** from the Boltzmann equation
- **Transport coefficients given explicitly** containing all information about particle interactions
- Mixed chemistry correlates particle flow → **coupled charge-transport**
- **Consistency** of equation of state, 1st- and 2nd-order transport coefficients **is important!**
- Implemented derived fluid dynamic theory in **(3+1)D-hydro code HYDRA**
- First preliminary calculations of **second-order coefficients** of massive mixture with more realistic cross sections **with TraCoLinR**
- **We plan to publish both codes**

Backup



Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

$10 + 4N_{\text{ch}}$ degrees of freedom, $4 + N_{\text{ch}}$ equations \rightarrow $6 + 3N_{\text{ch}}$ unknowns

q-th conserved charge (eg. B,Q,S)

Conservation of charge: $\partial_\mu N_q^\mu = 0$

$$N_q^\mu = \sum_i q_i N_i^\mu = n_q u^\mu + V_q^\mu$$

What needs to be known:

- Equation of state $P_0 = P_0(\epsilon, n_q)$, $T = T(\epsilon, n_q)$, $\alpha_q = \mu_q/T = \alpha_q(\epsilon, n_q)$
- Equations of motion for dissipative fields & transport coefficients $\Pi, V_q^\mu, \pi^{\mu\nu}$
- Initial state
- Final state: freeze-out and δf -correction

Deriving fluid dynamics from kinetic theory

Denicol et al., PRD 85, 114047 (2012)

On basis of DNMR theory: derivation from the Boltzmann equation with method of moments
Fotakis, Molnar, Niemi, Rischke, Greiner, Phys. Rev. D 106 (2022), 036009

Also refer to: Monnai et al., Nucl. Phys. A847:283-314 (2010) or Kikuchi et al., PRC 92, 064909 (2015)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$



2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

Irreducible **off-equilibrium** moments
obey Boltzmann eq.:

Problem: infinitely many coupled PDEs.

Aim: Truncate in a well-defined manner (“perturbation theory”)

$$f_{i,\mathbf{k}} = \overset{\text{equilibrium}}{f_{i,\mathbf{k}}^{(0)}} + \overset{\text{off-equilibrium}}{\delta f_{i,\mathbf{k}}}$$
$$\rho_{i,n}^{\mu\nu} = \int \frac{d^3\mathbf{k}}{(2\pi)^3 E_{i,\mathbf{k}}} E_{i,\mathbf{k}}^n k_i^{\langle\mu} k_i^{\nu\rangle} \delta f_{i,\mathbf{k}}$$

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Aim: Truncate in a well-defined manner (“perturbation theory”)

“Order-of-magnitude approximation”:
relate off-equilibrium moments to the dissipative fields

$$\rho_{i,n}^{\mu\nu} = \frac{\eta_{i,n}}{\eta} \pi^{\mu\nu} + \mathcal{O}(2)$$

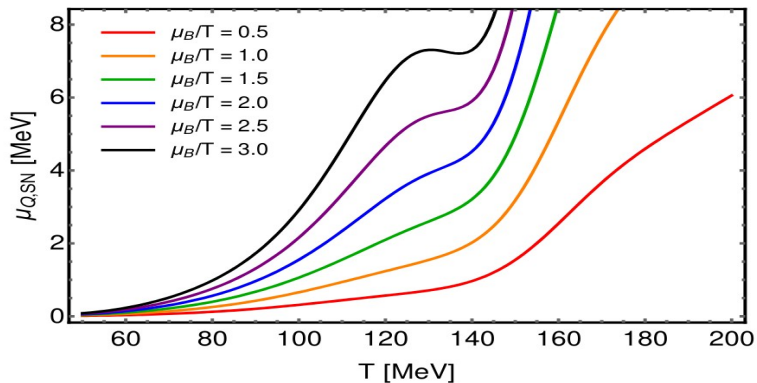
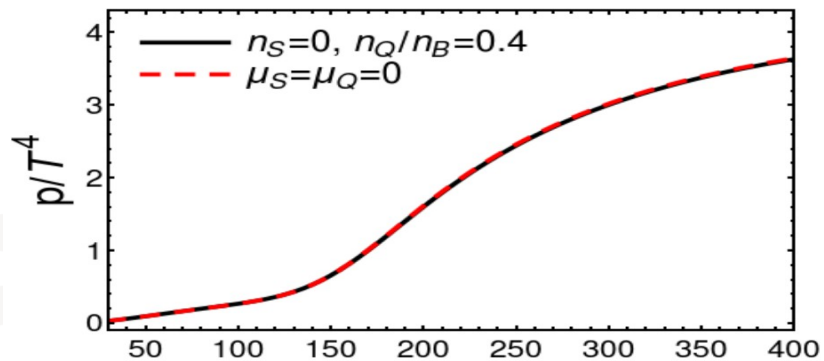
Counting scheme:

Gradients in velocity, temperature etc. $\sigma^{\mu\nu} \sim \mathcal{O}(1), \mathcal{O}(\text{Kn})$
Dissipative fields $\pi^{\mu\nu} \sim \mathcal{O}(1), \mathcal{O}(\text{Rn}^{-1})$

Equation of state with multiple conserved charges

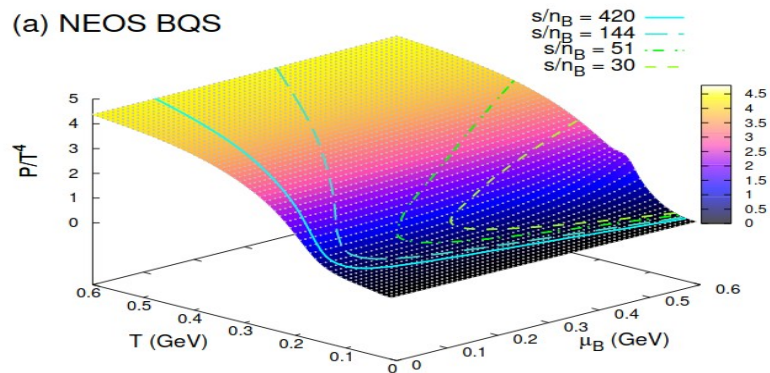
$$P_0(T) \rightarrow P_0(T, \mu_B, \mu_Q, \mu_S)$$

Noronha-Hostler et al., PRC 100, 064910 (2019)

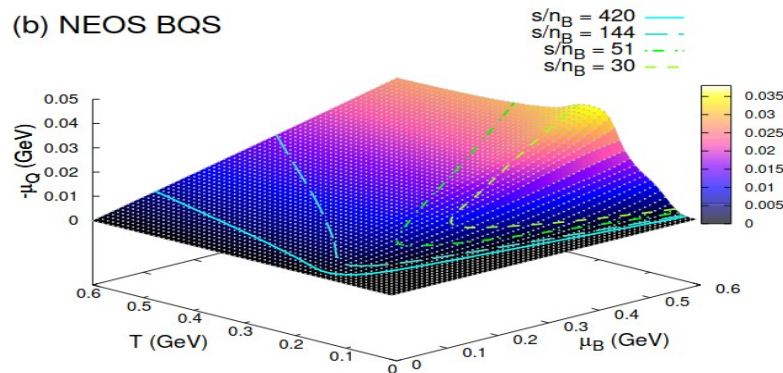


Monnai et al., PRC 100, 024907 (2019)

(a) NEOS BQS



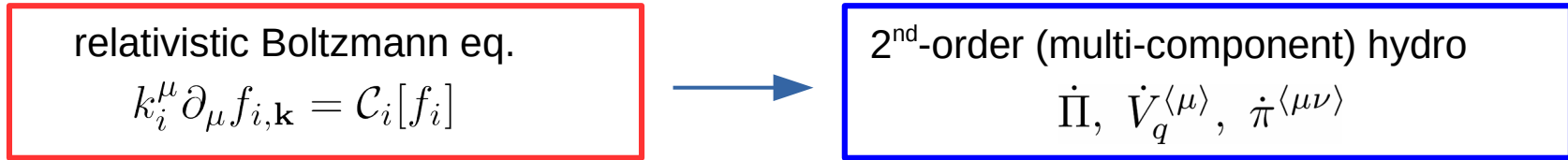
(b) NEOS BQS



Computation of transport coefficients (Example: diffusion coefficients)

On basis of DNMR theory: derivation from the Boltzmann equation with method of moments

Fotakis, Molnar, Niemi, Rischke, Greiner, Phys. Rev. D 106 (2022), 036009



$$C_{i,n-1}^{\langle\mu\rangle} \equiv \int \frac{d^3\mathbf{k}_i}{(2\pi)^3 E_{i,\mathbf{k}}} E_{i,\mathbf{k}}^{n-1} k_i^{\langle\mu\rangle} C_i[f_i]$$

$$= - \sum_{m=0}^{\infty} \sum_j C_{ij,nm}^{(1)} \rho_{j,m}^\mu + \text{non-linear terms}$$

Entries of „collision matrix“ (for diffusive moments)

$$\kappa_{qq'} = \sum_{n=0}^{\infty} \sum_{i,j=1}^{N_{\text{species}}} \left(C^{(1)} \right)_{ij,0n}^{-1} q_i \left(q'_j J_{j,n+1,1} - \frac{n_{q'}}{\epsilon + P_0} J_{j,n+2,1} \right)$$

Diffusion coefficient matrix! (equivalent to our PRL and PRD expression)

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)
Fotakis, Soloveva et al, PRD 104, 034014 (2021)

Yet another hydro code - „Hydra“

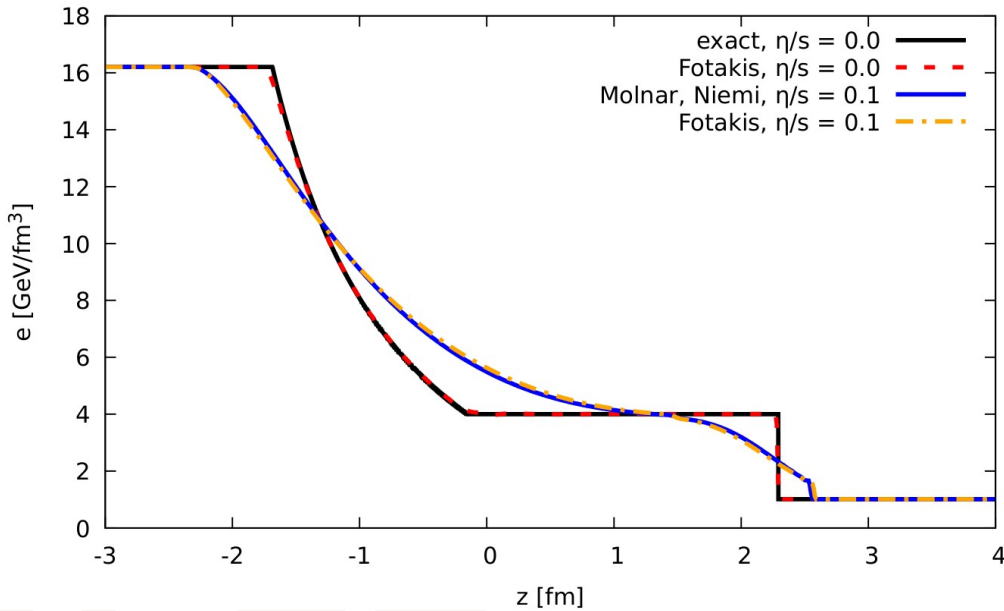
Core features:

- (3+1)D-hydro – optimized reduction to 2D and 1D
- (v)SHASTA solver
- Shear-stress and multiple conserved charges (2 charges)
- Ultrarelativistic, tabled and/or any user-defined equations of state
- DNMR theory, this theory, and/or any user-defined theory
- any (tabled, user-defined) transport coefficients
- Curve-linear geometry (so far Cartesian and Hyperbolic coordinates)
- state of the art unit and physical tests
- available in the CRC-TR211 collaboration soon (hopefully)

Yet another hydro code - „Hydra“

Core features:

- (3+1)D-hydro – optimized reduction to 2D and 1D
- (v)SHASTA solver



charges (2 charges)

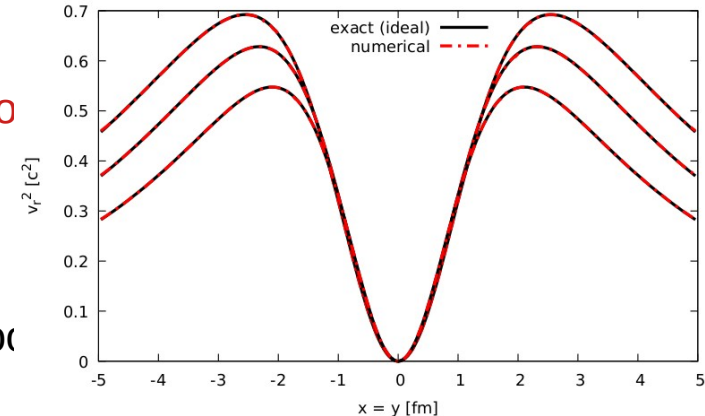
refined equations of

re-defined theory

coefficients

and Hyperbolic coc

is soon (hopefully)



- Hadronic system including lightest 19 species

$$\pi^\pm, \pi^0, K^\pm, K^0, \bar{K}^0, p, \bar{p}, n, \bar{n}, \Lambda^0, \bar{\Lambda}^0, \Sigma^0, \bar{\Sigma}^0, \Sigma^\pm, \bar{\Sigma}^\pm$$

- Assume classical statistics and non-interacting limit

$$P_0(T, \{\mu_q\}) \equiv \frac{1}{3} \sum_{i=1}^{N_{\text{species}}} \int \frac{dp^3}{(2\pi)^3 E_{i,p}} (E_{i,p}^2 - m_i^2) g_i \exp(-E_{i,p}/T + \sum_q q_i \alpha_q)$$

- Only assume **baryon number** and **strangeness**, **neglect electric charge**

- Tabulate state variables over energy density and net charge densities

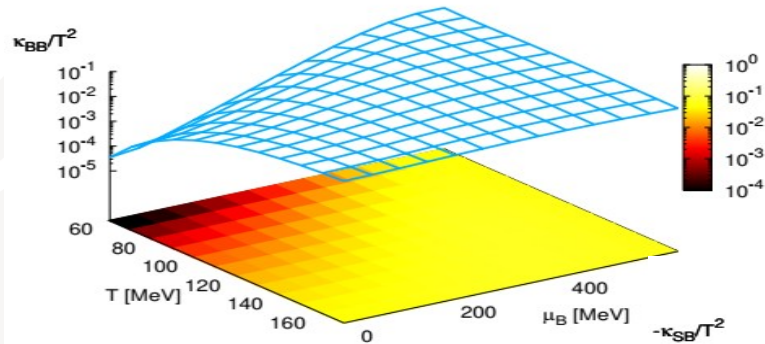
$$T \equiv T(\epsilon, \{n_q\}), \quad \mu_q \equiv \mu_q(\epsilon, \{n_q\}), \quad P_0 \equiv P_0(\epsilon, \{n_q\})$$

Diffusion coefficient matrix - details

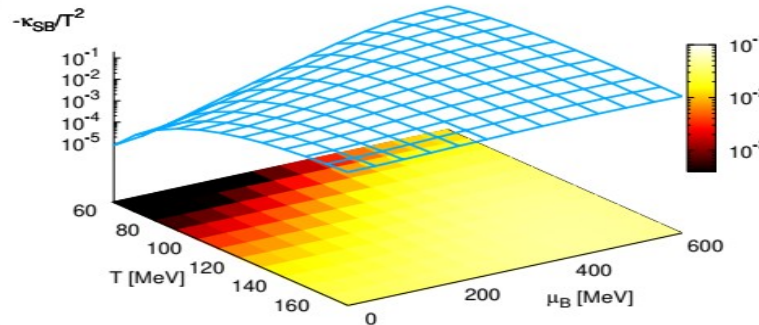
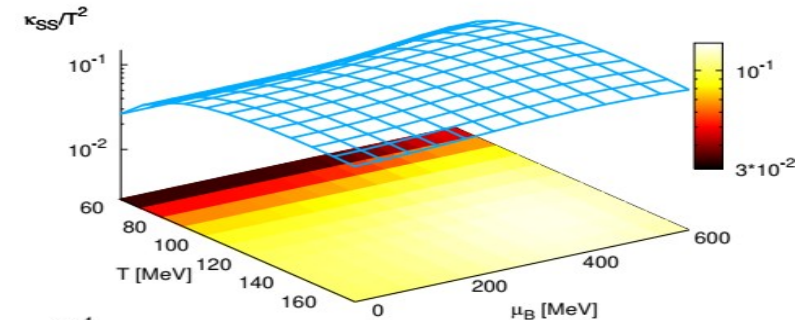
$$\begin{pmatrix} V_B^\mu \\ V_S^\mu \end{pmatrix} \sim \begin{pmatrix} \kappa_{BB} & \kappa_{BS} \\ \kappa_{SB} & \kappa_{SS} \end{pmatrix} \begin{pmatrix} \nabla^\mu \alpha_B \\ \nabla^\mu \alpha_S \end{pmatrix}$$

- Matrix is symmetric

L. Onsager, Phys. Rev. **37**, 405 (1931) & Phys. Rev. **38**, 2265 (1931)



- Elastic isotropic cross sections from PDG, SMASH, GiBUU, UrQMD



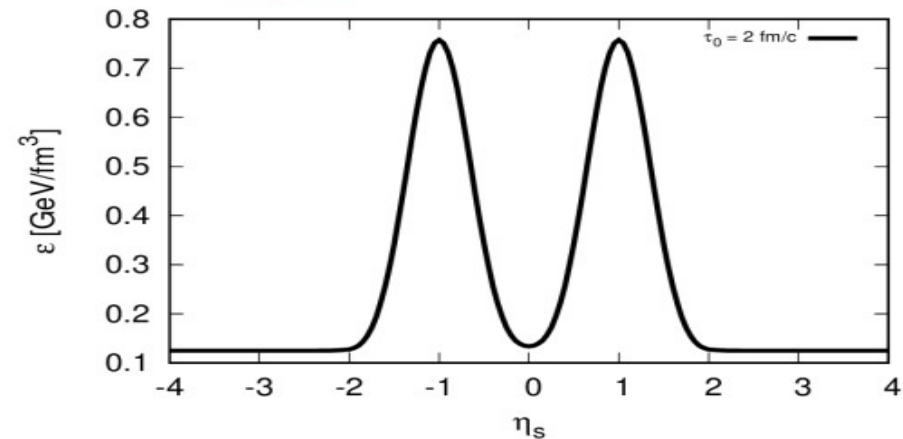
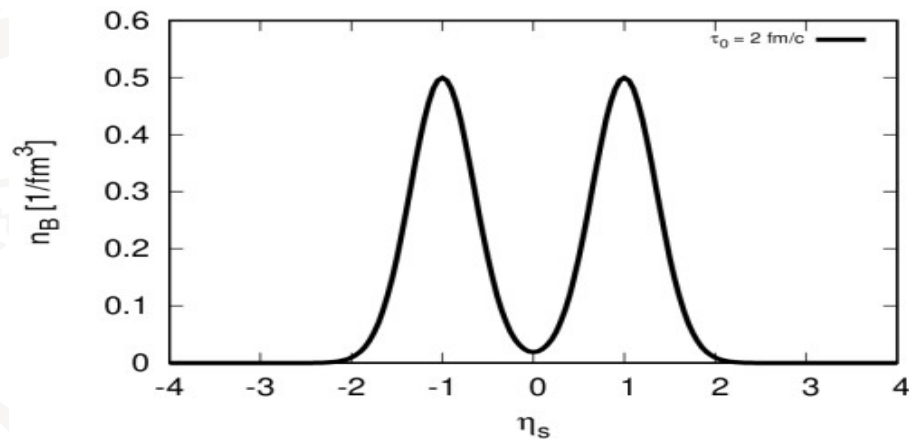
κ_{SB} is **negative** and has **similar magnitude** as κ_{BB}

⇒ significant coupling?

- Tabulate coefficient matrix over T, μ_B, μ_S

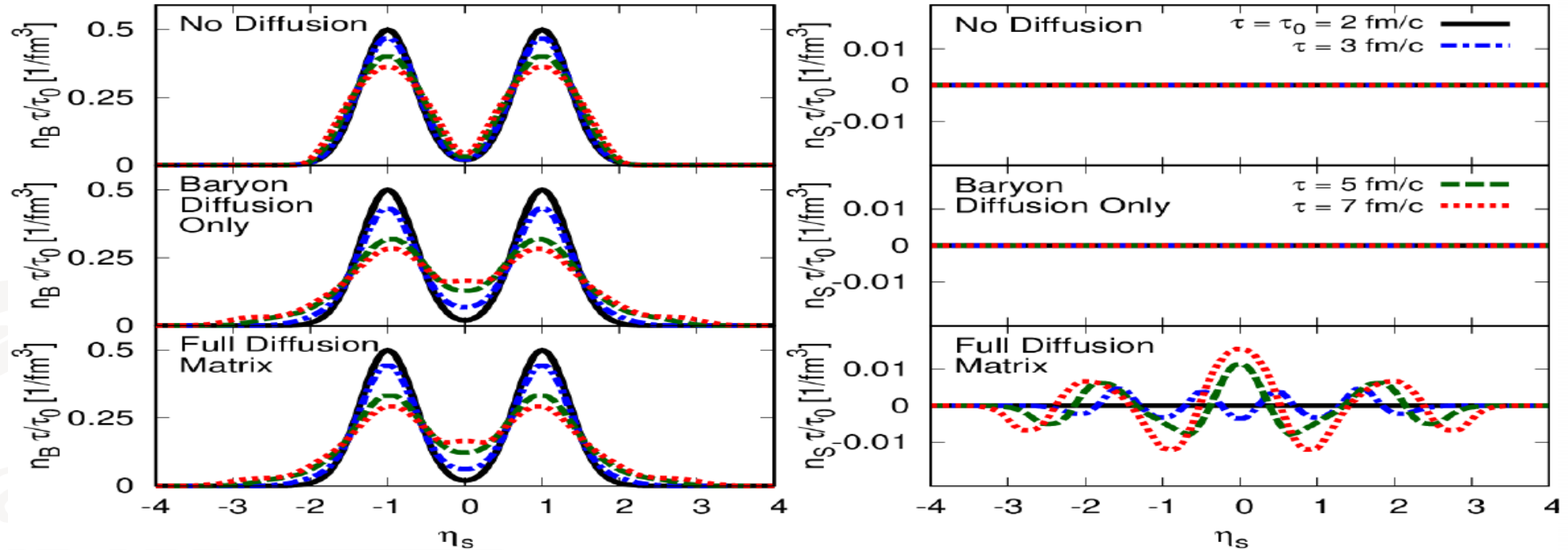
- $\mu_Q = 0$

- $\tau_0 = 2 \text{ fm}/c$
- Initially: no dissipation and only **Bjorken scaling flow**
- Temperature = 160 MeV
- Double-gaussian profile in net baryon number
- From **EoS**: get energy density



Coupled charge-transport

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)



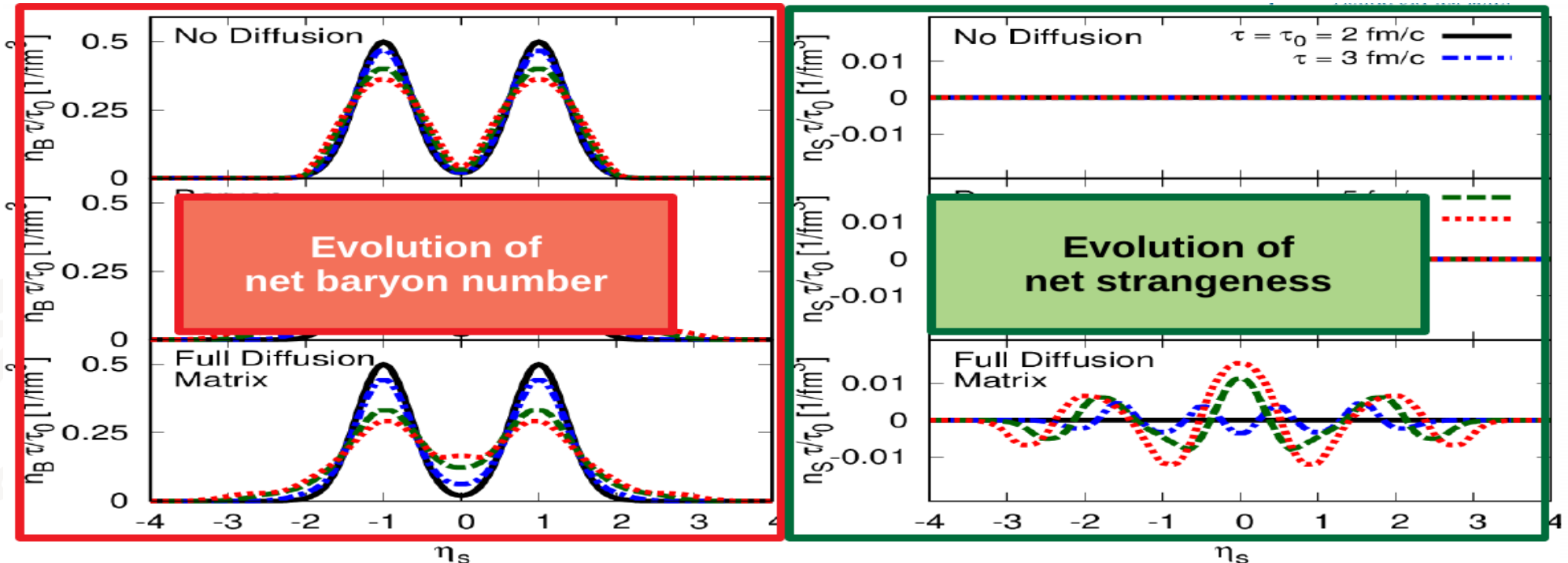
Simplistic case study: no viscosity, diffusion only, no 2nd-order terms, conserved B and S charge, classical, hadronic system (19 species), realistic binary elastic cross sections

Hydrodynamic (1+1)D-simulation

$$\Pi \equiv 0, \quad \pi^{\mu\nu} \equiv 0, \quad \tau_q \dot{V}_q^{\langle\mu} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \left(\frac{\mu_{q'}}{T} \right)$$

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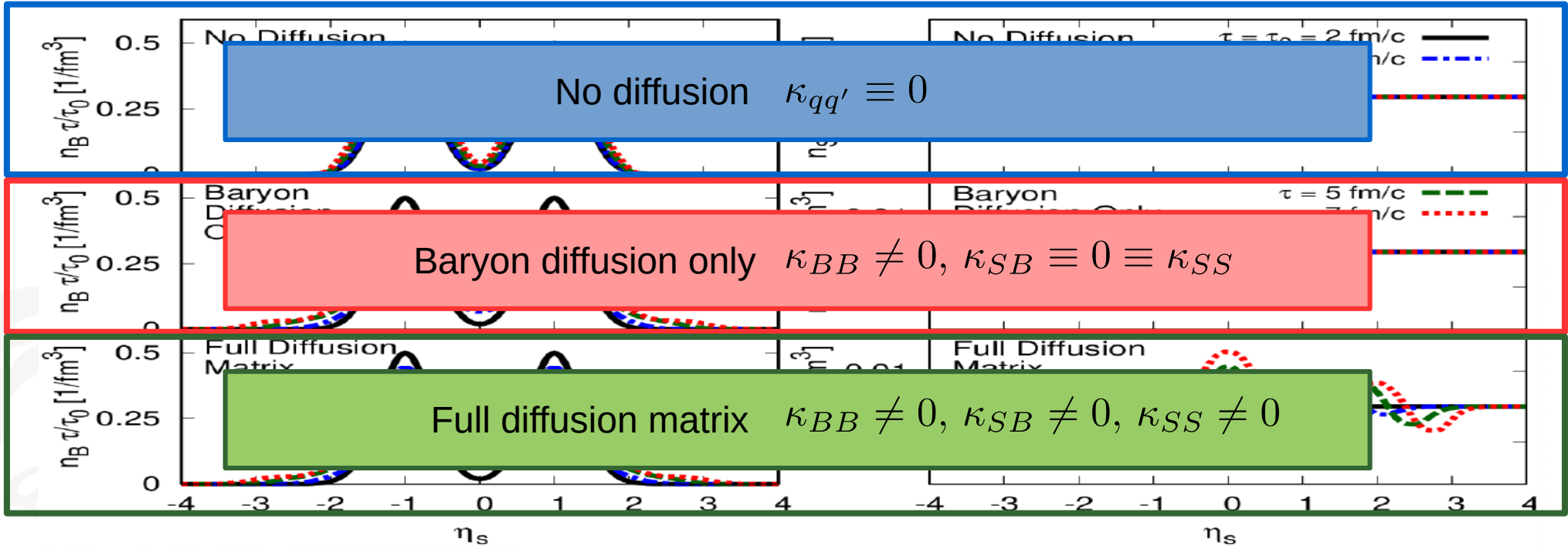
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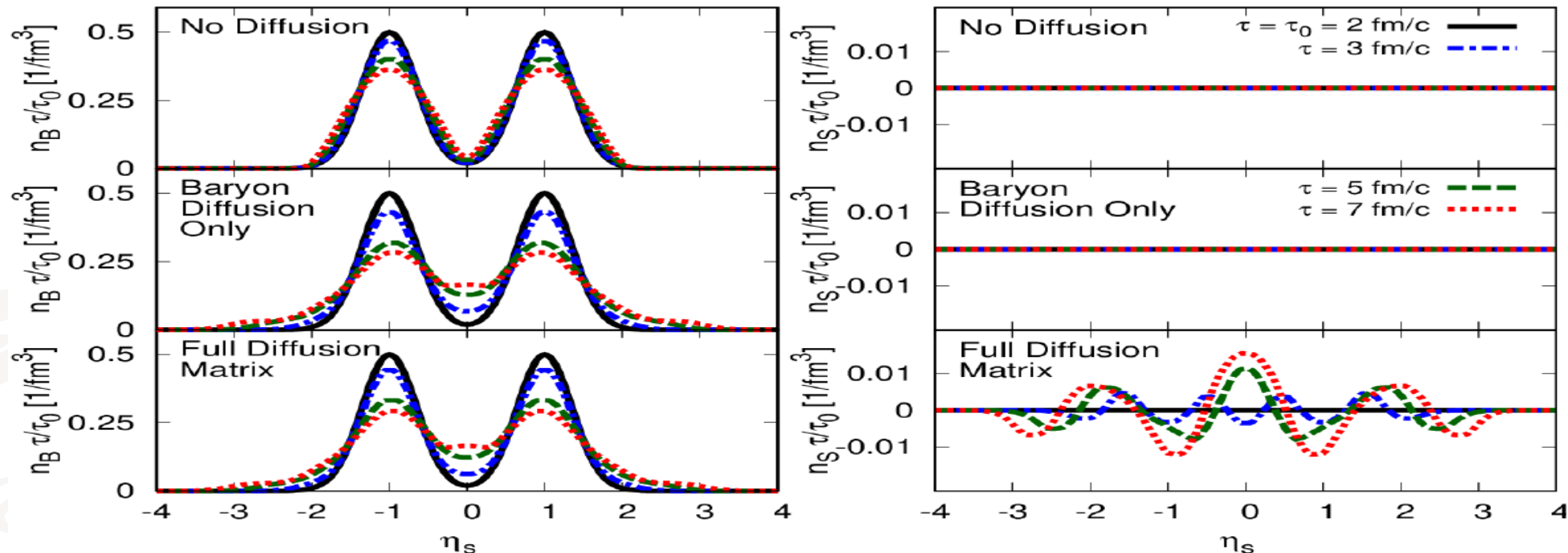
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Coupled charge-transport

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Mixed chemistry couples diffusion currents and introduces charge-correlation through EoS



Generation of domains of non-vanishing local net charge (here net strangeness)!

e.g.: $\mu_S \equiv \mu_S(\epsilon, n_B, n_S)$
 $\nabla^\mu \alpha_S \sim \nabla^\mu n_B$